

Matrix-valued finite elements

Jay Gopalakrishnan

Portland State University

In collaboration with

Kaibo Hu, Philip Lederer, Michael Neunteufel, Joachim Schöberl & others

FEM@LLNL Seminar, Livermore, CA, January 20, 2026

Thanks:



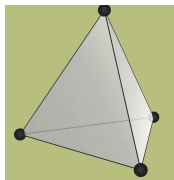
DMS

Outline

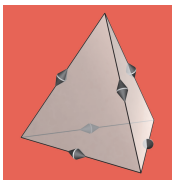
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements

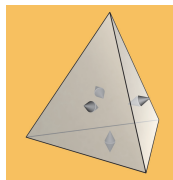
($\mathbb{R}$ -valued)



($\mathbb{V}$ -valued)



($\mathbb{V}$ -valued)



($\mathbb{R}$ -valued)

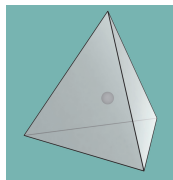
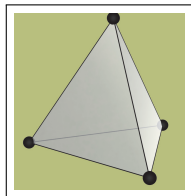


Figure Source: [Arnold+Logg 2014] poster "The Periodic Table of Finite Elements."

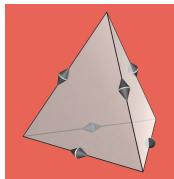
Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements

($\mathbb{R}$ -valued)

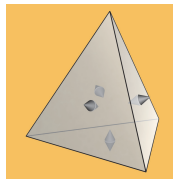


grad, continuous

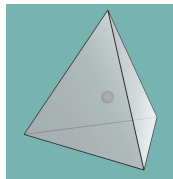
($\mathbb{V}$ -valued)



($\mathbb{V}$ -valued)



($\mathbb{R}$ -valued)



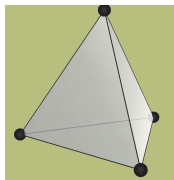
Lagrange finite element space [Courant 1943]

Natural operator: grad

Continuous across elements

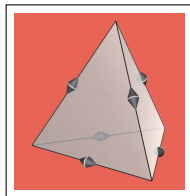
Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements

($\mathbb{R}$ -valued)



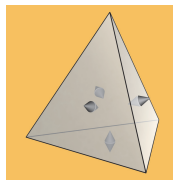
grad, continuous

($\mathbb{V}$ -valued)

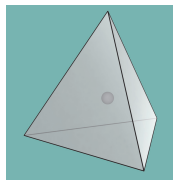


curl, " t -continuous"

($\mathbb{V}$ -valued)



($\mathbb{R}$ -valued)



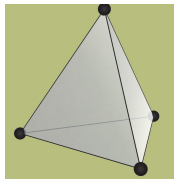
Nédélec finite element space [Nédélec 1980]

Natural operator: curl

Inter-element continuity: t -component

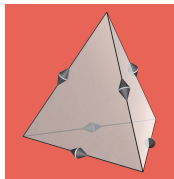
Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements

($\mathbb{R}$ -valued)



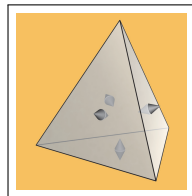
grad, continuous

($\mathbb{V}$ -valued)



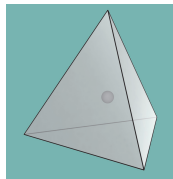
curl, " t -continuous"

($\mathbb{V}$ -valued)



div, " n -continuous"

($\mathbb{R}$ -valued)



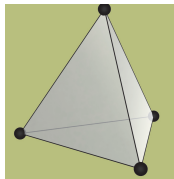
Raviart-Thomas finite element space [Raviart+Thomas 1977]

Natural operator: **div**

Inter-element continuity: **n -component**

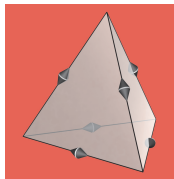
Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements

($\mathbb{R}$ -valued)



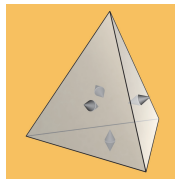
grad, continuous

($\mathbb{V}$ -valued)



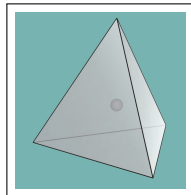
curl, " t -continuous"

($\mathbb{V}$ -valued)



div, " n -continuous"

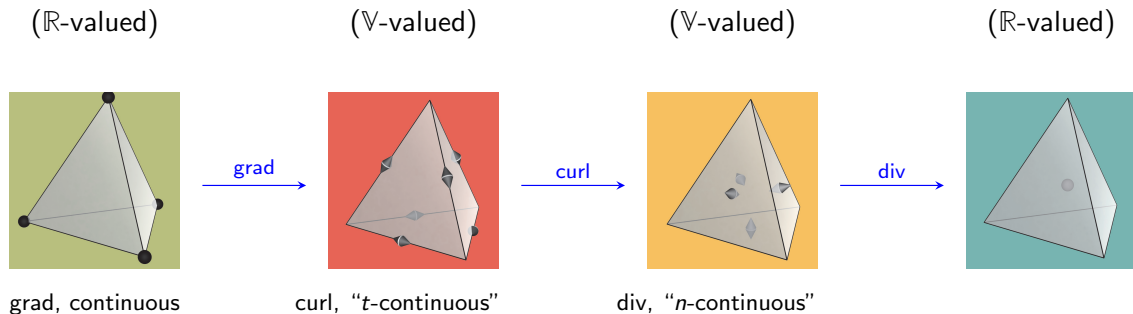
($\mathbb{R}$ -valued)



Discontinuous Galerkin finite element space

Inter-element continuity: None

Lowest order scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued finite elements



These finite element spaces are now collectively understood using the [de Rham complex](#).

Where do matrix-valued finite elements fit? And second-order differential operators?

Next example: **curl div** operator and "*nt-continuous*" matrix elements for viscous stresses.

Stokes system

Textbook version:

$$\begin{aligned} -\nu \Delta u + \nabla p &= f && \text{in } \Omega \\ \operatorname{div} u &= 0 && \text{in } \Omega \\ u &= 0 && \text{on } \partial\Omega \end{aligned}$$

Physical stress-based version:

$$\begin{aligned} \frac{1}{2\nu} \sigma - \varepsilon(u) &= 0 && \text{in } \Omega \\ \operatorname{div} \sigma - \nabla p &= -f && \text{in } \Omega \\ \operatorname{div} u &= 0 && \text{in } \Omega \\ u &= 0 && \text{on } \partial\Omega \end{aligned}$$

Notation:

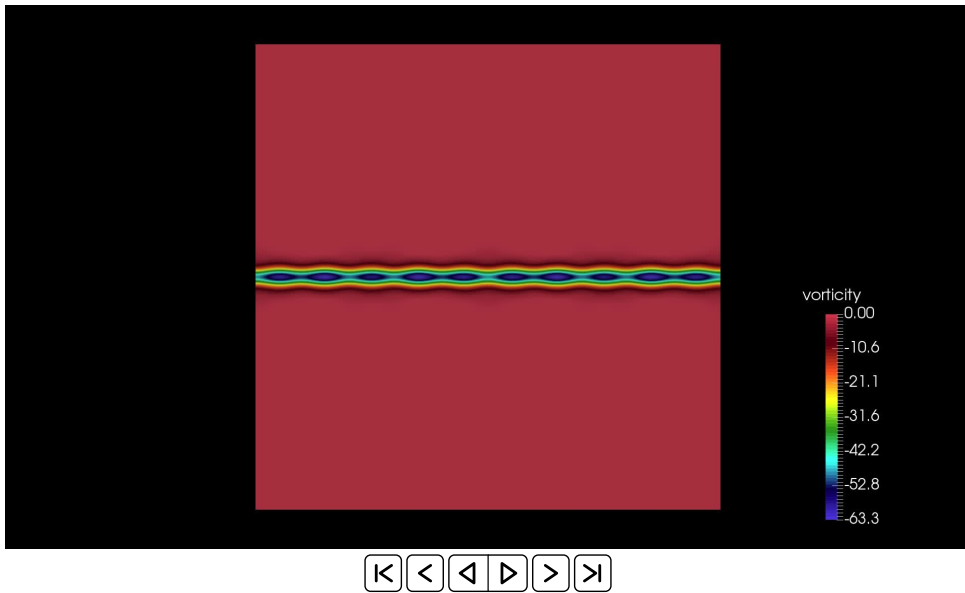
- $\nu(x)$ = kinematic viscosity
- $f(x)$ = body force
- $u(x)$ = fluid velocity
- $p(x)$ = kinematic pressure
- $\sigma(x)$ = viscous stress
- $\varepsilon(u) = \frac{1}{2}(\nabla u + (\nabla u)') = \operatorname{sym}(\nabla u)$.

What topology on σ pairs well with $H(\operatorname{div})$ -topology on u ?

Outline

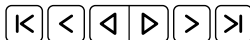
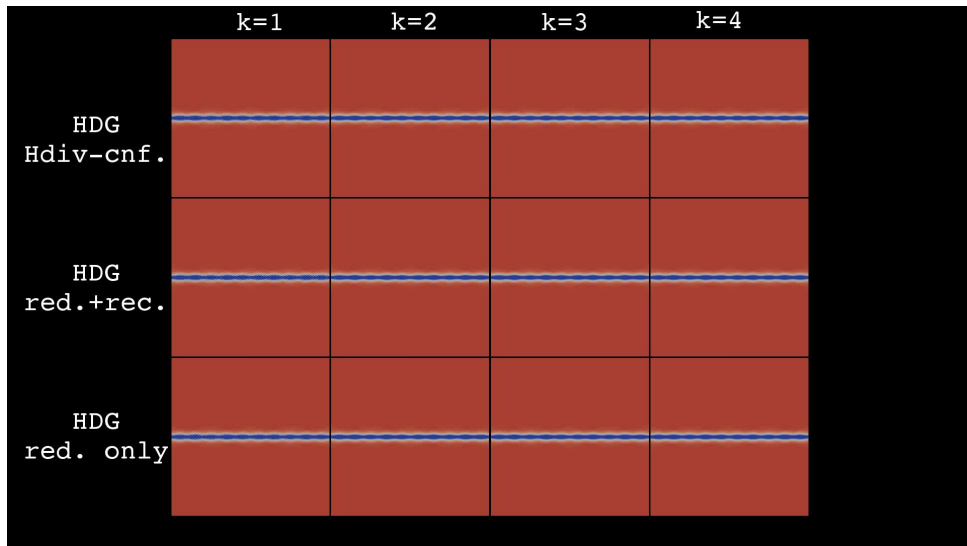
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method**
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

Kelvin-Helmholtz instability simulation



[Schroeder+John+Lederer+Lehrenfeld+Lube+Schöberl 2019]

Different results from different discretizations



Source: P. Lederer's thesis

Mass conservation

Mass is neither created nor destroyed.

- Mass = $\int \rho$, density integrated.
- Incompressibility $\implies$ the material derivative $\frac{D\rho}{Dt} = 0$.

Equivalent constraint on fluid velocity u :

- Continuity equation (mass conservation) $\implies \frac{D\rho}{Dt} + \rho \operatorname{div} u = 0$.
- Hence, mass conservation in incompressible fluids $\iff \boxed{\operatorname{div} u = 0}$.

Mass conservation

Mass is neither created nor destroyed.

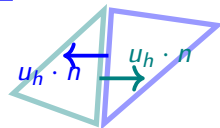
- Mass = $\int \rho$, density integrated.
- Incompressibility $\implies$ the material derivative $\frac{D\rho}{Dt} = 0$.

Equivalent constraint on fluid velocity u :

- Continuity equation (mass conservation) $\implies \frac{D\rho}{Dt} + \rho \operatorname{div} u = 0$.
- Hence, mass conservation in incompressible fluids $\iff \boxed{\operatorname{div} u = 0}$.

A piecewise polynomial fluid velocity approximation u_h has $\boxed{\operatorname{div} u_h = 0}$ if

- its divergence is zero pointwise within elements,
- and $\llbracket u_h \cdot n \rrbracket = 0$, i.e., u_h has at least *n-continuity*.



How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

$$(\operatorname{div} u_h, q_h) = 0 \quad \text{for all } q_h \text{ in the discrete pressure space } Q_h.$$

Hence u_h is exactly mass conserving if $\operatorname{div} V_h \subseteq Q_h$.

Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

H^1 -based

How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

$$(\operatorname{div} u_h, q_h) = 0 \quad \text{for all } q_h \text{ in the discrete pressure space } Q_h.$$

Hence u_h is exactly mass conserving if $\operatorname{div} V_h \subseteq Q_h$.

Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

$$H(\operatorname{div}) = \{u : u_i \in L_2, \sum_i \partial_i u_i \in L_2\}$$

H^1 -based

$$H^1 \otimes \mathbb{V} = \{u : u_i \in L_2, \partial_j u_i \in L_2\}$$

(This is vector H^1 .)

How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

$$(\operatorname{div} u_h, q_h) = 0 \quad \text{for all } q_h \text{ in the discrete pressure space } Q_h.$$

Hence u_h is exactly mass conserving if $\operatorname{div} V_h \subseteq Q_h$.

Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

When $V_h \subset H(\operatorname{div})$:

- $\oplus$ Range of divergence is simple.
- $\ominus$ Discretizing viscous term is complex.

H^1 -based

When $V_h \subset H^1 \otimes \mathbb{V}$:

- $\ominus$ Range of divergence is complex.
- $\oplus$ Viscous term easily discretized.

How can methods gain exact mass conservation?

Discrete version of the incompressibility constraint in a space V_h with n -continuity:

$$(\operatorname{div} u_h, q_h) = 0 \quad \text{for all } q_h \text{ in the discrete pressure space } Q_h.$$

Hence u_h is exactly mass conserving if $\operatorname{div} V_h \subseteq Q_h$.

Two approaches to methods with mass conservation:

$H(\operatorname{div})$ -based

[Cockburn+Kanschat+Schötzau 2005]

⋮

[Cockburn+G+Nguyen+Peraire+Sayas 2011]

[Linke+Merdon 2016]

[G+Lederer+Schöberl 2019] (MCS method)

H^1 -based

[Scott+Vogelius 1985]

⋮

[Guzmán+Scott 2019]

[Neilan 2020]

[Ainsworth+Parker 2021]

Punchlines of Mass-Conserving Stress-yielding (MCS) method

- *Structure-preservation:*
Numerical solutions are exactly mass conserving.
Scheme is pressure robust.
- *Viscous stress σ approximation:* New matrix-valued finite elements with *continuous shear component* (or *nt-component*) used for approximating σ . ($\llbracket t \cdot \sigma n \rrbracket \equiv \llbracket \sigma_{nt} \rrbracket = 0$.)
- *Optimally convergent in* velocity (u), pressure (p), viscous stress (σ), and vorticity (ω).
- *More features:*
 - ▶ No stabilization parameters.
 - ▶ Easy to incorporate stress boundary conditions.
 - ▶ No vertex unknowns, facet couplings only: permits easy hybridization.
 - ▶ Element mappings (even curvilinear) are straightforward.
 - ▶ Stable for ∇u -forms as well as the more physical $\varepsilon(u)$ -formulations.

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \operatorname{grad} p = 0$$

$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \sigma - \varepsilon(u) = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$u = 0 \quad \text{on } \partial\Omega$$

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \operatorname{grad} p = 0$$

$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \operatorname{dev} \sigma - \varepsilon(u) = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$u = 0 \quad \text{on } \partial\Omega$$

Notation:

$$\operatorname{dev} m := m - \frac{1}{3}(\operatorname{tr} m)\delta, \quad m \text{ in } \mathbb{R}^{3 \times 3}$$

$$\sigma = \operatorname{dev} \sigma, \text{ in } \mathbb{T}$$

$$\mathbb{T} := \operatorname{dev}(\mathbb{R}^{3 \times 3}) \equiv \{\text{traceless matrices}\}$$

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \operatorname{grad} p = 0$$

$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \operatorname{dev} \sigma - \nabla u + \omega = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$\sigma - \sigma' = 0 \quad \text{in } \Omega$$

$$u = 0 \quad \text{on } \partial\Omega$$

Notation:

$$\operatorname{dev} m := m - \frac{1}{3}(\operatorname{tr} m)\delta,$$

$$\sigma = \operatorname{dev} \sigma, \text{ in } \mathbb{T}$$

$$\omega := \text{vorticity, in } \mathbb{K}$$

$$m \text{ in } \mathbb{R}^{3 \times 3},$$

$$\mathbb{T} := \operatorname{dev}(\mathbb{R}^{3 \times 3}) \equiv \{\text{traceless matrices}\}$$

$$\mathbb{K} := \operatorname{skw}(\mathbb{R}^{3 \times 3}) \equiv \{\text{skew-symmetric matrices}\}$$

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \operatorname{grad} p = 0$$

$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \operatorname{dev} \sigma - \nabla u + \omega = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$\sigma - \sigma' = 0 \quad \text{in } \Omega$$

$$u = 0 \quad \text{on } \partial\Omega$$

Test with $\tau : \Omega \rightarrow \mathbb{T}$, $v : \Omega \rightarrow \mathbb{V}$, $q : \Omega \rightarrow \mathbb{R}$, $\eta : \Omega \rightarrow \mathbb{K}$:

$$\left(\frac{1}{2\nu} \operatorname{dev} \sigma, \operatorname{dev} \tau \right) + (u, \operatorname{div} \tau) + (\omega, \tau) = 0$$

$$(\operatorname{div} \sigma, v) + (p, \operatorname{div} v) = -(f, v)$$

$$(\operatorname{div} u, q) = 0$$

$$(\sigma, \eta) = 0.$$

Deriving the MCS formulation

Standard form:

$$-\nu \Delta u + \operatorname{grad} p = 0$$

$$\operatorname{div} u = 0$$

$$u = 0$$

With σ :

$$\frac{1}{2\nu} \operatorname{dev} \sigma - \nabla u + \omega = 0 \quad \text{in } \Omega$$

$$\operatorname{div} \sigma - \operatorname{grad} p = -f \quad \text{in } \Omega$$

$$\operatorname{div} u = 0 \quad \text{in } \Omega$$

$$\sigma - \sigma' = 0 \quad \text{in } \Omega$$

$$u = 0 \quad \text{on } \partial\Omega$$

Test with $\tau : \Omega \rightarrow \mathbb{T}$, $v : \Omega \rightarrow \mathbb{V}$, $q : \Omega \rightarrow \mathbb{R}$, $\eta : \Omega \rightarrow \mathbb{K}$:

$$\left(\frac{1}{2\nu} \operatorname{dev} \sigma, \operatorname{dev} \tau \right) + (u, \operatorname{div} \tau) + (\omega, \tau) = 0$$

$$(\operatorname{div} \sigma, v) + (p, \operatorname{div} v) = -(f, v)$$

$$(\operatorname{div} u, q) = 0$$

$$(\sigma, \eta) = 0.$$

Derivation is completed by replacing the integrals $(\operatorname{div} \sigma, v)$ by **functional actions** $\langle \operatorname{div} \sigma, v \rangle$.

The viscous stress space

MCS formulation: Find $\sigma \in \Sigma$, $u \in \dot{H}(\text{div})$, $p \in L_{2,\mathbb{R}}$, $\omega \in L_2 \otimes \mathbb{K}$ satisfying

$$\begin{aligned} \left(\frac{1}{2\nu} \text{dev } \sigma, \text{dev } \tau\right) + \langle \text{div } \tau, u \rangle + (\omega, \tau) &= 0, & \tau \in \Sigma, \\ \langle \text{div } \sigma, v \rangle + (p, \text{div } v) &= -\langle f, v \rangle, & v \in \dot{H}(\text{div}), \\ (\text{div } u, q) &= 0, & q \in L_{2,\mathbb{R}}, \\ (\sigma, \eta) &= 0, & \eta \in L_2 \otimes \mathbb{K}. \end{aligned}$$

Minimal requirement for any σ in Σ is $\text{div } \sigma \in \dot{H}(\text{div})^*$, the dual space of $\dot{H}(\text{div})$.

Theorem: $\dot{H}(\text{div})^* = H^{-1}(\text{curl})$.

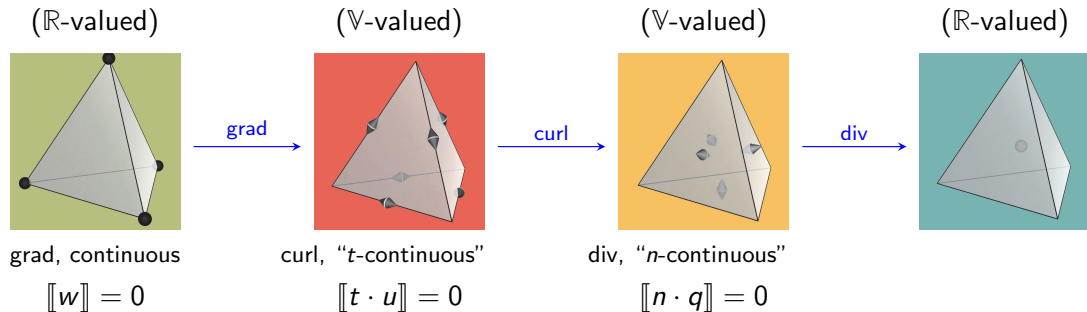
Space for viscous fluid stresses: $\Sigma = \{\tau \in L_2 \otimes \mathbb{T} : \text{curl div}(\tau) \in H^{-1}\}$.

Notation: ① $\mathbb{T} = \{\text{traceless matrices}\}$, ② $\mathbb{K} = \{\text{skew symmetric matrices}\}$, ③ $L_{2,\mathbb{R}} = \{p \in L_2 : (p, 1) = 0\}$,
④ $H^{-1}(\text{curl}) = \{\phi \in H^{-1} \otimes \mathbb{V} : \text{curl } \phi \in H^{-1} \otimes \mathbb{V}\}$, ⑤ $\dot{H}(\text{div}) = \{w \in H(\text{div}) : w \cdot n|_{\partial\Omega} = 0\}$.

Outline

- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements**
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

The *nt*-continuity



For the **curl div** operator, we make “*nt*-continuous” matrix elements for viscous stresses, i.e.,

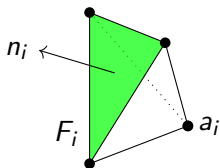
$$[[t \cdot \sigma n]] = 0$$

for any tangent vector t and normal vector n on element boundaries.

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n_t -component vanishes on ∂T except on one face?*

Tetrahedron T

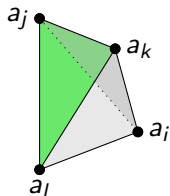


- Consider $\sigma = n_j \otimes (n_k \times n_l)$

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n -component vanishes on ∂T except on one face?*

Tetrahedron T

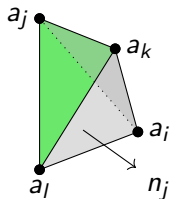


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- $\sigma n = 0$ on $F_k \cup F_l$

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n -component vanishes on ∂T except on one face?*

Tetrahedron T

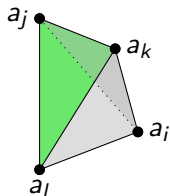


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- $\sigma n = 0$ on $F_k \cup F_l$ and σn_j is collinear with n_j , so $\sigma n_j \cdot t|_{F_j} = 0$.

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n -component vanishes on ∂T except on one face?*

Tetrahedron T

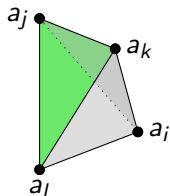


- Consider $\sigma = n_j \otimes (n_k \times n_l)$
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n -component vanishes on ∂T except on one face?*

Tetrahedron T

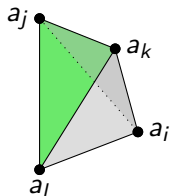


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose n -component vanishes on ∂T except on one face?*

Tetrahedron T

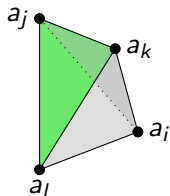


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$ and $\sigma = \text{dev}[n_l \otimes (n_j \times n_k)]$.
- Thus $t \cdot \sigma n = 0$ on all faces except F_i .

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose nt -component vanishes on ∂T except on one face?*

Tetrahedron T

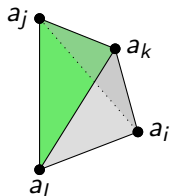


- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_l)]$ and $\sigma = \text{dev}[n_l \otimes (n_j \times n_k)]$.
- Taking all faces, we get a basis for $\mathbb{T} = \{m \in \mathbb{R}^{3 \times 3} : \text{tr } m = 0\}$.

Designing finite elements for viscous stresses

Lowest order case: *Is it possible to construct a constant matrix function whose nt -component vanishes on ∂T except on one face?*

Tetrahedron T



- Consider $\sigma = \text{dev}[n_j \otimes (n_k \times n_i)]$ and $\sigma = \text{dev}[n_i \otimes (n_j \times n_k)]$.
- Taking all faces, we get a basis for $\mathbb{T} = \{m \in \mathbb{R}^{3 \times 3} : \text{tr } m = 0\}$.

More general shape functions: $\lambda_i^{\alpha_i} \lambda_j^{\alpha_j} \lambda_k^{\alpha_k} \text{dev}[\nabla \lambda_i \otimes (\nabla \lambda_j \times \nabla \lambda_k)]$.

Unisolvent finite element

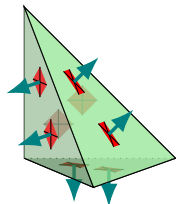
Ciarlet-style finite element definition:

- Geometry: Tetrahedron T
- Space: $P_k(T) \otimes \mathbb{T} = \{m : T \rightarrow \mathbb{T} \mid m_{ij} \text{ is a polynomial of degree } \leq k\}$ for any $k \geq 0$.
- Degrees of freedom:

$$\int_F \sigma_{nt} \cdot r, \quad \text{tangential } r \text{ with } r_i \in P_k(F) \text{ on each face } F,$$
$$\int_T \sigma : \varsigma, \quad \varsigma \in P_{k-1}(T) \otimes \mathbb{T}, \text{ if } k \geq 1.$$

Theorem: These degrees of freedom are *unisolvant* for $P_k(T) \otimes \mathbb{T}$.
They generate global degrees of freedom suitable for enforcing
nt-continuity.

[G+Lederer+Schöberl 2019]

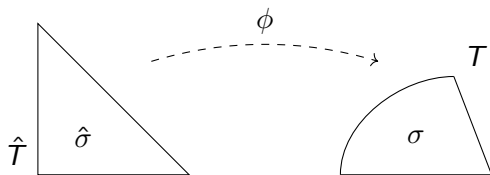


The $k = 0$ case

Mapping

- Let T = image of reference element $\hat{T}$ under (possibly curvilinear) map ϕ ,
- and define the **curl div pullback** by

$$\hat{\sigma} = \det(\nabla \phi) (\nabla \phi)^t (\sigma \circ \phi) (\nabla \phi)^{-t}.$$



Theorem: Let $F = \phi(\hat{F})$ for a facet $\hat{F}$ of $\hat{T}$ with normal $\hat{n}$ and $\hat{t} \in \hat{n}^\perp$. Let n and t be the mapped normal and tangent on F . Then

- There is a nonzero c_F depending only on $\phi|_{\hat{F}}$ such that $\hat{t} \cdot \hat{\sigma} \hat{n} = c_F (t \cdot \sigma n) \circ \phi$,
- The mapped $\hat{\sigma}$ is traceless if and only if σ is.

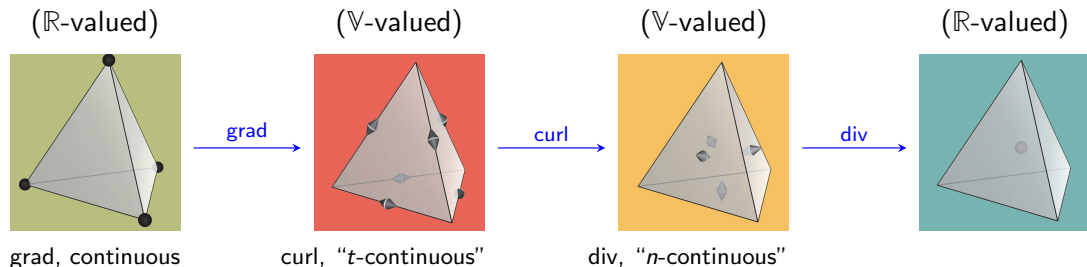
Other MCS variables have standard pullbacks (e.g., standard Piola for $H(\text{div})$ -velocity).

Thus all elements used in the MCS method have natural mappings.

Outline

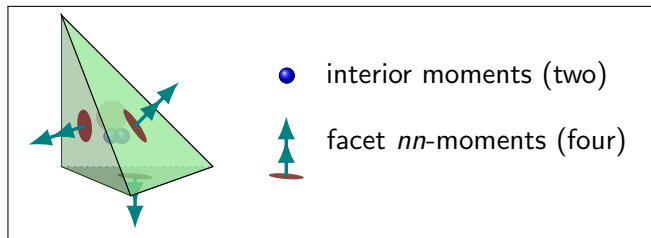
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements**
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram

Where do matrix-valued finite elements fit?



- Where does the nt -continuous $\mathbb{T}$ -valued matrix element fit?
- How about matrix elements with tt -continuity?
- How about matrix elements with nn -continuity?

More matrix-valued elements



Coupling dofs:

$$\sigma \mapsto \int_F n \cdot \sigma n$$

$\mathbb{S}$ -valued **HHJ element**, aka TDNNS element

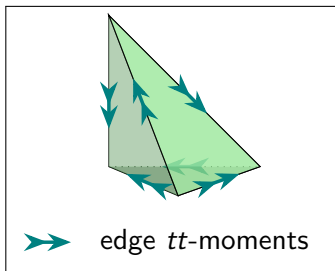
Natural operator: **div div**

Continuous **nn -component**

[[Hellan 1967](#)] [[Herrmann 1967](#)] [[Johnson 1973](#)]
[[Comodi 1989](#)] [[Pechstein+Schöberl 2011](#)]

Notation: ① $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ② $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.

More matrix-valued elements



Coupling dofs on each edge:

$$g \mapsto \int_E t \cdot gt$$

$\mathbb{S}$ -valued [Regge element](#)

Natural operator: $\text{curl}^\top \text{curl}$ or “inc”

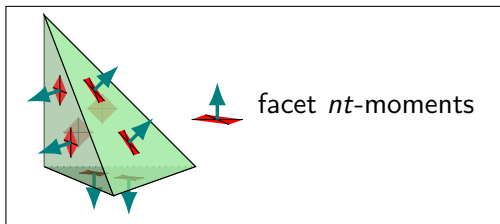
Continuous tt -component

[[Regge 1961](#)] [[Christiansen 2011](#)] [[Li 2018](#)]

[[G+Neunteufel+Schöberl+Wardetsky 2023](#)]

Notation: ❶ $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ❷ $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.

More matrix-valued elements



Two coupling dofs on each face:

$$\tau \mapsto \int_F t \cdot \tau n$$

$\mathbb{T}$ -valued viscous stress element

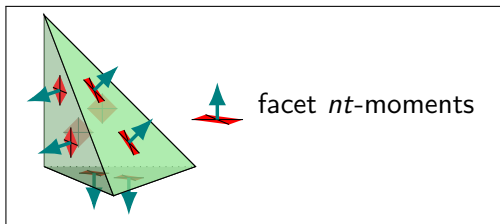
Natural operator: curl div

Continuous nt -component

[G+Lederer+Schöberl 2019]

Notation: ❶ $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ❷ $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.
❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} =$ space of traceless matrices.

More matrix-valued elements



Two coupling dofs on each face:

$$\tau \mapsto \int_F t \cdot \tau n$$

$\mathbb{T}$ -valued viscous stress element

Natural operator: curl div

Continuous nt -component

[G+Lederer+Schöberl 2019]

Recent understanding: An algebraic structure with matrix-valued Sobolev space functions connected by curl div , $\text{curl } \mathbb{T} \text{ curl}$, div div .

Notation: ❶ $\text{sym}(\tau) = \frac{1}{2}(\tau + \tau^\top)$ for matrices $\tau \in \mathbb{R}^{3 \times 3}$. ❷ $\mathbb{S} = \text{sym}(\mathbb{R}^{3 \times 3})$, space of symmetric matrices.
❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} =$ space of traceless matrices.

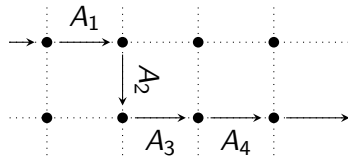
Outline

- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces**
- 6 Regular decompositions
- 7 A complementary diagram

Complexes from the 2-complex

We use “commutative diagrams”

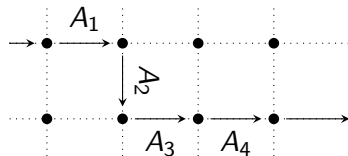
- whose “vertices” or “**objects**” are function spaces,
- whose “arrows” or “**morphisms**” are differential operators.



Complexes from the 2-complex

We use “commutative diagrams”

- whose “vertices” or “**objects**” are function spaces,
- whose “arrows” or “**morphisms**” are differential operators.



- Compositions of morphisms are called “**paths**.”
- A path is a **complex** if compositions of 2 successive morphisms vanish:

$$A_{j+1} \circ A_j = 0.$$

- A complex is **exact** if $\text{range}(A_j) = \ker(A_{j+1})$.
- A path is a **2-complex** if the compositions of 3 successive morphisms vanish: [Olver 1982]

$$A_{j+2} \circ A_{j+1} \circ A_j = 0.$$

Scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued fields in 3D

On a bounded contractible domain $\Omega \subset \mathbb{R}^3$ with Lipschitz boundary, let $\mathcal{D}$ denote the space of infinitely differentiable functions compactly supported on Ω . The **de Rham complex**

$$0 \longrightarrow \mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D} \longrightarrow \mathbb{R}$$

is exact.

Notation: $\mathbb{V} = \mathbb{R}^3$. Tensor product $X \otimes \mathbb{V}$ is identifiable with the Cartesian product $X \times X \times X$.

Scalar ($\mathbb{R}$) and vector ($\mathbb{V}$) valued fields in 3D

On a bounded contractible domain $\Omega \subset \mathbb{R}^3$ with Lipschitz boundary, let $\mathcal{D}$ denote the space of infinitely differentiable functions compactly supported on Ω . The **de Rham complex**

$$0 \longrightarrow \mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D} \longrightarrow \mathbb{R}$$

is exact, and so is the Sobolev de Rham complex,

$$0 \longrightarrow \dot{H}^1 \xrightarrow{\text{grad}} \dot{H}(\text{curl}) \xrightarrow{\text{curl}} \dot{H}(\text{div}) \xrightarrow{\text{div}} L_2 \longrightarrow \mathbb{R}$$

where

$$\begin{aligned} \dot{H}(\text{curl}) &= \overline{\mathcal{D} \otimes \mathbb{V}}^{\|\cdot\|_{H(\text{curl})}}, & \dot{H}(\text{div}) &= \overline{\mathcal{D} \otimes \mathbb{V}}^{\|\cdot\|_{H(\text{div})}} \\ \|u\|_{\dot{H}(\text{curl})}^2 &= \|u\|_{L_2}^2 + \|\text{curl } u\|_{L_2}^2, & \|q\|_{\dot{H}(\text{div})}^2 &= \|q\|_{L_2}^2 + \|\text{div } q\|_{L_2}^2. \end{aligned}$$

Notation: $\mathbb{V} = \mathbb{R}^3$. Tensor product $X \otimes \mathbb{V}$ is identifiable with the Cartesian product $X \times X \times X$.

$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

$$\mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D}$$

$$\mathcal{D} \otimes \mathbb{V}$$

$$\mathcal{D} \otimes \mathbb{S}$$

$$\mathcal{D} \otimes \mathbb{T}$$

$$\mathcal{D} \otimes \mathbb{V}$$

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of **symmetric** matrices.
❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3}$ = space of **traceless** matrices.

$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

$$\mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D}$$

$$\mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{def}} \mathcal{D} \otimes \mathbb{S} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{T} \xrightarrow{\text{div}} \mathcal{D} \otimes \mathbb{V}$$

- $\text{def } v \equiv \varepsilon(v) = \text{sym grad } v$, $\text{grad } v = \text{Jacobian matrix of vector field } v$,

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of symmetric matrices.
❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} = \text{space of traceless matrices}.$

$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

$$\mathcal{D} \xrightarrow{\text{grad}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{div}} \mathcal{D}$$

$$\mathcal{D} \otimes \mathbb{V} \xrightarrow{\text{def}} \mathcal{D} \otimes \mathbb{S} \xrightarrow{\text{curl}} \mathcal{D} \otimes \mathbb{T} \xrightarrow{\text{div}} \mathcal{D} \otimes \mathbb{V}$$

- $\text{def } v \equiv \varepsilon(v) = \text{sym grad } v$, $\text{grad } v = \text{Jacobian matrix of vector field } v$,
- curl and div on matrix-valued functions act row-wise.
- **curl** of a symmetric matrix is traceless.

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of symmetric matrices.
❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} = \text{space of traceless matrices}$.

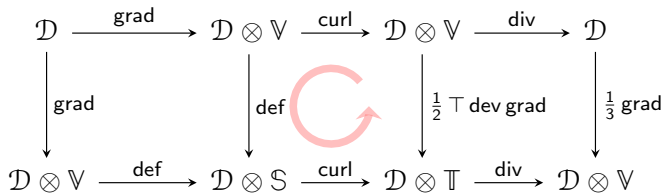
$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

$$\begin{array}{ccccccc}
 \mathcal{D} & \xrightarrow{\text{grad}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{div}} & \mathcal{D} \\
 \downarrow \text{grad} & & \downarrow \text{def} & & \downarrow \frac{1}{2} \mathbb{T} \text{ dev grad} & & \downarrow \frac{1}{3} \text{ grad} \\
 \mathcal{D} \otimes \mathbb{V} & \xrightarrow{\text{def}} & \mathcal{D} \otimes \mathbb{S} & \xrightarrow{\text{curl}} & \mathcal{D} \otimes \mathbb{T} & \xrightarrow{\text{div}} & \mathcal{D} \otimes \mathbb{V}
 \end{array}$$

- Connect vertically.

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of symmetric matrices.
 ❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3} =$ space of traceless matrices.

$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

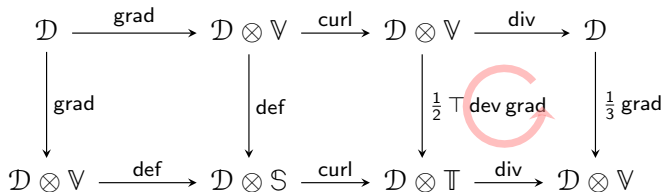


- The diagram commutes.

$$\text{curl def} = \frac{1}{2} \mathbb{T} \text{ dev grad curl}$$

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of symmetric matrices.
 ❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3}$ = space of traceless matrices.

$\mathbb{S}$ and $\mathbb{T}$ valued matrix fields

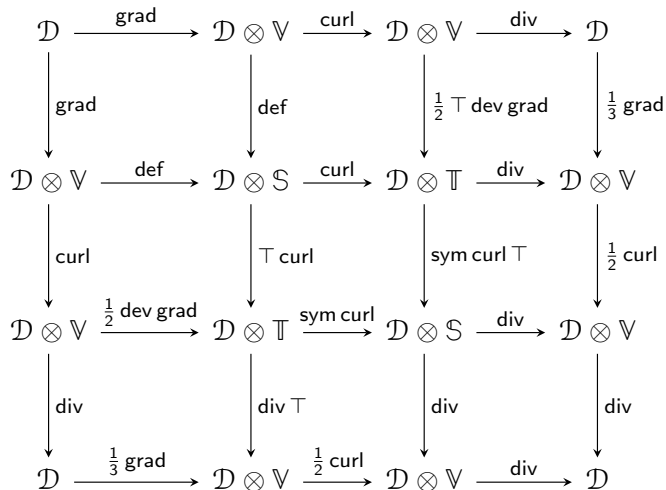


- The diagram commutes.

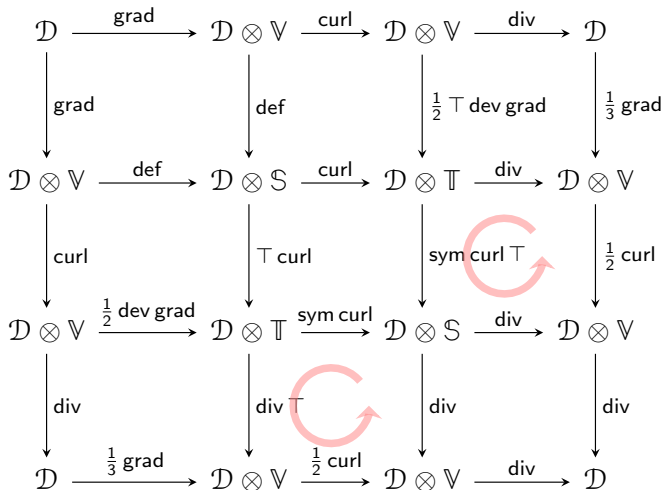
$$\begin{aligned}
 \text{curl def} &= \frac{1}{2} \mathbb{T} \text{ dev grad curl} \\
 \frac{1}{2} \text{div } \mathbb{T} \text{ dev grad } u &= \frac{1}{3} \text{grad div}
 \end{aligned}$$

Notation: ❶ $\text{sym } \tau = \frac{1}{2}(\tau + \tau^\top)$ for matrices τ . ❷ $\mathbb{S} = \text{sym } \mathbb{R}^{3 \times 3}$, space of symmetric matrices.
 ❸ $\text{dev } \tau = \tau - \frac{1}{3}(\text{tr } \tau)I$ for matrices τ , ❹ $\mathbb{T} = \text{dev } \mathbb{R}^{3 \times 3}$ = space of traceless matrices.

Completing the picture

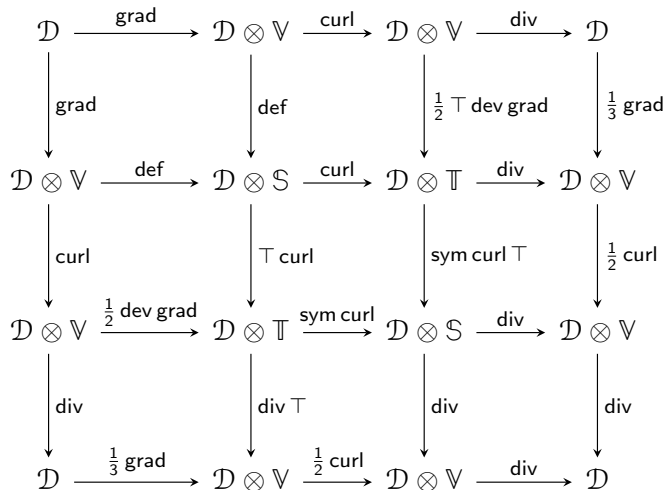


Completing the picture



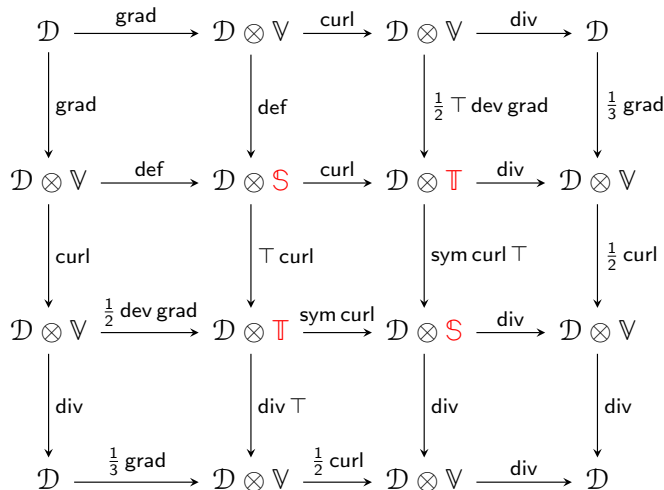
$$\text{div sym curl } \mathbb{T} = \frac{1}{2} \text{ curl div .}$$

Completing the picture



Note the symmetry in the diagram about the diagonal.

Completing the picture



The four middle spaces extend to new Sobolev spaces of matrix fields.

New and old Sobolev spaces

Let

$$H^{-1} = (\dot{H}^1)^*,$$

$$\dot{H}^{-1} = (H^1)^*.$$

Define new Sobolev spaces of matrix fields:

$$\tilde{H}_{cc} := \{g \in \tilde{H}^{-1} \otimes \mathbb{S} : \operatorname{curl} g \in \tilde{H}^{-1} \otimes \mathbb{T}, \text{curl } \top \operatorname{curl} g \in \tilde{H}^{-1} \otimes \mathbb{S}\}.$$

$$\tilde{H}_{cd} := \{\tau \in \tilde{H}^{-1} \otimes \mathbb{T} : \operatorname{div} \tau \in \tilde{H}^{-1} \otimes \mathbb{V}, \operatorname{sym} \operatorname{curl} \tau \in \tilde{H}^{-1} \otimes \mathbb{S}, \text{curl } \operatorname{div} \tau \in \tilde{H}^{-1} \otimes \mathbb{V}\}.$$

$$\tilde{H}_{cd\top} := \top \tilde{H}_{cd}.$$

$$\tilde{H}_{dd} := \{\sigma \in \tilde{H}^{-1} \otimes \mathbb{S} : \operatorname{div} \sigma \in \tilde{H}^{-1} \otimes \mathbb{V}, \operatorname{div} \operatorname{div} \sigma \in \tilde{H}^{-1}\}.$$

Define $H_{cc}, H_{cd}, H_{cd\top}, H_{dd}$ similarly, replacing $\tilde{H}^{-1}$ by H^{-1} .

Two de Rham complexes

Recall the standard Sobolev de Rham complex:

$$0 \longrightarrow \dot{H}^1 \xrightarrow{\text{grad}} \dot{H}(\text{curl}) \xrightarrow{\text{curl}} \dot{H}(\text{div}) \xrightarrow{\text{div}} L_{2,\mathbb{R}}.$$

A version with lower regularity:

$$L_{2,\mathbb{R}} \xrightarrow{\text{grad}} \hat{H}_{\mathcal{RT}}^{-1}(\text{curl}) \xrightarrow{\text{curl}} \hat{H}_{\mathcal{ND}}^{-1}(\text{div}) \xrightarrow{\text{div}} \hat{H}_{\mathcal{P}_1}^{-1}.$$

Notation:

$$L_{2,\mathbb{R}} = \{u \in L_2 : \int_{\Omega} u = 0\},$$

$$\hat{H}_{\mathcal{RT}}^{-1}(\text{curl}) = \{v \in \hat{H}^{-1} : \text{curl } v \in \hat{H}^{-1}, v(r) = 0 \text{ for all } r \in \mathcal{RT}\},$$

$$\hat{H}_{\mathcal{ND}}^{-1}(\text{div}) = \{q \in \hat{H}^{-1} : \text{div } q \in \hat{H}^{-1}, q(r) = 0 \text{ for all } r \in \mathcal{ND}\},$$

$$\hat{H}_{\mathcal{P}_1}^{-1} = \{w \in \hat{H}^{-1} : w(p) = 0 \text{ for all } p \in \mathcal{P}_1\},$$

$$\mathcal{RT} = \{a + bx : a \in \mathbb{V}, b \in \mathbb{R}\},$$

$$\mathcal{ND} = \{a + d \times x : a, d \in \mathbb{V}\},$$

$$\mathcal{P}_k = \text{polynomials of degree } \leq k.$$

Two de Rham complexes

Recall the standard Sobolev de Rham complex:

$$0 \longrightarrow \dot{H}^1 \xrightarrow{\text{grad}} \dot{H}(\text{curl}) \xrightarrow{\text{curl}} \dot{H}(\text{div}) \xrightarrow{\text{div}} L_{2,\mathbb{R}}.$$

A version with lower regularity:

$$L_{2,\mathbb{R}} \xrightarrow{\text{grad}} \hat{H}_{\mathcal{RT}}^{-1}(\text{curl}) \xrightarrow{\text{curl}} \hat{H}_{\mathcal{ND}}^{-1}(\text{div}) \xrightarrow{\text{div}} \hat{H}_{\mathcal{P}_1}^{-1}.$$

$$(Why\ remove\ \mathcal{RT}? \quad u \in L_{2,\mathbb{R}} \implies (\text{grad } u)(a + bx) = -u(\text{div}(a + bx)) = -b \int_{\Omega} u = 0.)$$

Notation:

$$L_{2,\mathbb{R}} = \{u \in L_2 : \int_{\Omega} u = 0\},$$

$$\hat{H}_{\mathcal{RT}}^{-1}(\text{curl}) = \{v \in \hat{H}^{-1} : \text{curl } v \in \hat{H}^{-1}, v(r) = 0 \text{ for all } r \in \mathcal{RT}\},$$

$$\hat{H}_{\mathcal{ND}}^{-1}(\text{div}) = \{q \in \hat{H}^{-1} : \text{div } q \in \hat{H}^{-1}, q(r) = 0 \text{ for all } r \in \mathcal{ND}\},$$

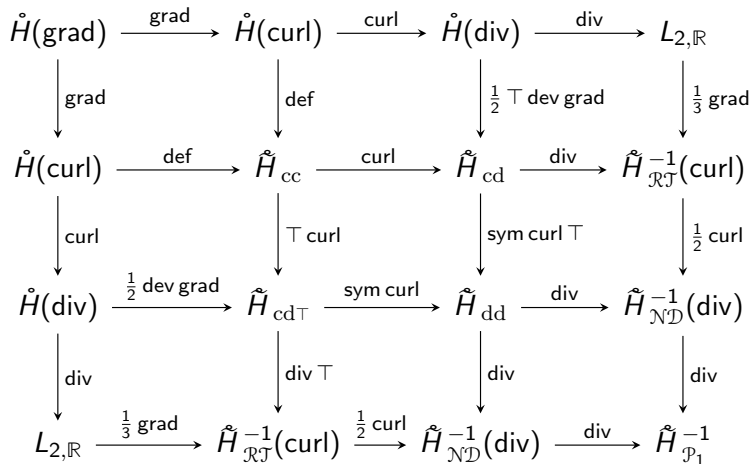
$$\hat{H}_{\mathcal{P}_1}^{-1} = \{w \in \hat{H}^{-1} : w(p) = 0 \text{ for all } p \in \mathcal{P}_1\},$$

$$\mathcal{RT} = \{a + bx : a \in \mathbb{V}, b \in \mathbb{R}\},$$

$$\mathcal{ND} = \{a + d \times x : a, d \in \mathbb{V}\},$$

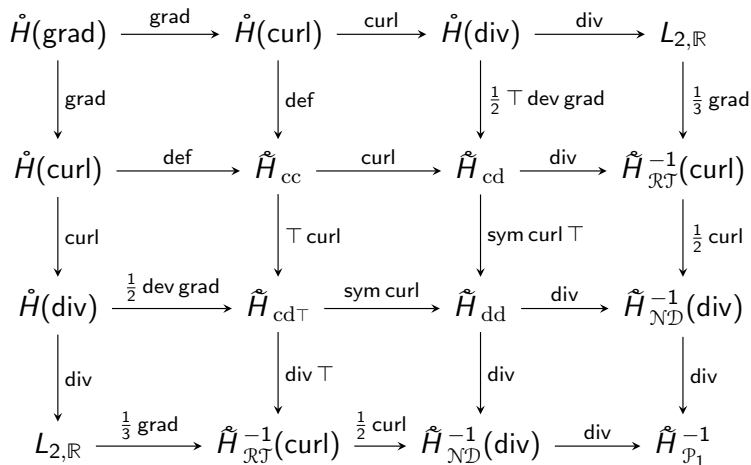
$$\mathcal{P}_k = \text{polynomials of degree } \leq k.$$

The diagram with Sobolev spaces



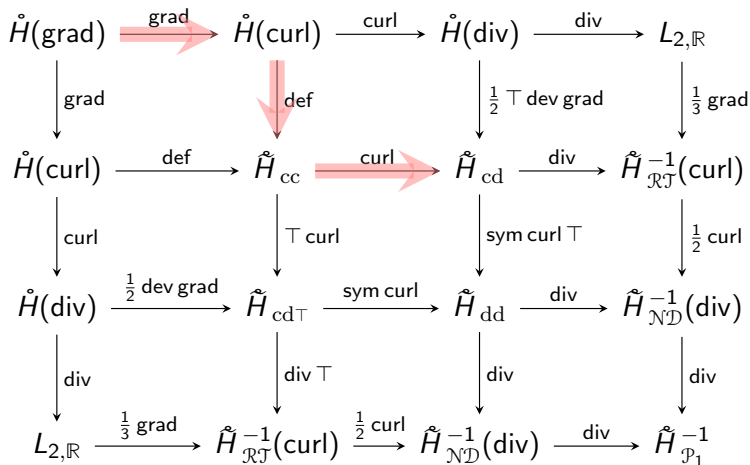
Theorem: The diagram **commutes** and every differential operator in it is **continuous** (with respect to the norms of the indicated domains and codomains) and has **closed range**.

2-complexes



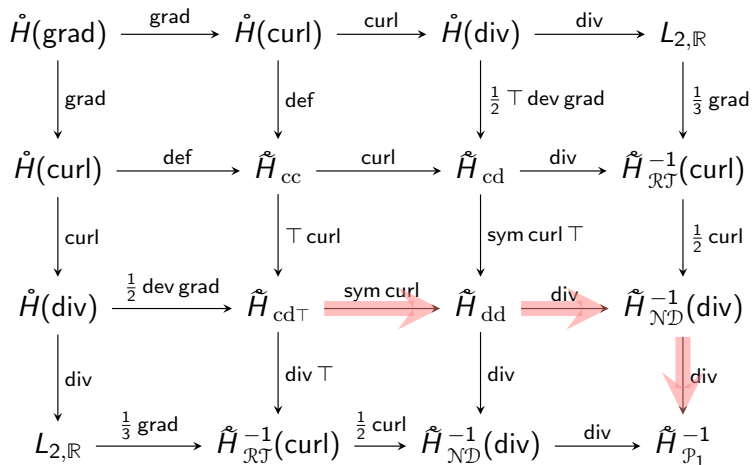
Theorem: All paths in the diagram are 2-complexes.

2-complexes



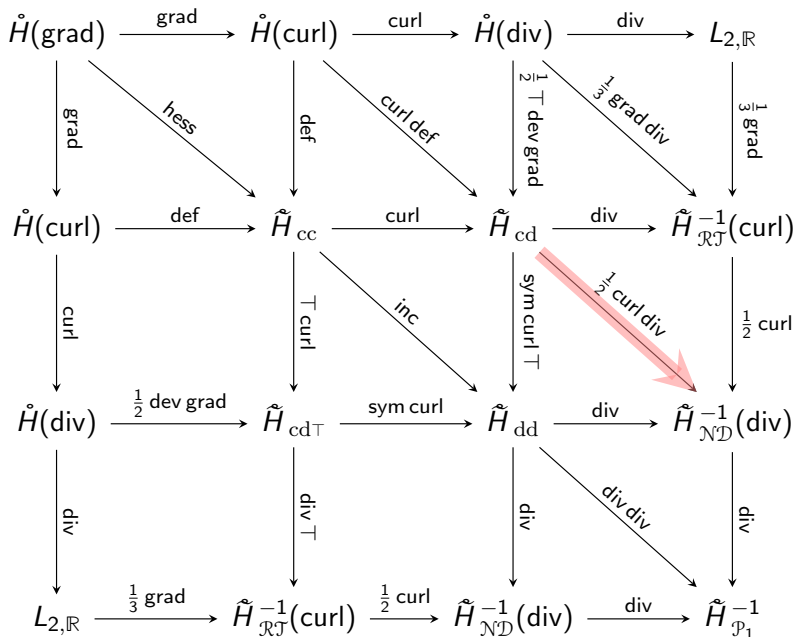
E.g., $\text{curl} \circ \text{def} \circ \text{grad } w = \frac{1}{2} \top \text{ dev grad curl grad } w = 0$.

2-complexes



E.g., $\text{div} \circ \text{div} \circ \text{sym curl } \tau = \frac{1}{2} \text{div curl div } \top \tau = 0$.

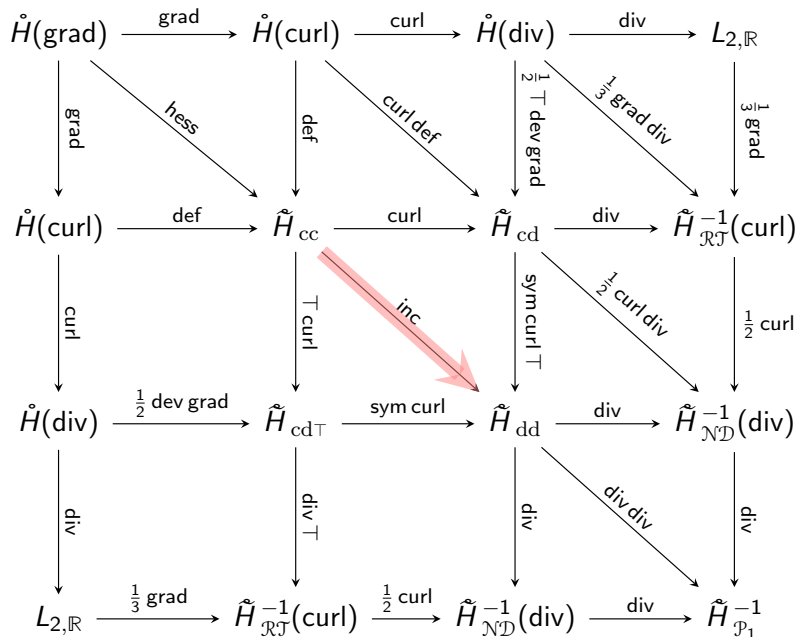
Diagonal second-order operators



Operators that appeared

- in MCS:
curl div

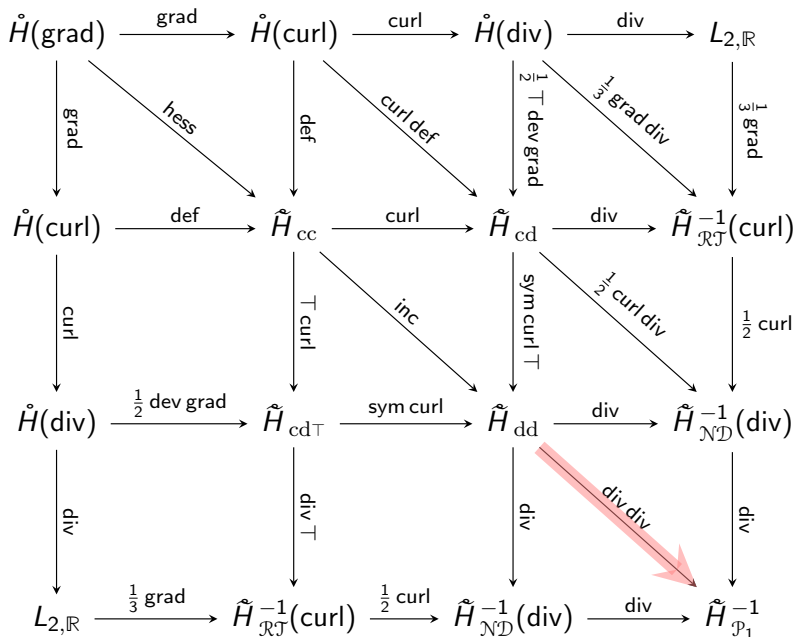
Diagonal second-order operators



Operators that appeared

- in MCS:
curl div
- in curvature:
inc = curl $\top$ curl

Diagonal second-order operators



Operators that appeared

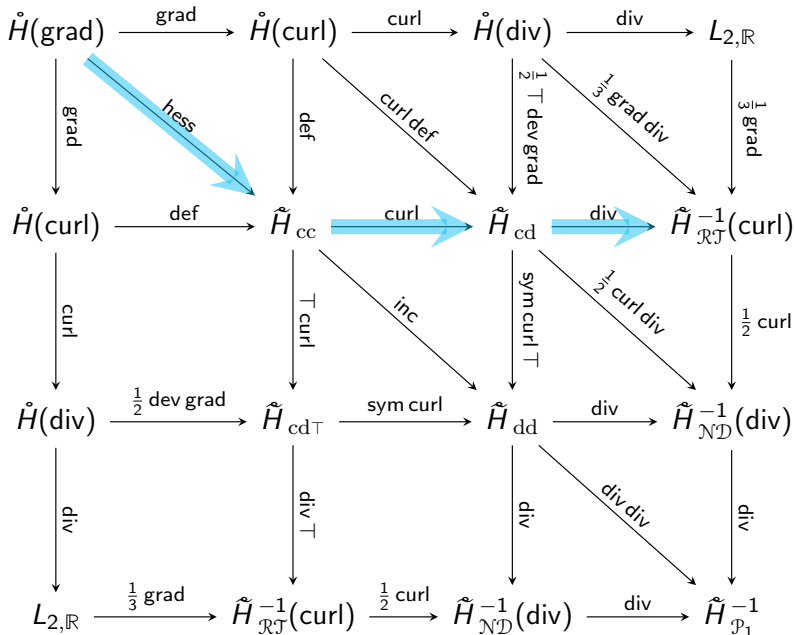
- in MCS:
curl div
- in curvature:
inc = curl $\top$ curl
- in TDNNS:
div div

Complexes

Theorem:

The following paths are **exact complexes**:

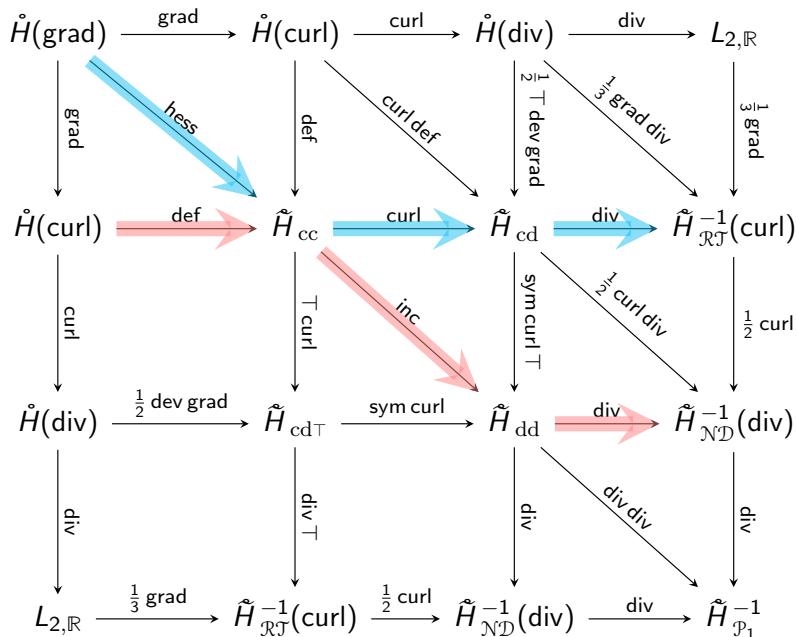
The hessian complex



Theorem:

The hessian complex

The elasticity complex



Complexes

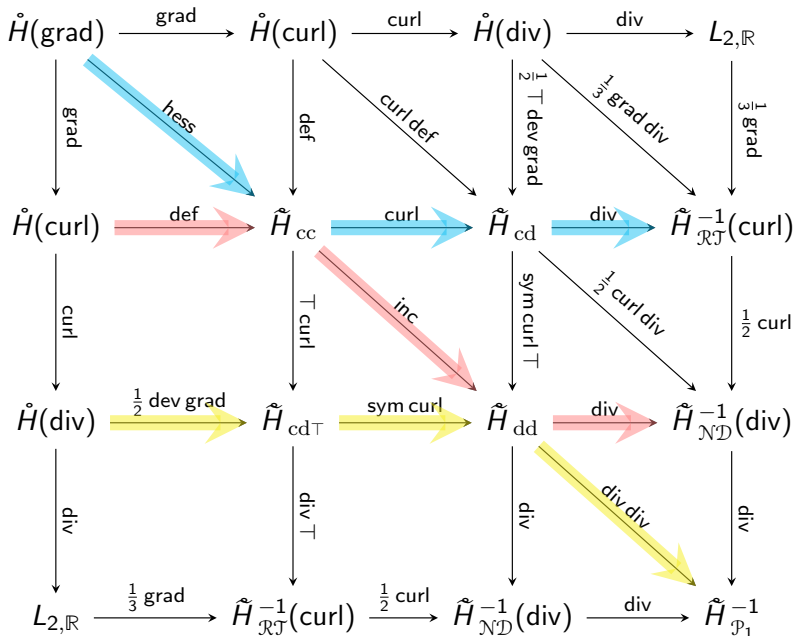
Theorem:

The following paths are **exact complexes**:

The hessian complex

The elasticity complex

The div div complex



Outline

- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions**
- 7 A complementary diagram

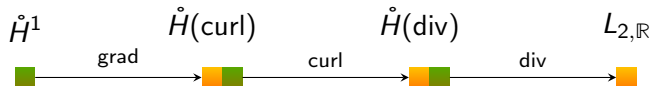
Generalized Bogovskiĭ operators

- A classical regular right inverse of $\text{div}(\cdot)$ was developed by [Bogovskiĭ 1979]. It was generalized to r -forms in n -dimensions by [Costabel+McIntosh 2010]. In 3D,

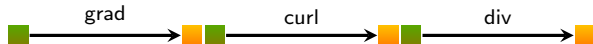
$$\dot{H}^1 \xrightleftharpoons[T_g]{\text{grad}} \dot{H}(\text{curl}) \xrightleftharpoons[T_c]{\text{curl}} \dot{H}(\text{div}) \xrightleftharpoons[T_d]{\text{div}} L_{2,\mathbb{R}},$$

where T_g , T_c , and T_d map into **more regular** subspaces with $\dot{H}^1$ components.

- Consequently, the spaces can be split into **smooth** and **rough** parts:



Regular decompositions of the spaces



Green square $\dot{H}^1$ (the good)
Orange square L_2 (the bad)

The standard case

There are continuous linear maps

$$S_c : \dot{H}(\text{curl}) \rightarrow \dot{H}^1 \text{ and}$$

$$S_d : \dot{H}(\text{div}) \rightarrow \dot{H}^1 \otimes \mathbb{V} \text{ such that}$$

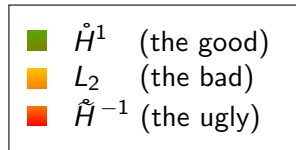
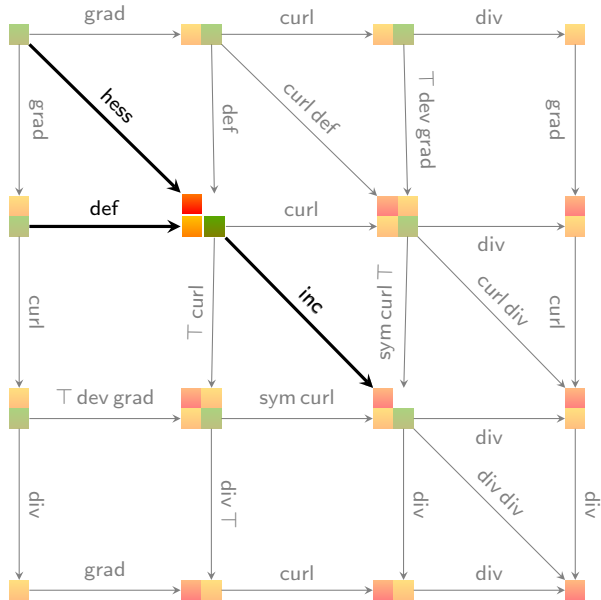
$$\forall u \in \dot{H}(\text{curl}) :$$

$$u = \text{grad}(S_c u) + T_c \text{curl } u,$$

$$\forall q \in \dot{H}(\text{div}) :$$

$$q = \text{curl}(S_d q) + T_d \text{div } q.$$

Regular decompositions of the spaces

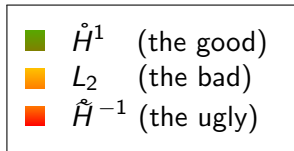
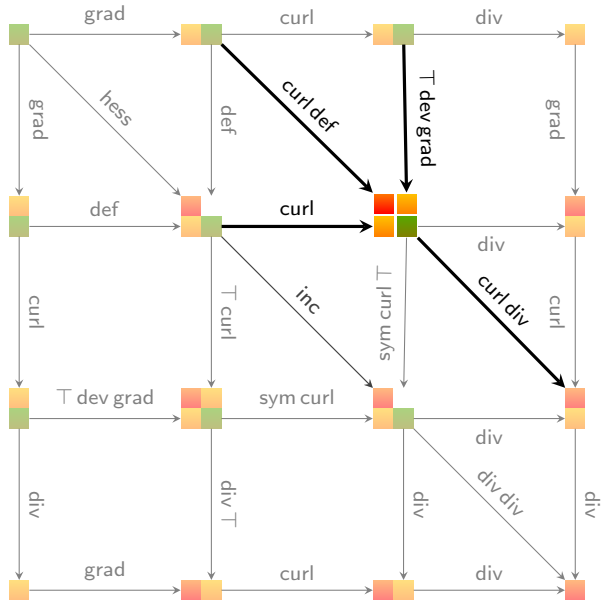


The $\hat{H}_{cc}$ case

Theorem: Any g in $\hat{H}_{cc}$ admits the decomposition

$$g = \text{hess}(S_{cc}^{(1)} g) + \text{def}(S_{cc}^{(2)} g) + D_{cc} \text{inc } g.$$

Regular decompositions of the spaces

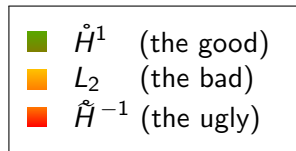
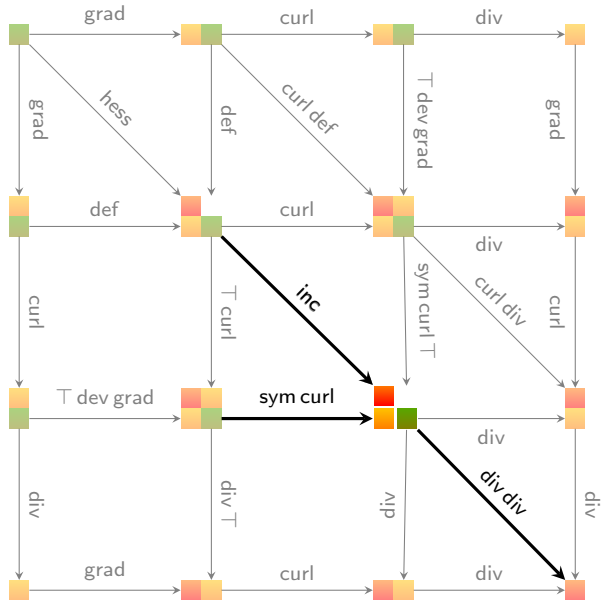


The $\hat{H}_{cd}$ case

Theorem: Any τ in $\hat{H}_{cd}$ admits the decomposition

$$\begin{aligned} \tau = & \text{curl def}(S_{cd}^{(1)} \tau) \\ & + \text{curl}(S_{cd}^{(2)} \tau) \\ & + \text{T dev grad}(S_{cd}^{(3)} \tau) \\ & + D_{cd} \text{curl div } \tau. \end{aligned}$$

Regular decompositions of the spaces



The $\hat{H}_{dd}$ case

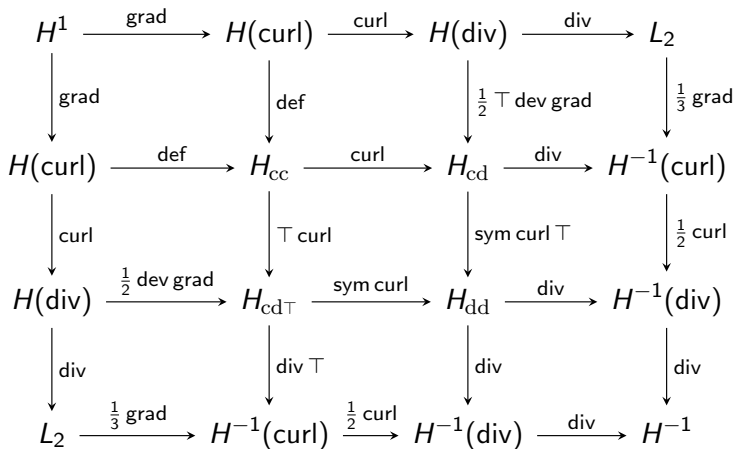
Theorem: Any σ in $\hat{H}_{dd}$ admits the decomposition

$$\begin{aligned} \sigma = & \text{inc}(S_{dd}^{(1)} \sigma) \\ & + \text{sym curl}(S_{dd}^{(2)} \sigma) \\ & + D_{dd} \text{div div } \sigma. \end{aligned}$$

Outline

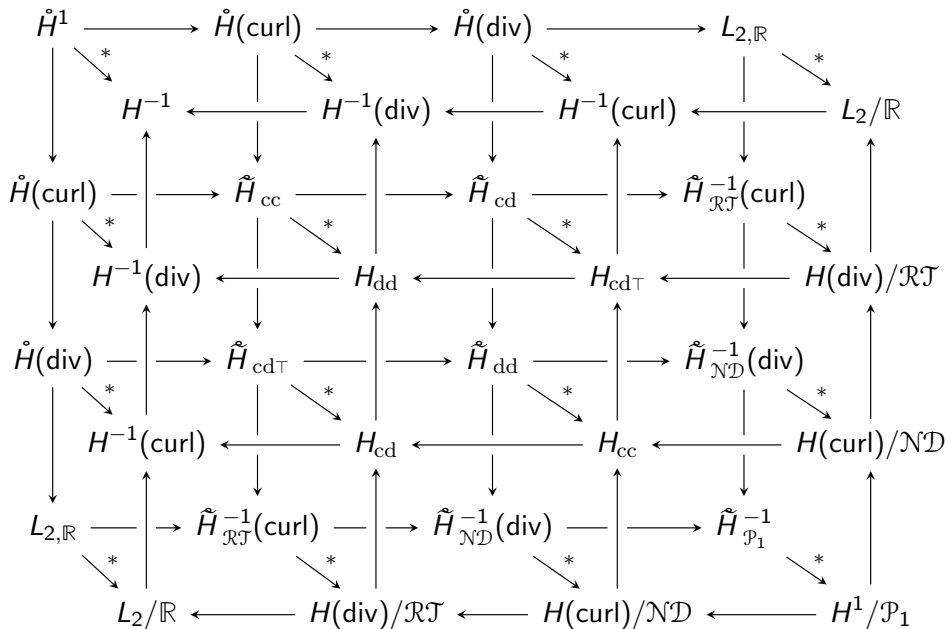
- 1 Introduction
- 2 Mass-Conserving Stress-yielding (MCS) method
- 3 Viscous stress elements
- 4 Other matrix finite elements
- 5 A unifying 2-complex of Sobolev spaces
- 6 Regular decompositions
- 7 A complementary diagram**

Diagram for the case **without boundary conditions**



Prior results extend to this diagram also. (We can also start with $H^1/\mathcal{P}_1$ instead of H^1 and mod out appropriate subspaces as in the next slide.)

Duality



Duality identities

$$① \quad (\mathring{H}^1)^* = H^{-1}$$

$$② \quad (H^1)^* = \mathring{H}^{-1}$$

$$③ \quad \mathring{H}(\operatorname{div})^* = H^{-1}(\operatorname{curl})$$

$$④ \quad \mathring{H}(\operatorname{curl})^* = H^{-1}(\operatorname{div})$$

$$⑤ \quad H(\operatorname{div})^* = \mathring{H}^{-1}(\operatorname{curl})$$

$$⑥ \quad H(\operatorname{curl})^* = \mathring{H}^{-1}(\operatorname{div})$$

$$⑦ \quad \mathring{H}_{\operatorname{cc}}^* = H_{\operatorname{dd}}$$

$$⑧ \quad \mathring{H}_{\operatorname{cd}}^* = H_{\operatorname{cd}\top}$$

$$⑨ \quad \mathring{H}_{\operatorname{dd}}^* = H_{\operatorname{cc}}$$

Conclusion

- Operator curl div is natural for **viscous stresses** of Stokes flow with $H(\text{div})$ velocity.
- Simple ***nt*-continuous** matrix finite elements can be designed: [G+Lederer+Schöberl 2019].
- The **MCS method** yield optimal rates for fluid stresses and is Reynolds robust.
- MCS convergence analysis: [G+Lederer+Schöberl 2020], [G+Kogler+Lederer+Schöberl 2023].
- Other *nn*- and *tt*-continuous elements motivate study of accompanying Sobolev spaces.
- They feature 2nd order operators **curl div**, **div div**, and **inc** on matrix fields.
- Two commuting diagrams that give **2-complexes** were seen. One is **dual** to the other.
- They contain new Sobolev spaces H_{cc} , H_{cd} , H_{dd} of matrix fields with low-regularity.
- Through **regular decompositions**, we understand these spaces better.
- Recent ArXiv preprint: [G+Hu+Schöberl 2025].