The Shifted Boundary Method: A Framework for Immersed Computational Mechanics

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Acknowledgments

Collaborators on the work presented here (present and past):

Alex Main (former Postdoc, now Research Scientist at Ansys) Ting Song (former Ph.D. now Research Scientist at ExxonMobil) Oriol Colomés (former Postdoc, now Faculty at TU Delft in The Netherlands) Léo Nouveau (former Postdoc, now Faculty at U. Rennes in France) Nabil Atallah (former Ph.D. student, now Postdoc at LLNL) Kangan Li (Ph.D. student, just graduated, *looking for jobs*)



Motivation I

Complex geometry is still a key challenge in engineering simulations

- Engineering designs are often characterized by geometric complexity
- The merging of topology optimization and advanced manufacturing (e.g., additive manufacturing) may exacerbate this trend

Structural joint (welded)

Structural joint (3D-printed)







Geometric model of a vascular stent

Motivation II



Motivation III

Imaging-to-computing, an *emerging* field:

An efficient transition from geometries reconstructed from images to computation may impact and transform many fields of application:

- Biomedical engineering: CT-scans are given as pixilated data or STL format (collections of triangular facets and their nodal coordinates).
 Body-fitted meshing can be quite hard to perform.
- Subsurface imaging and computing (meshing requires considerable effort in reservoir engineering applications)
- Additive manufacturing simulations (e.g., 3D-printing). The typical file format for 3D-printers is again STL

In these examples the geometric information is not very precise and/or consistent (surfaces with *gaps* and *overlaps*, typical of computer graphics, STL = set of *disconnected* triangular faces)



Two commonly used computational strategies:

- 1. Body-fitted grids. The grid conforms to the boundary geometry of the shape to be simulated.
 - <u>Advantages</u>: Easier treatment of the boundary conditions (and boundary layers)
 - <u>Limitations</u>: Requires more advanced meshing for complex geometry, or re-meshing in problems with large deformations
- 2. Embedded/immersed grids. The shape to be simulated is fully or partially embedded (or immersed) into a regular background grid.
 - <u>Advantages</u>: Generality of the method, especially if coupling heterogeneous computational frameworks, rapid prototyping
 - Limitations: More complex enforcement of boundary conditions

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Existing Embedded Boundary Methods

Unfitted Finite Element Methods (excluding GFEM-XFEM)

- Embedded methods of finite element type (a.k.a. *cutFEMs, unfitted FEMs, Finite Cell Method, Embedded Splines, IGA-Immersogeometric etc.*) often rely on XFEM methodologies to integrate on cut cells, *Inverse Lax-Wendroff procedure* (DG) [Burman, Hansbo, Larson, Massing, Cirak, Kamenski, Schillinger, Yan, Parvizian, Düster, Rank, Wall, Annavarapu, Dolbow, Harari, Badia, Rossi, C-W. Shu, Masud, et. al., etc.]
- Unfitted/embedded FEMs typically utilize Lagrange multipliers or Nitsche variational formulations
- CutFEMs/unfitted FEMs require data structures and special quadratures to integrate on geometrically complex cut cells
- The small cut-cell problem: Integration over cut cells introduces additional interface degrees of freedom that may yield stability problems, very small time-steps or poor matrix conditioning. [Burman & Hansbo Appl. Num. Math. (2012)]. Solution: ghost penalty, and related methods

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Key ideas: [with Alex Main at Duke University: the origins of the method]

• Use a *purely* embedded approach





- Use a *purely* embedded approach
- Use the Nitsche framework to impose boundary conditions weakly





- Use a *purely* embedded approach
- Use the Nitsche framework to impose boundary conditions weakly
- Apply boundary conditions on a surrogate boundary, near the true boundary





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- Apply boundary conditions on a surrogate boundary, near the true boundary
- Appropriately modify the boundary condition to account for the discrepancy between surrogate and true boundary





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The Shifted Boundary Method (SBM)



The extension map M & a distance vector function d $M : \tilde{\Gamma} \to \Gamma$ $\tilde{x} \mapsto x$ $d_M(\tilde{x}) = x - \tilde{x} = [M - I](\tilde{x})$

Extension of functions defined on boundaries: $\bar{\psi}(\tilde{x}) \equiv \psi(M(\tilde{x}))$



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The (Base) Nitsche Method

A prototypical example: The Poisson problem

Strong form of the equation $\Delta u + f = 0$ on Ω u = g on $\Gamma = \partial \Omega$



Augmented Lagrangian formulation

$\int_{\Omega} \nabla w \cdot \nabla u$	1 1	$=\int_{\Omega} w f$
5 22		J 12

Nitsche's formulation

$$\int_{\Omega} \nabla w \cdot \nabla u - \int_{\Gamma} w \, \frac{\partial u}{\partial n} - \int_{\Gamma} \frac{\partial w}{\partial n} \, (u - g) + \int_{\Gamma} \alpha \, w \, (u - g) = \int_{\Omega} w \, f$$

The Shifted Nitsche Method

A prototypical example: The Poisson problem

Weak form of the equation with weak boundary conditions (Nitsche method):

$$\int_{\Omega} \nabla w \cdot \nabla u - \int_{\Gamma} w \frac{\partial u}{\partial n} - \int_{\Gamma} \frac{\partial w}{\partial n} (u - g) + \int_{\Gamma} \alpha w (u - g) = \int_{\Omega} w f$$
Shifted Nitsche method:

$$\int_{\overline{\Omega}} \nabla w \cdot \nabla u - \int_{\overline{\Gamma}} w \frac{\partial u}{\partial n} - \int_{\overline{\Gamma}} \frac{\partial w}{\partial n} (u - \overline{g}) + \int_{\overline{\Gamma}} \alpha w (u - \overline{g}) = \int_{\overline{\Omega}} w f$$

$$\int_{\overline{\Omega}} \frac{n}{\overline{\Gamma}} = \partial \overline{\Omega}$$
Surrogate computational domain
[with Alex Main at Duke University: the origins of the method]

The Shifted Nitsche Method

A prototypical example: The Poisson problem

Weak form of the equation with weak boundary conditions (Nitsche method):

$$\int_{\tilde{\Omega}} \nabla w \cdot \nabla u - \int_{\tilde{\Gamma}} w \frac{\partial u}{\partial n} - \int_{\tilde{\Gamma}} \frac{\partial w}{\partial n} (u - \tilde{g}) + \int_{\tilde{\Gamma}} \alpha w (u - \tilde{g}) = \int_{\tilde{\Omega}} w f$$
$$\int_{\tilde{\Omega}} \nabla w \cdot \nabla u - \int_{\tilde{\Gamma}} w \frac{\partial u}{\partial n} - \int_{\tilde{\Gamma}} \frac{\partial w}{\partial n} (u + \nabla u \cdot d - g) + \int_{\tilde{\Gamma}} \alpha w (u + \nabla u \cdot d - g) = \int_{\tilde{\Omega}} w f$$



Numerical Results: Poisson Problem

Numerical convergence test with an exact solution on a circular domain





Condition number (estimated!):

- About 3 times larger than for the body-fitted method
- Same scaling with h⁻² has for the body-fitted method

Numerical Results: Poisson Problem

[with Nabil Atallah (at Duke) & Claudio Canuto, Math. Dept., Politecnico di Torino]





Piecewise-linear triangular finite elements



Mesh Size	No. of surrogate edges with $\boldsymbol{v} \cdot \boldsymbol{n} \leq 0$	% of total surrogate edges
4.00E-02	1	4.35%
2.00E-02	1	2.33%
1.00E-02	5	5.43%
5.00E-03	9	5.06%
2.50E-03	23	6.35%
1.25E-03	38	5.38%

The SBM for Domains with Corners

[with Nabil Atallah (at Duke) & Claudio Canuto, Math. Dept., Politecnico di Torino] Solution: Consider a surrogate edge \tilde{E} with \tilde{n} as its unit normal.



The three-dimensional case is similar

N. Atallah, C. Canuto, G. Scovazzi, "Analysis of the Shifted Boundary Method for the Poisson problem in domains with corners," Mathematics of Computation

Coercivity of the SBM variational form

We have a coercivity estimate:

$$a_u^h(u^h, u^h) \ge C_a \|u^h\|_a^2, \qquad \forall u^h \in V^h(\tilde{\Omega}),$$

where

$$\|u^{h}\|_{a}^{2} = \|\nabla u^{h}\|_{0,\tilde{\Omega}}^{2} + \|\sqrt{1/h} \left(u^{h} + \nabla u^{h} \cdot \boldsymbol{d}\right)\|_{0,\tilde{\Gamma}_{D}}^{2}$$



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$$\|u^{h}\|_{a}^{2} = \|\nabla u^{h}\|_{0;\tilde{\Omega}}^{2} + \|\sqrt{1/h} \left(u^{h} + \nabla u^{h} \cdot \boldsymbol{d}\right)\|_{0;\tilde{\Gamma}_{D}}^{2}$$

Boundedness (Continuity)

We can also prove continuity (boundedness): $V(\tilde{\Omega};h) = V^{h}(\tilde{\Omega}) + V(\tilde{\Omega}) \subset H^{2}(\tilde{\Omega};\tilde{\mathscr{T}}^{h})$ $a^{h}_{u}(u,w^{h}) \leq C_{\mathscr{A}} \parallel w^{h} \parallel_{a} \parallel u \parallel_{V(\tilde{\Omega};h)}, \qquad \parallel v \parallel^{2}_{V(\tilde{\Omega};h)} = \parallel v \parallel^{2}_{a} + |hw|^{2}_{2;\tilde{\Omega};\tilde{\mathscr{T}}^{h}}$

for all $u \in V(\tilde{\Omega}; h)$ and for all $w^h \in V^h(\tilde{\Omega})$.

Consistency error due to the Taylor expansions:

$$\begin{aligned} a_{u}^{h}(u,w^{h}) - l_{u}^{h}(w^{h}) &= -(w^{h}, \underbrace{\Delta u + f}_{0})_{\tilde{\Omega}} - \langle \sqrt{h} \nabla w^{h} \cdot \tilde{\boldsymbol{n}}, \sqrt{1/h} \rangle_{\tilde{\Gamma}_{D}} \\ &+ \langle \alpha \sqrt{1/h} \left(w^{h} + \nabla w^{h} \cdot \boldsymbol{d} \right), \sqrt{1/h} \rangle_{\tilde{\Gamma}_{D}} \\ &\leq C_{PBL} ||w^{h}||_{a} h_{\tilde{\Gamma}_{D}} \hat{h}_{\tilde{\Gamma}_{D}}^{3/2} ||\nabla(\nabla u)||_{0,\Omega} . \end{aligned}$$



Consistency error due to the Taylor expansions:

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Optimal convergence estimate in the energy norm (using Strang's second lemma)

$$\begin{aligned} \| u - u^{h} \|_{V(\tilde{\Omega};h)} &\leq \left(1 + \frac{\|a_{u}^{h}\|_{V(h) \times V^{h}(\tilde{\Omega})}}{C_{a}} \right) \inf_{w^{h} \in V^{h}(\tilde{\Omega})} \| u - w^{h} \|_{V(\tilde{\Omega};h)} \\ &+ \frac{1}{C_{a}} \sup_{w^{h} \in V^{h}(\tilde{\Omega})} \frac{|l^{h}(w^{h}) - a_{u}^{h}(u, w^{h})|}{\|w^{h}\|_{V(\tilde{\Omega};h)}} \,. \end{aligned}$$

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Consistency error due to the Taylor expansions:

$$\begin{aligned} a_{u}^{h}(u,w^{h}) - l_{u}^{h}(w^{h}) &= -(w^{h}, \underbrace{\Delta u + f}_{0})_{\tilde{\Omega}} - \langle \sqrt{h} \nabla w^{h} \cdot \tilde{\boldsymbol{n}}, \sqrt{1/h} \rangle_{\tilde{\Gamma}_{D}} \\ &+ \langle \alpha \sqrt{1/h} \left(w^{h} + \nabla w^{h} \cdot \boldsymbol{d} \right), \sqrt{1/h} \rangle_{\tilde{\Gamma}_{D}} \\ &\leq C_{PBL} ||w^{h}||_{a} h_{\tilde{\Gamma}_{D}} \hat{h}_{\tilde{\Gamma}_{D}}^{3/2} ||\nabla(\nabla u)||_{0,\Omega} . \end{aligned}$$

Optimal convergence estimate in the energy norm (using Strang's second lemma)

$$\begin{split} \| u - u^{h} \|_{V(\tilde{\Omega};h)} &\leq \left(1 + \frac{\|a_{u}^{h}\|_{V(h) \times V^{h}(\tilde{\Omega})}}{C_{a}} \right) \inf_{w^{h} \in V^{h}(\tilde{\Omega})} \| u - w^{h} \|_{V(\tilde{\Omega};h)} \\ &+ \frac{1}{C_{a}} \sup_{w^{h} \in V^{h}(\tilde{\Omega})} \frac{|l^{h}(w^{h}) - a_{u}^{h}(u, w^{h})|}{\|w^{h}\|_{V(\tilde{\Omega};h)}} \,. \end{split}$$

Suboptimal (not in practice!) duality (L²) estimate: $\|u - u^h\|_{0,\tilde{\Omega}} \leq C_0 h_{\tilde{\Omega}}^{3/2} \|\nabla(\nabla u)\|_{0,\tilde{\Omega}}$

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Numerical Analysis: Stokes Flow Problem [with Nabil Atallah (at Duke) & Claudio Canuto, Math. Dept., Politecnico di Torino] Stability (LBB)

 $\mathcal{B}([\boldsymbol{u}^h, \boldsymbol{p}^h]; [\boldsymbol{w}^h, \boldsymbol{q}^h]) \geq \alpha_{LBB} \| [\boldsymbol{u}^h, \boldsymbol{p}^h] \|_{\mathcal{B}} \| [\boldsymbol{w}^h, \boldsymbol{q}^h] \|_{\mathcal{B}}$

An LBB inf-sup condition can be derived in the case of the Stokes' operator

Convergence (in natural norm)

$$\| [\boldsymbol{u}, p] - [\boldsymbol{u}^h, p^h] \|_{\boldsymbol{W}(\tilde{\Omega}; h)} \le C h_{\tilde{\Omega}} \left(\| \nabla (\nabla \boldsymbol{u}) \|_{0, \Omega} + \| \nabla p \|_{0, \tilde{\Omega}} \right)$$

The proof is analogous to the one for the Poisson problem, using the *inf-sup* LBB condition and Strang's second lemma.

Duality estimates (L²-estimates for the velocity field)

$$\|\boldsymbol{u} - \boldsymbol{u}^h\|_{0;\tilde{\Omega}} \le C_0 h_{\tilde{\Omega}}^{3/2} \left(\|\nabla(\nabla \boldsymbol{u})\|_{0,\tilde{\Omega}} + \|\nabla p\|_{0,\tilde{\Omega}} \right)$$

Analogous but more complicated proof than in the Poisson case. We observe quadratic convergence for the velocity, in practical calculations

Similar results hold for the Darcy flow equations (mixed-Poisson)

Numerical Analysis: Linear Elasticity

[with Nabil Atallah (at Duke) & Claudio Canuto, Math. Dept., Politecnico di Torino] Stability (coercivity by Korn's inequality)

 $\mathscr{B}_{\text{SBM}}^{[\tilde{\Omega}]}([\boldsymbol{u}^{h},\boldsymbol{\varepsilon}^{h}];[\boldsymbol{u}^{h},\boldsymbol{\varepsilon}^{h}]) \geq c_{\text{SBM}} \left| \left| \left| \left[\boldsymbol{u}^{h},\boldsymbol{\varepsilon}^{h} \right] \right| \right| \right|_{\text{SBM};\tilde{\Omega}}^{2}$

 $\begin{aligned} ||| [\boldsymbol{u}^{h}, \boldsymbol{\varepsilon}^{h}] |||_{_{\text{SBM}}}^{2} &= \left(1 - \tau_{\varepsilon}/2\right) \|\boldsymbol{\mathsf{C}}^{1/2} \nabla^{s} \boldsymbol{u}^{h}\|_{0;\tilde{\Omega}_{h} \setminus \tilde{\Omega}_{h}^{\ell}}^{2} + \tau_{\varepsilon} \|\boldsymbol{\mathsf{C}}^{1/2} \nabla^{s} \boldsymbol{u}^{h}\|_{0;\tilde{\Omega}_{h}^{\ell}}^{2} + \tau_{\varepsilon} \rho_{\min} \left(l(\tilde{\Omega})^{-2} \|\boldsymbol{u}^{h}\|_{H^{1}(\tilde{\Omega}_{h})}^{2} + \|h^{1/2} \nabla \boldsymbol{u}^{h}\|_{0;\tilde{\Gamma}_{D,h}}^{2}\right) \\ &+ \left(1 - \tau_{\varepsilon}\right) \left(\|\boldsymbol{\mathsf{C}}^{1/2} \,\boldsymbol{\varepsilon}^{h}\|_{0;\tilde{\Omega}_{h}^{\ell}}^{2} + \|h^{1/2} \,(\boldsymbol{\mathsf{C}}^{1/2} \,\boldsymbol{\varepsilon}^{h})\|_{0;\tilde{\Gamma}_{D,h}}^{2}\right) + \|\rho_{\mathsf{c}}^{-1/2} \,\tau_{\boldsymbol{u}}^{1/2} \,\nabla \cdot (\boldsymbol{\mathsf{C}} \,\boldsymbol{\varepsilon}^{h})\|_{0;\tilde{\Omega}_{h}^{\ell}}^{2} + \|\rho_{\mathsf{c}}^{1/2} \,\boldsymbol{h}^{-1/2} \,\boldsymbol{S}_{D,h} \,\boldsymbol{u}^{h}\|_{0;\tilde{\Gamma}_{D,h}}^{2} \end{aligned}$

Convergence (energy norm)

$$|||[\boldsymbol{e}_{\boldsymbol{u}},\boldsymbol{e}_{\boldsymbol{\varepsilon}}]|||_{\text{SBM}} \lesssim \rho_{\max}^{1/2} \left(h_{\tilde{\Omega}_{h}} |\boldsymbol{u}|_{2;\tilde{\Omega}_{h}} + h_{\tilde{\Gamma}_{D,h}} |\boldsymbol{u}|_{2;\Omega \setminus \tilde{\Omega}_{h}} + h_{\tilde{\Gamma}_{D,h}}^{3/2} l(\tilde{\Omega}_{h})^{1/2} |\boldsymbol{\varepsilon}|_{2;\tilde{\Omega}_{h}^{\ell}} \right)$$

Optimal convergence is achieved.

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Convergence (L2-estimates for the displacement field)

 $\|\boldsymbol{e}_{\boldsymbol{u}}\|_{0;\tilde{\Omega}_{h}} \lesssim \rho_{\max}^{1/2} \rho_{\min}^{-1/2} h_{\tilde{\Omega}_{h}}^{3/2} l(\tilde{\Omega}_{h})^{1/2} (|\boldsymbol{u}|_{2;\Omega} + |\boldsymbol{\varepsilon}|_{2;\tilde{\Omega}_{h}^{\ell}})$

We observe quadratic convergence for the displacement, in practical calculations

Porous Media Flow (Mixed Formulation)



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 $\boldsymbol{\beta} \cdot \boldsymbol{n} = h_N \text{ on } \Gamma_N$

Porous Media Flow (Mixed Formulation)

~20 million elements setup time < 1 minute Surrogate geometry $(\boldsymbol{\Lambda}^{-1}\boldsymbol{\beta} + \boldsymbol{\nabla} p = 0 \quad \text{in } \Omega$ $\nabla \cdot \boldsymbol{\beta} = \phi \quad \text{in } \Omega$ [with Nabil Atallah and in collaboration with G. Katul, CEE, Duke University] $p = p_D$ on Γ_D

 $\boldsymbol{\beta} \cdot \boldsymbol{n} = h_N \text{ on } \Gamma_N$

Porous Media Flow (Mixed Formulation)



Solution (flux)

[with Nabil Atallah and in collaboration with G. Katul, CEE, Duke University]

 $\begin{pmatrix} \mathbf{\Lambda}^{-1} \boldsymbol{\beta} + \boldsymbol{\nabla} p = 0 & \text{in } \Omega \\ \boldsymbol{\nabla} \cdot \boldsymbol{\beta} = \phi & \text{in } \Omega \\ p = p_D & \text{on } \Gamma_D \\ \boldsymbol{\beta} \cdot \boldsymbol{n} = h_N & \text{on } \Gamma_N \end{cases}$

The Shifted Boundary Method

Convergence test (Navier-Stokes): Decelerating cylinder test



	With Extrapolation	Without
	With Extrapolation	Extrap.
Mesh Size	2 nd order	1 st order
0.05	0.1338	2.548
0.025	4.143 * 10^-2	0.9565
0.0125	1.067 * 10^-2	0.4761
0.00625	3.422 * 10^-3	0.265
0.003125	9.576 *10^-4	0.1404
К	1.78	1

Exact solution:

$$u_{\phi}(\mathbf{r}, \mathbf{t}) = \frac{C}{J_{1}\left(\sqrt{\frac{\mathrm{i}w\rho}{\mu}R}\right)} J_{1}\left(\sqrt{\frac{\mathrm{i}w\rho}{\mu}r}\right) e^{(-iwt)}$$



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Turbulent Flows

Formulations based on turbulent viscosities:

Find $\mathbf{u} \in V^h(\tilde{\Omega})$ and $p \in Q^h(\tilde{\Omega})$ such that, $\forall \mathbf{u} \in V^h(\tilde{\Omega})$ and $\forall q \in Q^h(\tilde{\Omega})$,

 $0 = \mathbb{NS}[\mu + \mu_T](\boldsymbol{u}, p; \boldsymbol{w}, q) + \mathbb{STAB}[\mu + \mu_T](\boldsymbol{u}, p; \boldsymbol{w}, q)$

- Spalart-Allmaras (SA) model with the Shifted Boundary Method are very similar to the Navier-Stokes equations
- Implicit LES is performed through the VMS stabilization/modeling approach

Wall model for the velocity boundary conditions:
$$u = g - \nabla u \cdot d \implies u = g - u_{wall}(d, \nabla u, \ldots)$$
 $u_{\tau} = \sqrt{\frac{\tau_w}{\rho}}$, $u^+ = \frac{1}{\kappa} \log y^+ + C^+$ $y^+ = \frac{u_{\tau}y}{\gamma}$, $u^+ = \frac{u_{\tau}y}{\kappa}$, $u^+ = \frac{u}{u_{\tau}}$,

Flow Over a Circular Cylinder

A classical test to validate algorithms for laminar/turbulent flow

Re	St (vs. reference)	C_D (vs. reference)	Reference source	
20	-	2.09 (1.99)	[7]	
100	0.167 (0.164,0.157)	1.35 (1.34)	[7, 39]	
300	0.211 (0.203, 0.215)	1.38 (1.37)	[39]	
3900	0.203 (0.203)	1.04 (1.00)	[7]	
^[7] Beaudan & Moin (1995)				



Lift coefficient, Re = 250

Flow Over a Circular Cylinder at Re=3,900

A classical test to validate algorithms for turbulent flow simulation



Q-criterion isosurfaces (used to visualize vortex structures)

Flow Over a Circular Cylinder at Re=3,900

A classical test to validate algorithms for turbulent flow simulation



Q-criterion isosurfaces (used to visualize vortical structures)

Rentered by LitineViz.

A More Complicated Shape at Re=3,900

A differential geometry monster (the Monkey Trefoil)



A More Complicated Shape at Re=3,900

A differential geometry monster (the Monkey Trefoil)



Q-criterion isosurfaces (used to visualize vortical structures)
The SBM for Solid Mechanics An alternative cycle to CAD-Meshing-Computing



Original geometry



The SBM for Solid Mechanics

An alternative cycle to CAD-Meshing-Computing



The background domain



The SBM for Solid Mechanics An alternative cycle to CAD-Meshing-Computing



The background domain and the immersed original geometry



The SBM for Solid Mechanics An alternative cycle to CAD-Meshing-Computing



The initial set of active elements (with boundary conditions sidesets)

The SBM for Solid Mechanics



Deformed configuration of the set of active elements



The SBM for Solid Mechanics An alternative cycle to CAD-Meshing-Computing



The intersection of the immersed geometry with the grid



The SBM for Solid Mechanics An alternative cycle to CAD-Meshing-Computing



Deformation of the intersected immersed geometry



Credit of SBM for Solid Mechanics



Ph.D. thesis of Nabil M. Atallah



The SBM for Solid Mechanics

Key challenges:

- Imposition of traction (Neumann) boundary conditions
- Numerical stability of Neumann conditions
- Effect of geometric complexity (edges/corners, etc.) on numerical stability
- Computational cost vs. standard Finite Element Method



The SBM for Solid Mechanics

Static linear elasticity problem

$$0 = b + \nabla \cdot \sigma(\nabla^{s} u) \quad \text{(governing eq.)}$$

$$u = u_{D} \quad \text{(displacement BCs)}$$

$$\sigma(\nabla^{s} u)n = t_{N} \quad \text{(traction BCs)}$$

$$= 0 \text{ for piecewise linear fcn.s}$$

$$\sigma(\varepsilon) = \mathbb{C} \varepsilon \quad \nabla^{s} u = 1/2 \left(\nabla u + \nabla^{t} u\right)$$



Neumann [Traction] BC

Idea: Recast the linear elasticity problem in mixed form:

 $0 = b + \nabla \cdot \sigma(\varepsilon) \text{ (governing eq.)}$ $\varepsilon = \nabla^{s} u \text{ (strain eq.)}$ $u + (\nabla u)d = u_{D} \text{ (displacement BCs)}$ $(\sigma + \nabla \sigma d) n = (C \varepsilon + \nabla(C \varepsilon) d) n = t_{N} \text{ (traction BCs)}$ $\neq 0$

Efficient Implementation of SBM for Solid Mechanics

Apply the mixed formulation only on a layer one-element thick



The SBM for Domains with Corners

How to handle multiple boundary conditions near corners/edges?





Mapping boundary conditions

Cook's Membrane (Compressible Elasticity)



Young's modulus $E = 200 \cdot 10^9$ Pa and Poisson's ratio $\nu = 0.3$.

Cook's Membrane (Compressible Elasticity)

Tip displacement convergence



2,500

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Three-Dimensional Performance Assessment



Young's modulus= 200.10⁶ MPa

Poisson ratio = 0.3

Boundary conditions

- Strong enforcement of Dirichlet conditions on outer boundary
- Traction conditions on inner boundaries

Three-Dimensional Performance Assessment



SBM costs you like a standard body-fitted FEM, but... no meshing overhead

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Pulling/Twisting of Three-Dimensional Stent

No body force

Boundary conditions

• *u* = 0 (left),

Du

- u = {2.5 y, -2.5 x, -0.5}(right)
- Traction free BC at all other boundaries

Poisson ratio = 0.452





Pulling/Twisting of Three-Dimensional Stent

Solution (displacement)







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Rock REV: L = 0.505, W =0.3, D = 0.27

No body force

Boundary conditions

- **u** = 0 (bottom),
- *u* = -5E-04 (top)
- Traction free BC at all other boundaries

Poisson ratio = 0.452

Total elements = 5.2 million Total active elements = 3.46 million Total cores = 240

Original geometry

Rock REV: L = 0.505, W =0.3, D = 0.27

Setup time:

- Compute intersection between STL surface and background grid.
- Generate surrogate domain and boundaries.
- Calculate distance vector **d**.

Setup time = 220 seconds!

Surrogate domain (blue)

Rock REV: L = 0.505, W =0.3, D = 0.27

Setup time:

- Compute intersection between STL surface and background grid.
- Generate surrogate domain and boundaries.
- Calculate distance vector **d**.

Setup time = 220 seconds!

Solution (displacement)

Rock REV: L = 0.50

No body force

Boundary conditions

- *u* = 0 (left),
- *u* = -5E-04 (right)
- Traction free BC at other boundaries

Original geometry

J₂ stress

Rock REV: L = 0.505, W = 0.3, D = 0.27

No body force

Boundary conditions

- *u* = 0 (left),
- *u* = -5E-04 (right)
- Traction free BC at all other boundaries

Rock REV: L = 0.505, W = 0.3, D = 0.27

No body force

Boundary conditions

- u = 0 (left),
- *u* = -5E-04 (right)
- Traction free BC at all other boundaries

J₂ stress along X-Y plane

Rock REV: L = 0.505, W = 0.3, D = 0.27

No body force

Boundary conditions

- *u* = 0 (left),
- *u* = -5E-04 (right)
- Traction free BC at all other boundaries

J₂ stress along X-Z plane

Rock REV: L = 0.505, W = 0.3, D = 0.27

No body force

Boundary conditions

- *u* = 0 (left),
- *u* = -5E-04 (right)
- Traction free BC at all other boundaries

J₂ stress along Y-Z plane

"Sponge-like" boundary test

Geometry

Surrogate boundary



traction-free + flux free

"Sponge-like" boundary test



Temperature T, t = 5.03



"Sponge-like" boundary test



Displacement $\boldsymbol{u}, t = 5.03$

Displacement $\boldsymbol{u}, t = 10$

Simulation on a complex structure (Naval Research Laboratory):



Fixed temperature BCs + clamped displacements

Simulation on a complex structure (Naval Research Laboratory):

Temperature



Fixed temperature BCs + clamped displacements

Simulation on a complex structure (Naval Research Laboratory):



Simulation on a complex structure (Naval Research Laboratory):



Simulation on a complex structure (Naval Research Laboratory):

Displacement, z-component 110

The Shifted Fracture Method

[Collaboration with Antonio Rodriguez-Ferran Universitat Politècnica de Catalunya]

The Shifted Fracture Method (SFM) in a nutshell:

- An attempt to overcome the challenges in the X-FEM/G-FEM/PU-FEM approach maintaining a comparable accuracy
- A surrogate fracture is introduced and constituted of edges/faces of the mesh that are in some sense close to the true fracture
- Fracture interface conditions are computed on the *surrogate fracture*
- The value of interface conditions is *shifted* (or perturbed) by means of Taylor expansions, to preserve accuracy and avoid mesh dependencies

Advantages of the Shifted Fracture Method:

- No cut cells: simplified data-structure and quadratures
- No small-cut cell problem affecting stability and matrix conditioning
- SFM solutions are mesh independent, unlike node-release techniques
- Continuous representation of the fracture surface in three dimensions

Credit for Shifted Fracture Method & Thermomechanics



Ph.D. thesis of Kangan Li (... looking for jobs)[in collaboration with Antonio Rodriguez-FerranCUNIVERSITYUniversitat Politècnica de Catalunya]
Overview: Cohesive Zone Models $[x] = x^{-} - x^{+} \{x\}_{\lambda} = \lambda x^{+} + (1 - \lambda)x^{-}$

Governing equations

Linear elasticity in mixed form $0 = \boldsymbol{b} + \nabla \cdot \boldsymbol{\sigma}(\boldsymbol{\epsilon}) ,$

 $\boldsymbol{\varepsilon} = \nabla^{s} \boldsymbol{u}$.



Cohesive zone model [Pandolfi-Ortiz, IJNME 44 (1999)]

CZMs were initially proposed by Dugdale, Barenblatt, Hillerborg, Bazant, et al.:

$$t \equiv t^{+} = \sigma(\varepsilon)^{+} \cdot n^{+} = -\sigma(\varepsilon)^{-} \cdot n^{-} = -t^{-} \text{ on } \Gamma_{c}$$

$$t = t_{coh}(w) \text{ on } \Gamma_{c}$$

$$w(u) = u^{-} - u^{+} = \llbracket u \rrbracket$$

$$t_{eq}(w_{eq}) = \begin{cases} t_{cr} - Tw_{eq} & \text{for } 0 < w_{eq} = w^{*} < w_{cr} , & \dot{w}^{*} > 0 , \\ \frac{t^{*}}{w^{*}} w_{eq} , & \text{for } 0 \le w_{eq} < w^{*} < w_{cr} , & \dot{w}^{*} = 0 , \\ 0, & \text{for } w_{eq} \ge w_{cr} , \end{cases}$$

$$t_{coh} = t_{n}n + t_{s}\tau$$

$$t_{n} = t_{eq}(w_{eq})\frac{w_{n}}{w_{eq}} , \\ t_{n} = t_{eq}(w_{eq})\frac{w_{n}}{w_{eq}} , \\ t_{s} = \beta^{2} t_{eq}(w_{eq})\frac{w_{s}}{w_{eq}} , \\ t_{s} = \beta^{2} t_{eq}(w_{eq})\frac{w_{s}}{w_{eq}} , \end{cases}$$

Overview: Cohesive Zone Models

X-FEM/PU-FEM variational formulation [Moës-Belytschko or Wells-Sluys]

Variational equations

$$(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \Gamma_c} + (\boldsymbol{\psi}, \boldsymbol{\varepsilon} - \nabla^s \boldsymbol{u})_{\Omega \setminus \Gamma_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \boldsymbol{t}_{coh} \rangle_{\Gamma_c} - (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \Gamma_c} - \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N} = 0$$

Euler-Lagrange equations on the fracture surface

 $\{\!\!\{\sigma\}\!\!\}_{\lambda} n - t_{coh} = 0, \quad \text{on } \Gamma_c \ \text{Cohesive traction} \\ \[\![\sigma]\!] n = 0, \quad \text{on } \Gamma_c \ \text{Normal stress continuity/balance}$



$$\|\chi\| = \chi^{-} - \chi^{+} \qquad \|\chi\|_{\lambda} = \lambda \chi^{+} + (1 - \lambda)\chi$$

$$n^{+} \equiv n$$

$$\Gamma_{c}^{-} \qquad \Gamma_{coh}^{-} \qquad A$$

 Γ_c^+

m ... m

Geometric setting of the Shifted Fracture Approach

Surrogate fracture: the set of edges/faces closest to the true crack (in red)

d: the distance vector of the surrogate crack from the true crack

 \tilde{A} : the surrogate crack tip, corresponding to A



Idea of the Shifted Fracture Method: use Taylor expansions to evaluate *approximate* fracture conditions on the surrogate fracture

$$\bar{t} = t_{coh}(w(\bar{u})) = \bar{\sigma}\bar{n} \approx (\sigma + \nabla\sigma d) \bar{n} \approx t_{coh}(w(u + \nabla u d))$$

on $\widetilde{\Gamma}_{coh}$.

Derivation of the variational equations

Function spaces (piecewise linear FEMs, discontinuous across the <u>surrogate</u> crack) $\mathscr{S}_{u}^{h}(\Omega \setminus \tilde{\Gamma}_{c}) = \{ \boldsymbol{v} \in (C^{0}(\Omega \setminus \tilde{\Gamma}_{c}))^{d} : \boldsymbol{v}_{|\Gamma_{D}} = \boldsymbol{u}_{D}, \boldsymbol{v} \mid_{T} \in (\mathscr{P}^{1}(T))^{d}, \forall T \in \mathscr{T}_{h} \},$ $\mathscr{S}_{\varepsilon}^{h}(\Omega \setminus \tilde{\Gamma}_{c}) = \{ \boldsymbol{\omega} \in (C^{0}(\Omega \setminus \tilde{\Gamma}_{c}))^{d \times d} : \boldsymbol{\omega} \mid_{T} \in (\mathscr{P}^{1}(T))^{d \times d}, \boldsymbol{\omega} = \boldsymbol{\omega}^{t}, \forall T \in \mathscr{T}_{h} \},$

 $\begin{aligned} \mathcal{V}_{\boldsymbol{u}}^{h}(\Omega \setminus \tilde{\Gamma}_{c}) &= \{ \boldsymbol{v} \in (C^{0}(\Omega \setminus \tilde{\Gamma}_{c}))^{d} : \boldsymbol{v}_{|\Gamma_{D}} = \boldsymbol{0} , \boldsymbol{v} \mid_{T} \in (\mathscr{P}^{1}(T))^{d}, \forall T \in \mathcal{T}_{h} \} , \\ \mathcal{V}_{\epsilon}^{h}(\Omega \setminus \tilde{\Gamma}_{c}) &= \mathcal{S}_{\epsilon}^{h}(\Omega \setminus \tilde{\Gamma}_{c}) . \end{aligned}$

Galerkin projection:

$$\begin{aligned} -\left(\boldsymbol{\phi},\nabla\cdot\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\right)_{\Omega\setminus\tilde{\Gamma_c}} &= \left(\boldsymbol{\phi},\boldsymbol{b}\right)_{\Omega\setminus\tilde{\Gamma_c}},\\ \left(\boldsymbol{\psi},\boldsymbol{\varepsilon}\right)_{\Omega\setminus\tilde{\Gamma_c}} &= \left(\boldsymbol{\psi},\nabla^s\boldsymbol{u}\right)_{\Omega\setminus\tilde{\Gamma_c}},\end{aligned}$$

Integrating by parts ...

 $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \{\!\!\{ \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \}\!\!\}_{\lambda} \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c}$

 $+ \langle \{\!\!\{ \boldsymbol{\phi} \}\!\!\}_{1-\lambda}, [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!\!] \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} = (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \tilde{\Gamma}_c} + \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N},$



Shifted interface conditions

 $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}_{c}} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \{\!\!\{ \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \}\!\!\}_{\lambda} \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_{c}} + \langle \{\!\!\{ \boldsymbol{\phi} \}\!\!\}_{1-\lambda}, \llbracket \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \rrbracket \!\!] \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_{c}} = (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \tilde{\Gamma}_{c}} + \langle \boldsymbol{\phi}, \boldsymbol{t}_{N} \rangle_{\Gamma_{N}}$





Shifted interface conditions

 $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \{\!\!\{\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\}\!\!\}_{\lambda} \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} + \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \llbracket \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \rrbracket \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} = (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \tilde{\Gamma}_c} + \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N}$

 $\{\!\{\sigma(\varepsilon)\}\!\}_{\lambda} \tilde{n} = \{\!\{\sigma(\varepsilon)\}\!\}_{\lambda} ((\tilde{n} \cdot n)n + (\tilde{n} \cdot \tau)\tau) = (\tilde{n}(\tilde{x}) \cdot \bar{n}(\tilde{x})) \{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda} \bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x}) \cdot \bar{\tau}(\tilde{x})) \{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda} \bar{\tau}(\tilde{x}) \}$





Shifted interface conditions

 $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \, \{\!\!\{\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\}\!\!\}_{\lambda} \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} + \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \, \llbracket \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \rrbracket \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} = (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \tilde{\Gamma}_c} + \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N}$

 $\{\!\{\sigma(\varepsilon)\}\!\}_{\lambda}\tilde{n} = \{\!\{\sigma(\varepsilon)\}\!\}_{\lambda}((\tilde{n}\cdot n)n + (\tilde{n}\cdot\tau)\tau)$ = $(\tilde{n}(\tilde{x})\cdot\bar{n}(\tilde{x}))\{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda}\bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x})\cdot\bar{\tau}(\tilde{x}))\{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda}\bar{\tau}(\tilde{x})$

 $\{\!\!\{\sigma(\varepsilon(\tilde{x}))\}\!\!\}_{\lambda} \bar{n}(\tilde{x}) = \{\!\!\{\sigma \bar{n}\}\!\!\}_{\lambda} = \{\!\!\{(\bar{\sigma} - \nabla \sigma d) \bar{n}\}\!\!\}_{\lambda} \\ = \{\!\!\{\bar{t} - (\nabla \sigma d) \bar{n}\}\!\!\}_{\lambda} = \bar{t} - \{\!\!\{(\nabla \sigma d)\}\!\!\}_{\lambda} \bar{n},$





Shifted interface conditions

 $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \, \{\!\!\{\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\}\!\!\}_{\lambda} \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} + \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \, \llbracket \boldsymbol{\sigma}(\boldsymbol{\varepsilon}) \rrbracket \tilde{\boldsymbol{n}} \rangle_{\tilde{\Gamma}_c} = (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \tilde{\Gamma}_c} + \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N}$

 $\{\!\{\sigma(\varepsilon)\}\!\}_{\lambda}\tilde{n} = \{\!\{\sigma(\varepsilon)\}\!\}_{\lambda}((\tilde{n}\cdot n)n + (\tilde{n}\cdot\tau)\tau)$ = $(\tilde{n}(\tilde{x})\cdot\bar{n}(\tilde{x}))\{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda}\bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x})\cdot\bar{\tau}(\tilde{x}))\{\!\{\sigma(\varepsilon(\tilde{x}))\}\!\}_{\lambda}\bar{\tau}(\tilde{x})$

$$\{\!\!\{\sigma(\varepsilon(\tilde{x}))\}\!\!\}_{\lambda} \bar{n}(\tilde{x}) = \{\!\!\{\sigma \bar{n}\}\!\!\}_{\lambda} = \{\!\!\{\bar{\sigma} - \nabla \sigma d\} \bar{n}\}\!\!\}_{\lambda}$$

$$= \{\!\!\{\bar{t} - (\nabla \sigma d) \bar{n}\}\!\!\}_{\lambda} = \bar{t} - \{\!\!\{(\nabla \sigma d)\}\!\!\}_{\lambda} \bar{n},$$

$$\bar{t} = t_{coh}(w(\bar{u})) \approx t_{coh}(w(u + \nabla u d)), \text{ on } \tilde{\Gamma}_{coh}$$





Shifted interface conditions

$$(\nabla \phi, \sigma(\varepsilon))_{\Omega \setminus \tilde{\Gamma}_{c}} + \langle \llbracket \phi \rrbracket, \llbracket \sigma(\varepsilon) \rrbracket_{\lambda} \tilde{n} \rangle_{\tilde{\Gamma}_{c}} + \langle \llbracket \phi \rrbracket_{1-\lambda}, \llbracket \sigma(\varepsilon) \rrbracket \tilde{n} \rangle_{\tilde{\Gamma}_{c}} = (\phi, b)_{\Omega \setminus \tilde{\Gamma}_{c}} + \langle \phi, t_{N} \rangle_{\Gamma_{N}}$$

$$\frac{\llbracket \sigma(\varepsilon) \rrbracket_{\lambda} \tilde{n}}{= \llbracket \sigma(\varepsilon) \rrbracket_{\lambda} ((\tilde{n} \cdot n)n + (\tilde{n} \cdot \tau)\tau)$$

$$= (\tilde{n}(\tilde{x}) \cdot \bar{n}(\tilde{x})) \llbracket \sigma(\varepsilon(\tilde{x})) \rrbracket_{\lambda} \bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x}) \cdot \bar{\tau}(\tilde{x})) \llbracket \sigma(\varepsilon(\tilde{x})) \rrbracket_{\lambda} \bar{\tau}(\tilde{x})$$

$$\frac{\llbracket \sigma(\varepsilon(\tilde{x})) \rrbracket_{\lambda} \bar{n}(\tilde{x}) = \llbracket \sigma \bar{n} \rrbracket_{\lambda} = \llbracket (\bar{\sigma} - \nabla \sigma d) \bar{n} \rrbracket_{\lambda}$$

$$= \llbracket \bar{t} - (\nabla \sigma d) \bar{n} \rrbracket_{\lambda} = \bar{t} - \llbracket (\nabla \sigma d) \rrbracket_{\lambda} \bar{n},$$

$$\bar{t} = t_{coh}(w(\bar{u})) \approx t_{coh}(w(u + \nabla u d)), \text{ on } \tilde{\Gamma}_{coh}$$

$$\bar{u} \approx u + \nabla u d,$$

$$\bar{\sigma} \approx \sigma + \nabla \sigma d,$$

$$\langle \llbracket \phi \rrbracket, \llbracket \sigma(\varepsilon) \rrbracket_{\lambda} \tilde{n} \rangle_{\tilde{\Gamma}_{c}} = \langle \llbracket \phi \rrbracket, (\tilde{n} \cdot \bar{n}) (\bar{t} - \llbracket (\nabla \sigma(\varepsilon) d) \rrbracket_{\lambda} \bar{n}) + (\tilde{n} \cdot \bar{\tau}) \llbracket \sigma(\varepsilon) \rrbracket_{\lambda} \bar{\tau} \rangle_{\tilde{\Gamma}_{c}}$$

 $= \langle \llbracket \phi \rrbracket, (\tilde{n} \cdot \bar{n}) t_{coh}(w(u + \nabla u \, d)) \rangle_{\bar{\Gamma}_{coh}} - \langle \llbracket \phi \rrbracket, (\tilde{n} \cdot \bar{n}) (\{ [\nabla \sigma(\varepsilon) \, d) \}_{\lambda} \bar{n}) - (\tilde{n} \cdot \bar{\tau}) \{ \{ \sigma(\varepsilon) \}_{\lambda} \bar{\tau} \rangle_{\bar{\Gamma}_{c}} \}$



Shifted interface conditions

$$\begin{split} (\nabla\phi,\sigma(\varepsilon))_{\Omega\setminus\tilde{\Gamma}_{c}} + \langle \llbracket\phi\rrbracket, \llbracket\sigma(\varepsilon)\rrbracket_{\lambda}\tilde{n}\rangle_{\tilde{\Gamma}_{c}} + \langle \llbracket\phi\rrbracket_{1-\lambda}, \llbracket\sigma(\varepsilon)\rrbracket\tilde{n}\rangle_{\tilde{\Gamma}_{c}} = (\phi,b)_{\Omega\setminus\tilde{\Gamma}_{c}} + \langle\phi,t_{N}\rangle_{\Gamma_{N}} \\ & \underbrace{\llbracket\sigma(\varepsilon)\rrbracket_{\lambda}\tilde{n}} = \llbracket\sigma(\varepsilon)\rrbracket_{\lambda}((\tilde{n}\cdot n)n + (\tilde{n}\cdot\tau)\tau) \\ & = (\tilde{n}(\tilde{x})\cdot\bar{n}(\tilde{x})) \llbracket\sigma(\varepsilon(\tilde{x}))\rrbracket_{\lambda}\bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x})\cdot\bar{\tau}(\tilde{x})) \llbracket\sigma(\varepsilon(\tilde{x}))\rrbracket_{\lambda}\bar{\tau}(\tilde{x}) \\ & \underbrace{\llbracket\sigma(\varepsilon(\tilde{x}))\rrbracket_{\lambda}\bar{n}(\tilde{x})} = \llbracket\sigma\bar{n}\rrbracket_{\lambda} = \llbracket(\bar{\sigma}-\nabla\sigma d)\bar{n}\rrbracket_{\lambda} \\ & = \llbracket\bar{t}-(\nabla\sigma d)\bar{n}\rrbracket_{\lambda} = \llbracket(\bar{\sigma}-\nabla\sigma d)\rrbracket_{\lambda}\bar{n}, \\ & \bar{t}=t_{coh}(w(\bar{u})) \approx t_{coh}(w(u+\nabla u d)), \quad \text{on }\tilde{\Gamma}_{coh} \\ & \langle \llbracket\phi\rrbracket, \llbracket\sigma(\varepsilon)\rrbracket_{\lambda}\tilde{n}\rangle_{\Gamma_{c}} = \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})(\bar{t}-\{(\nabla\sigma(\varepsilon)d)\rrbracket_{\lambda}\bar{n}) + (\tilde{n}\cdot\bar{\tau})\{[\sigma(\varepsilon)\rrbracket_{\lambda}\bar{\tau}\rangle_{\Gamma_{c}} \\ & = \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})t_{coh}(w(u+\nabla u d))\rangle_{\Gamma_{ch}} - \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})(\{(\nabla\sigma(\varepsilon)d)\rrbracket_{\lambda}\bar{n}) - (\tilde{n}\cdot\bar{\tau})\{[\sigma(\varepsilon)\rrbracket_{\lambda}\bar{\tau}\rangle_{\Gamma_{c}} \end{split}$$

$$\begin{split} \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!]\tilde{\boldsymbol{n}}\rangle_{\bar{\Gamma}_c} &= \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \, (\tilde{\boldsymbol{n}} \cdot \bar{\boldsymbol{n}}) [\!\![\bar{\boldsymbol{\sigma}}(\boldsymbol{\varepsilon}) - \nabla \boldsymbol{\sigma}(\boldsymbol{\varepsilon})d]\!\!]\bar{\boldsymbol{n}} + (\tilde{\boldsymbol{n}} \cdot \bar{\boldsymbol{\tau}}) [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!]\bar{\boldsymbol{\tau}}\rangle_{\bar{\Gamma}_c} \\ &= \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \, -(\tilde{\boldsymbol{n}} \cdot \bar{\boldsymbol{n}}) [\!\![\nabla \boldsymbol{\sigma}(\boldsymbol{\varepsilon})d]\!]\bar{\boldsymbol{n}} + (\tilde{\boldsymbol{n}} \cdot \bar{\boldsymbol{\tau}}) [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!]\bar{\boldsymbol{\tau}}\rangle_{\bar{\Gamma}_c} \\ &= \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, \, [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!]\tilde{\boldsymbol{n}} - (\tilde{\boldsymbol{n}} \cdot \bar{\boldsymbol{n}}) [\!\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon}) + \nabla \boldsymbol{\sigma}(\boldsymbol{\varepsilon})d]\!]\bar{\boldsymbol{n}}\rangle_{\bar{\Gamma}_c} \ . \end{split}$$

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Shifted interface conditions

$$\begin{split} (\nabla\phi,\sigma(\varepsilon))_{\Omega\setminus\tilde{\Gamma}_{c}} &+ \langle \llbracket\phi\rrbracket, \llbracket\sigma(\varepsilon) \rrbracket_{\lambda}\tilde{n}\rangle_{\tilde{\Gamma}_{c}} + \langle \llbracket\phi\rrbracket_{1-\lambda}, \llbracket\sigma(\varepsilon)\rrbracket\tilde{n}\rangle_{\tilde{\Gamma}_{c}} = (\phi,b)_{\Omega\setminus\tilde{\Gamma}_{c}} + \langle\phi,t_{N}\rangle_{\Gamma_{N}} \\ & \underbrace{\llbracket\sigma(\varepsilon) \rrbracket_{\lambda}\tilde{n}} = \llbracket\sigma(\varepsilon) \rrbracket_{\lambda}((\tilde{n}\cdot n)n + (\tilde{n}\cdot\tau)\tau) \\ &= (\tilde{n}(\tilde{x})\cdot\bar{n}(\tilde{x})) \llbracket\sigma(\varepsilon(\tilde{x})) \rrbracket_{\lambda}\bar{n}(\tilde{x}) + (\tilde{n}(\tilde{x})\cdot\bar{\tau}(\tilde{x})) \llbracket\sigma(\varepsilon(\tilde{x})) \rrbracket_{\lambda}\bar{\tau}(\tilde{x}) \\ & \underbrace{\llbracket\sigma(\varepsilon) \rrbracket_{\lambda}\tilde{n}(\tilde{x})}_{\delta} = \llbracket\sigma\bar{n} \rrbracket_{\lambda} = \llbracket(\bar{\sigma}-\nabla\sigma d)\bar{n} \rrbracket_{\lambda} \\ &= \llbracket\bar{t} - (\nabla\sigma d)\bar{n} \rrbracket_{\lambda} = [\bar{t} - \{(\nabla\sigma d)\}_{\lambda}\bar{n}, \\ & \bar{t} = t_{coh}(w(\bar{u})) \approx t_{coh}(w(u+\nabla u d)), \text{ on } \tilde{\Gamma}_{coh} \\ & \langle \llbracket\phi\rrbracket, [\tilde{n}\circ\varepsilon) \rrbracket_{\lambda}\tilde{n}\rangle_{\bar{\Gamma}_{c}} = \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})(\bar{t} - \{(\nabla\sigma(\varepsilon)d)\}_{\lambda}\bar{n}) + (\tilde{n}\cdot\bar{\tau}) [\![\sigma(\varepsilon)]_{\lambda}\bar{\tau}\rangle_{\bar{\Gamma}_{c}} \\ &= \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})t_{coh}(w(u+\nabla u d))\rangle_{\bar{\Gamma}_{coh}} - \langle \llbracket\phi\rrbracket, (\tilde{n}\cdot\bar{n})(\{(\nabla\sigma(\varepsilon)d)\}_{\lambda}\bar{n}) - (\tilde{n}\cdot\bar{\tau}) [\![\sigma(\varepsilon)]_{\lambda}\bar{\tau}\rangle_{\bar{\Gamma}_{c}} \end{split}$$

 $[\![\bar{\sigma}(\varepsilon)]\!]\,\bar{n}=0$

$$\begin{split} \left\langle \left\{\!\left\{\boldsymbol{\phi}\right\}\!\right\}_{1-\lambda}, \left[\!\left[\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\right]\!\right]\!\tilde{\boldsymbol{n}}\right\rangle_{\bar{\Gamma}_{c}} &= \left\langle \left\{\!\left\{\boldsymbol{\phi}\right\}\!\right\}_{1-\lambda}, \left(\tilde{\boldsymbol{n}}\cdot\bar{\boldsymbol{n}}\right)\!\left[\!\left[\bar{\boldsymbol{\sigma}}(\boldsymbol{\varepsilon})\right]\!-\nabla\boldsymbol{\sigma}(\boldsymbol{\varepsilon})d\right]\!\right]\!\tilde{\boldsymbol{n}} + \left(\tilde{\boldsymbol{n}}\cdot\bar{\boldsymbol{\tau}}\right)\left[\!\left[\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\right]\!\right]\!\tilde{\boldsymbol{\tau}}\right\rangle_{\bar{\Gamma}_{c}} \\ &= \left\langle \left\{\!\left\{\boldsymbol{\phi}\right\}\!\right\}_{1-\lambda}, -\left(\tilde{\boldsymbol{n}}\cdot\bar{\boldsymbol{n}}\right)\left[\!\left[\nabla\boldsymbol{\sigma}(\boldsymbol{\varepsilon})d\right]\!\right]\!\tilde{\boldsymbol{n}} + \left(\tilde{\boldsymbol{n}}\cdot\bar{\boldsymbol{\tau}}\right)\left[\!\left[\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\right]\!\right]\!\tilde{\boldsymbol{\tau}}\right\rangle_{\bar{\Gamma}_{c}} \\ &= \left\langle \left\{\!\left\{\boldsymbol{\phi}\right\}\!\right\}_{1-\lambda}, \left[\!\left[\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\right]\!\right]\!\tilde{\boldsymbol{n}} - \left(\tilde{\boldsymbol{n}}\cdot\bar{\boldsymbol{n}}\right)\left[\!\left[\boldsymbol{\sigma}(\boldsymbol{\varepsilon})\!+\nabla\boldsymbol{\sigma}(\boldsymbol{\varepsilon})d\right]\!\right]\!\tilde{\boldsymbol{n}}\right\rangle_{\bar{\Gamma}_{c}} \\ \end{split}$$

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X-FEM/PU-FEM vs shifted variational formulations

X-FEM/PU-FEM Euler-Lagrange equations:

 $\{\!\!\{\boldsymbol{\sigma}\}\!\!\}_{\lambda} \boldsymbol{n} - \boldsymbol{t}_{coh} = \boldsymbol{0}, \quad \text{on } \boldsymbol{\Gamma}_{c}$ $[\![\boldsymbol{\sigma}]\!] \boldsymbol{n} = \boldsymbol{0}, \quad \text{on } \boldsymbol{\Gamma}_{c}$



Shifted Fracture Method Euler-Lagrange equations:

$$\begin{split} & \underbrace{\tilde{n} \cdot n} \left(\left\{ \left\{ S(\sigma(\varepsilon)) \right\}_{\lambda} n - t_{coh}(w(S(u))) = 0, & \text{on } \tilde{\Gamma}_{c} \right\} \right\}_{\lambda} n - t_{coh}(w(S(u))) = 0, & \text{on } \tilde{\Gamma}_{c} \end{split} \\ & \underbrace{\tilde{n} \cdot n} \left[\left[S(\sigma(\varepsilon)) \right] \right] n = 0, & \text{on } \tilde{\Gamma}_{c} \end{aligned} \\ \end{split}$$

X-FEM/PU-FEM weak form:

$$(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \Gamma_c} + (\boldsymbol{\psi}, \boldsymbol{\varepsilon} - \nabla^s \boldsymbol{u})_{\Omega \setminus \Gamma_c} + \langle \llbracket \boldsymbol{\phi} \rrbracket, \boldsymbol{t}_{coh} \rangle_{\Gamma_c} - (\boldsymbol{\phi}, \boldsymbol{b})_{\Omega \setminus \Gamma_c} - \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N} = 0$$

Shifted Fracture Method weak form (VMS stabilization terms omitted):

- $(\nabla \boldsymbol{\phi}, \boldsymbol{\sigma}(\boldsymbol{\varepsilon}))_{\Omega \setminus \tilde{\Gamma}} + (\boldsymbol{\psi}, \boldsymbol{\varepsilon} \nabla^{s} \boldsymbol{u})_{\Omega \setminus \tilde{\Gamma}} + \langle \llbracket \boldsymbol{\phi} \rrbracket, (\tilde{\boldsymbol{n}} \cdot \boldsymbol{n}) t_{coh}(\boldsymbol{w}(\boldsymbol{u} + \nabla \boldsymbol{u} \, \boldsymbol{d})) \rangle_{\tilde{\Gamma}_{c}}$
- $+ \langle \llbracket \phi \rrbracket, \{ \! \{ \sigma(\varepsilon) \} \! \}_{\lambda} \tilde{n} \{ \! \{ \sigma(\varepsilon) + \nabla \sigma(\varepsilon) d \} \! \}_{\lambda} (\tilde{n} \cdot n) n \rangle_{\tilde{\Gamma}_{\alpha}}$
- $+ \langle \{\!\!\{\boldsymbol{\phi}\}\!\!\}_{1-\lambda}, [\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon})]\!]\tilde{\boldsymbol{n}} (\tilde{\boldsymbol{n}} \cdot \boldsymbol{n}) [\![\boldsymbol{\sigma}(\boldsymbol{\varepsilon}) + \nabla \boldsymbol{\sigma}(\boldsymbol{\varepsilon})\boldsymbol{d}]\!]\boldsymbol{n} \rangle_{\tilde{\Gamma}_c} \langle \boldsymbol{\phi}, \boldsymbol{b} \rangle_{\Omega \setminus \tilde{\Gamma}_c} \langle \boldsymbol{\phi}, \boldsymbol{t}_N \rangle_{\Gamma_N} = 0$

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The Cohesive Shifted Fracture Approach An efficient implementation

Find $[\boldsymbol{u}, \boldsymbol{\varepsilon}] \in V_{\boldsymbol{u},h}(\Omega) \times V_{\boldsymbol{\varepsilon},h}(\tilde{\Omega}_{\boldsymbol{\ell}})$ such that, $\forall [\boldsymbol{\phi}, \boldsymbol{\psi}] \in V_{\boldsymbol{u},h}(\Omega) \times V_{\boldsymbol{\varepsilon},h}(\tilde{\Omega}_{\boldsymbol{\ell}})$,

 $\mathscr{B}_{\rm SFM}([\boldsymbol{u},\boldsymbol{\varepsilon}];[\boldsymbol{\phi},\boldsymbol{\psi}]) = \mathscr{L}_{\rm SFM}([\boldsymbol{\phi},\boldsymbol{\psi}]),$

 $\mathscr{B}_{\rm SFM}([\boldsymbol{u},\boldsymbol{\varepsilon}];[\boldsymbol{\phi},\boldsymbol{\psi}]) = (\mathbf{C}\,\nabla^{s}\boldsymbol{u},\nabla^{s}\boldsymbol{\phi})_{\Omega\setminus\tilde{\Omega}_{\ell}} + \mathscr{B}_{\rm SFM}^{[\Omega_{\ell}\setminus\Gamma_{c}]}([\boldsymbol{u},\boldsymbol{\varepsilon}];[\boldsymbol{\phi},\boldsymbol{\psi}]) ,$ $\mathscr{L}_{\rm SFM}([\boldsymbol{\phi},\boldsymbol{\psi}]) = (\boldsymbol{\phi},\boldsymbol{b})_{\Omega\setminus\tilde{\Omega}_{\ell}} + \langle \boldsymbol{\phi},\boldsymbol{t}_{N}\rangle_{\Gamma_{N}} + \mathscr{L}_{\rm SFM}^{[\tilde{\Omega}_{\ell}\setminus\tilde{\Gamma}_{c}]}([\boldsymbol{\phi},\boldsymbol{\psi}]) .$

Idea: The SFM mixed formulation is applied only on a layer of elements near the fracture, and a primal formulation is applied everywhere else

Euler-Lagrange coupling conditions

 $\begin{aligned} -\nabla \cdot \boldsymbol{\sigma}(\nabla^{s}\boldsymbol{u}) &= \boldsymbol{b} , & \text{in } \Omega \setminus \tilde{\Omega}_{\ell} ,\\ \boldsymbol{\sigma}(\nabla^{s}\boldsymbol{u}^{+}) \, \boldsymbol{n}^{+} + \boldsymbol{\sigma}(\boldsymbol{\varepsilon})^{-} \, \boldsymbol{n}^{-} &= \boldsymbol{0} , & \text{on } \tilde{\Gamma}_{\ell} ,\\ (\boldsymbol{C} \, \nabla^{s}\boldsymbol{u}) \, \boldsymbol{n} &= \boldsymbol{t}_{N} , & \text{on } \Gamma_{N} \cap \partial(\Omega \setminus \tilde{\Omega}_{\ell}) , \end{aligned}$

Three-point bending test: true and surrogate crack paths



Note: The true crack path is directly estimated using the maximum principal stress criterion (i.e., it is not given *a priori*)



Numerical Examples Three-point bending test: Mixed vs. efficient implementations



Meshes	Number of elements	Number of nodes	Average element size around crack path
Mesh 1	1,565	851	5 mm
Mesh 2	3,477	1,828	2.5 mm
Mesh 3	10,513	5,385	1.25 mm
Mesh 4	19,318	9,774	0.625 mm

Three-point bending test: SFM vs. node-release technique



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Three-point bending test: SFM vs. node-release technique



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Three-point bending test: SFM vs. node-release technique



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Three-point bending test for fracture toughness = 50 & 200 (ductile)



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L-shaped panel:



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L-shaped panel:



(a) SFM true crack path for the L-shape panel test with SIF.

(b) Load-deflection curve for the L-shape panel test with SIF.

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Numerical Examples L-shaped panel:



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Four-point bending test:





Four-point bending test:



(a) SFM true crack path computed with the maximal tensile principal stress criterion.

(b) SFM true crack path computed with the stress intensity factor criterion.

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Four-point bending test:



(c) Load-deflection curves for the loads F_1 and F_2 (maximal tensile principal stress criterion).



(d) Load-deflection curves for the loads F_1 and F_2 (SIF criterion).

Numerical Examples Four-point bending test: SFM vs. node-release technique



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Three-dimensional pull-out test



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Three-dimensional pull-out test



Note: The surrogate and estimate of the true fractures are always continuous



Three-dimensional pull-out test



Note: The surrogate and estimate of the true fractures are always continuous





Three-dimensional pull-out test



(a) Linear and exponential cohesive zone model.

(b) SFM: Linear cohesive zone model.

Meshes	Number of elements	Number of nodes	Average element size around crack path
Mesh 1	10,690	2,176	50 mm
Mesh 2	13,122	2,745	30 mm
Mesh 3	23,443	4,791	15 mm

The Shifted Boundary Method

Theoretical developments for the following equations:

Poisson (with Alex Main, Nabil Atallah & Claudio Canuto) Darcy (with Nabil Atallah & Claudio Canuto) Stokes (with Nabil Atallah & Claudio Canuto) Advection-diffusion (with Alex Main)

Linear Elasticity (wih Nabil Atallah)

Computational developments for the following equations:

Thermo-mechanics (with Kangan Li)

Navier-Stokes (with Alex Main)

Free-surface flow (<u>not discussed here</u>, with Alex Main, Léo Nouveau and Oriol Colomés) Acoustics and shallow water equations (<u>not discussed here</u>, with Ting Song and Alex Main) Fracture mechanics (with Kangan Li, Nabil Atallah, and Antonio Rodriguez-Ferran)

Under development, ongoing work:

Nonlinear mechanics (with Nabil Atallah, Vladimir Tomov & Bojan Lazarov) Shock hydrodynamics (with Nabil Atallah, Vladimir Tomov & Bojan Lazarov) Higher-order SBM (with Nabil Atallah, Claudio Canuto, Vladimir Tomov & Bojan Lazarov) Moving boundaries (with Danjie Xu, Oriol Colomés, Léo Nouveau, & Vladimir Tomov)