

# A Continuous Interior Penalty Method Framework for Sixth Order Cahn-Hilliard-type Equations with applications to microstructure evolution and microemulsions

Natasha S. Sharma

Department of Mathematical Sciences

The University of Texas at El Paso

Funding: NSF DMS-2110774

Joint work with Amanda Diegel (Mississippi State University)



FEM@LLNL Seminar Series  
July 18, 2023

# Sixth Order Cahn-Hilliard-type Equations

Applications

# Sixth Order Cahn-Hilliard-type Equations

Applications

# Sixth Order Cahn-Hilliard-type Equations

## Applications

- Microstructure evolution specifically formation of defects within microstructures.



# Sixth Order Cahn-Hilliard-type Equations

## Applications

- Microstructure evolution specifically formation of defects within microstructures.
- Microemulsions capturing the dynamics of phase transitions in a oil-water-surfactant mixtures with applications to designing drug delivery systems.

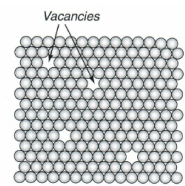
# Sixth Order Cahn-Hilliard-type Equations

## Applications

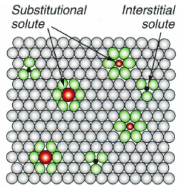
- Microstructure evolution specifically formation of defects within microstructures.
- Microemulsions capturing the dynamics of phase transitions in a oil-water-surfactant mixtures with applications to designing drug delivery systems.

Goal: Develop a Continuous Interior Penalty Method framework to solve these Sixth-Order Phase Field Models.

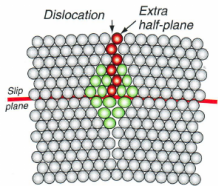
# Defects in Crystalline Materials



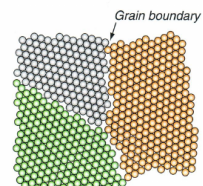
(a)



(b)



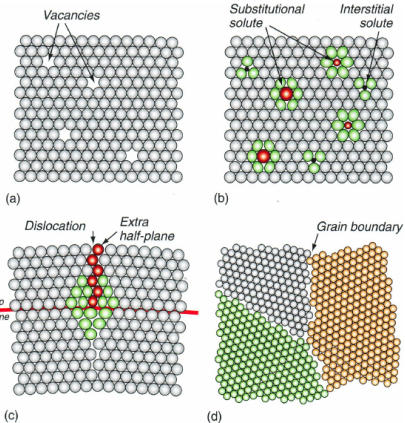
(c)



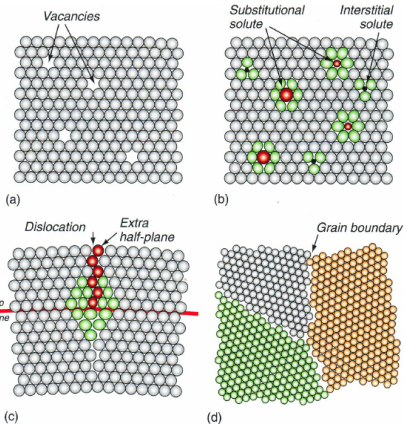
(d)

# Defects in Crystalline Materials

## Defects control



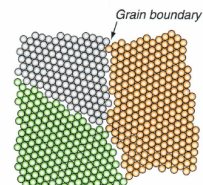
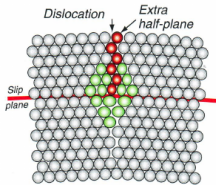
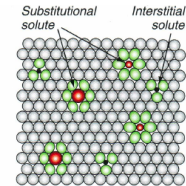
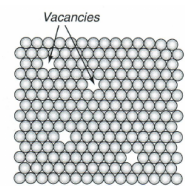
# Defects in Crystalline Materials



## Defects control

- 1 electrical conductivity whether they make efficient solar panels,

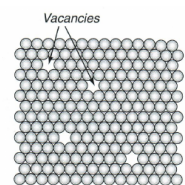
# Defects in Crystalline Materials



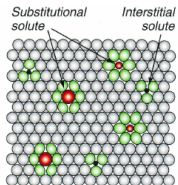
## Defects control

- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity

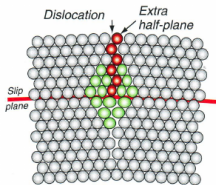
# Defects in Crystalline Materials



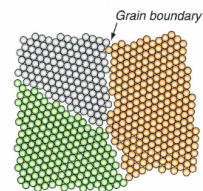
(a)



(b)



(c)

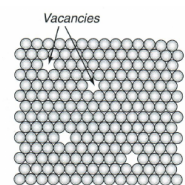


(d)

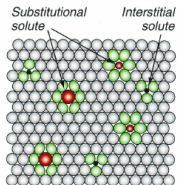
## Defects control

- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material

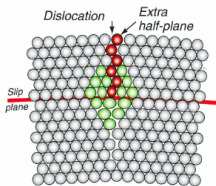
# Defects in Crystalline Materials



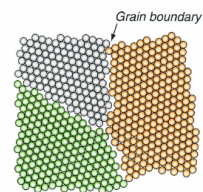
(a)



(b)



(c)



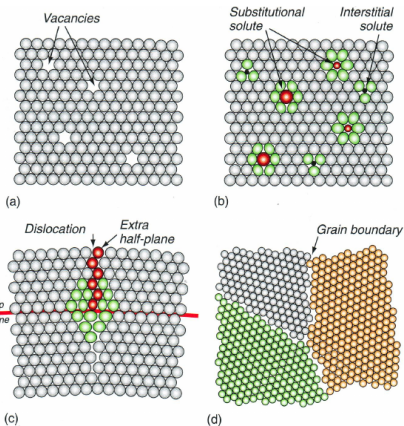
(d)

## Defects control

- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material



# Defects in Crystalline Materials

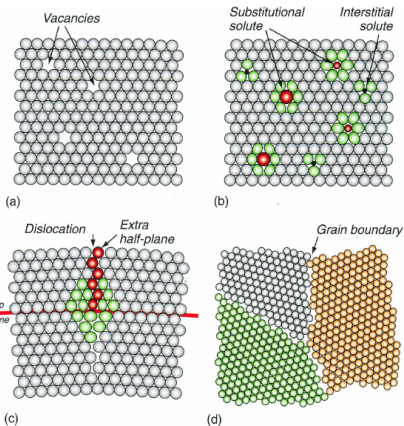


## Defects control

- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material

**Goal:** Control/predict the formation and evolution of defects.

# Defects in Crystalline Materials



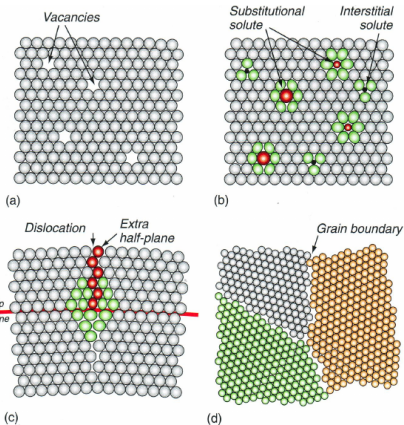
## Defects control

- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material

**Goal:** Control/predict the formation and evolution of defects.

**Approach:**

# Defects in Crystalline Materials



## Defects control

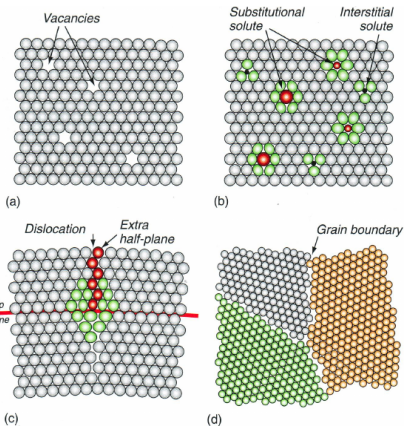
- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material

**Goal:** Control/predict the formation and evolution of defects.

## Approach:

- Use **phase field crystal equation** as our atomistic model.

# Defects in Crystalline Materials



## Defects control

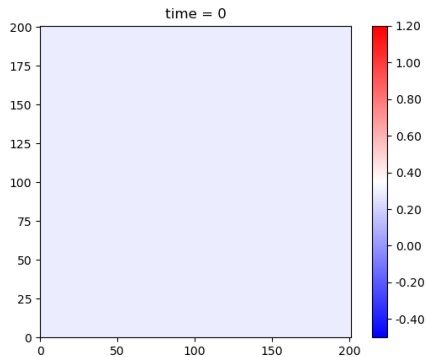
- 1 electrical conductivity whether they make efficient solar panels,
- 2 chemical reactivity
- 3 tensile strength of a material

**Goal:** Control/predict the formation and evolution of defects.

## Approach:

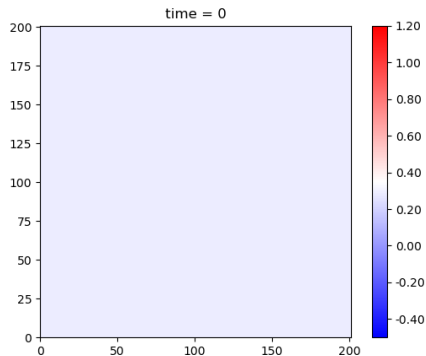
- Use **phase field crystal equation** as our atomistic model.
- Develop an accurate, efficient, easy-to-compute numerical scheme.

# Phase Field Crystal Equation



# Phase Field Crystal Equation

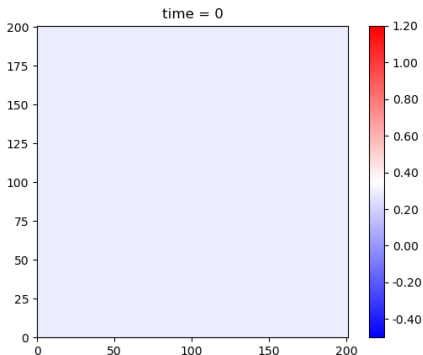
Two-phase system



# Phase Field Crystal Equation

Two-phase system

$\varphi$ : number density of atoms in the material occupying  $\Omega$  with

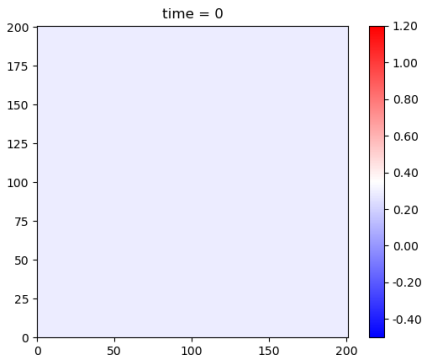


# Phase Field Crystal Equation

Two-phase system

$\varphi$ : number density of atoms in the material occupying  $\Omega$  with

- ◇ liquid phase characterized by a constant value of  $\varphi$



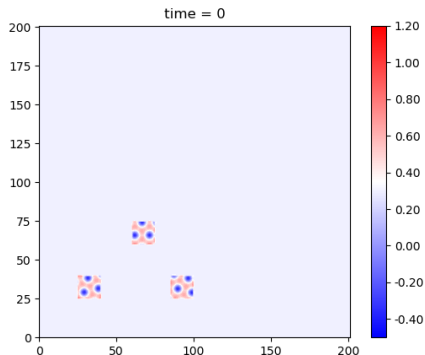


# Phase Field Crystal Equation

## Two-phase system

$\varphi$ : number density of atoms in the material occupying  $\Omega$  with

- ◇ liquid phase characterized by a constant value of  $\varphi$
- ◇ solid phase characterized by a spatially varying periodic function  $\varphi$  that inherits the symmetry and periodicity of the crystal lattice



# Phase Field Crystal Equation

Conservation Law:

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla \mu)$$

# Phase Field Crystal Equation

Conservation Law:

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla \mu)$$

- $\mathcal{M}$ : mobility coefficient assumed to be constant

# Phase Field Crystal Equation

Conservation Law:

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla \mu)$$

- $\mathcal{M}$ : mobility coefficient assumed to be constant
- $\mu = \delta_{\varphi} E$ : chemical potential

# Phase Field Crystal Equation

Conservation Law:

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla \mu)$$

- $\mathcal{M}$ : mobility coefficient assumed to be constant
- $\mu = \delta_{\varphi} E$ : chemical potential

- $E(\varphi) = \int_{\Omega} \frac{\varphi^4}{4} - |\nabla \varphi|^2 + \frac{1-\varepsilon}{2} \varphi^2 + \frac{1}{2} (\Delta \varphi)^2 \, dx$

# Phase Field Crystal Equation

Conservation Law:

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla \mu)$$

- $\mathcal{M}$ : mobility coefficient assumed to be constant
- $\mu = \delta_{\varphi} E$ : chemical potential

- $E(\varphi) = \int_{\Omega} \frac{\varphi^4}{4} - |\nabla \varphi|^2 + \frac{1-\varepsilon}{2} \varphi^2 + \frac{1}{2} (\Delta \varphi)^2 \, dx$

Phase Field Crystal Equation (Elder et al. 2004)

$$\frac{\partial \varphi}{\partial t} = \nabla \cdot (\mathcal{M} \nabla (\varphi^3 + 2\Delta \varphi + (1 - \varepsilon)\varphi + \Delta^2 \varphi)) \quad \text{on } \Omega \times (0, T).$$

# PFC equation in the mixed form

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0,$$

# PFC equation in the mixed form

Compliment

$$\begin{aligned}\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) &= 0, \\ \varphi^3 + (1 - \varepsilon)\varphi + 2\Delta \varphi + \Delta^2 \varphi - \mu &= 0,\end{aligned}$$

with either periodic boundary conditions or natural boundary conditions

$$\partial_n \varphi = \partial_n \Delta \varphi = \partial_n \mu = 0,$$



# PFC equation in the mixed form

Compliment

$$\begin{aligned}\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) &= 0, \\ \varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu &= 0,\end{aligned}$$

with either periodic boundary conditions or natural boundary conditions

$$\partial_n \varphi = \partial_n \Delta\varphi = \partial_n \mu = 0,$$

and the initial value:

$$\varphi(0) = \varphi_0.$$

# PFC equation in the mixed form

Compliment

$$\begin{aligned}\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) &= 0, \\ \varphi^3 + (1 - \varepsilon)\varphi + 2\Delta \varphi + \Delta^2 \varphi - \mu &= 0,\end{aligned}$$

with either periodic boundary conditions or natural boundary conditions

$$\partial_n \varphi = \partial_n \Delta \varphi = \partial_n \mu = 0,$$

and the initial value:

$$\varphi(0) = \varphi_0.$$

Notation:

- $H^s(\Omega)$  denote the Sobolev spaces of order  $s \geq 1$ ,
- $Z := \{z \in H^2(\Omega) \mid \mathbf{n} \cdot \nabla z = 0 \text{ on } \partial\Omega\}$ .

# Weak Formulation

Find  $(\varphi, \mu)$  such that

$$\begin{aligned}\varphi &\in L^\infty(0, T; Z) \cap L^2(0, T; H^3(\Omega)), \\ \partial_t \varphi &\in L^2(0, T; H_N^{-1}(\Omega)), \\ \mu &\in L^2(0, T; H^1(\Omega)),\end{aligned}$$

and for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$\left( (\varphi)^3 + (1 - \epsilon) \varphi, \psi \right) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

$$\text{with } a(u, v) := \int_{\Omega} (\nabla^2 u : \nabla^2 v) \, dx,$$

# Weak Formulation

Find  $(\varphi, \mu)$  such that

$$\begin{aligned}\varphi &\in L^\infty(0, T; Z) \cap L^2(0, T; H^3(\Omega)), \\ \partial_t \varphi &\in L^2(0, T; H_N^{-1}(\Omega)), \\ \mu &\in L^2(0, T; H^1(\Omega)),\end{aligned}$$

and for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$\left( (\varphi)^3 + (1 - \epsilon)\varphi, \psi \right) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

with  $a(u, v) := \int_{\Omega} (\nabla^2 u : \nabla^2 v) \, dx$ ,  $\varphi(0) = \varphi_0 \in H^4(\Omega)$  such that  $\varphi_0$  satisfies the boundary conditions.

[Pawlow et al., 2013]

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- Space discretization: Choose a suitable discretization for the higher order term.

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- Space discretization: Choose a suitable discretization for the higher order term.
- Time discretization: Classical methods give us **conditional** solvability and stability.

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- Space discretization: Choose a suitable discretization for the higher order term.
- Time discretization: Classical methods give us **conditional** solvability and stability.

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$



## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- Finite Element Method for a Mixed form of three second-order equations: Backofen, Rätz, and Voigt 2007

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- **Finite Element Method for a Mixed form of three second-order equations:** Backofen, Rätz, and Voigt 2007
- **Finite Difference Method:** Wise, Wang, and Lowengrub 2009; Dong, Feng, Wang, Wise and Zhang, 2018

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- **Finite Element Method for a Mixed form of three second-order equations:** Backofen, Rätz, and Voigt 2007
- **Finite Difference Method:** Wise, Wang, and Lowengrub 2009; Dong, Feng, Wang, Wise and Zhang, 2018
- **$C^1$  Finite Element Method:** Gomez, Nogueira, 2012

## Phase Field Crystal Equation

$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta\varphi + \Delta^2\varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- **Finite Element Method for a Mixed form of three second-order equations:** Backofen, Rätz, and Voigt 2007
- **Finite Difference Method:** Wise, Wang, and Lowengrub 2009; Dong, Feng, Wang, Wise and Zhang, 2018
- **$C^1$  Finite Element Method:** Gomez, Nogueira, 2012
- **Local Discontinuous Galerkin Method:** Guo and Xu, 2016

## Phase Field Crystal Equation

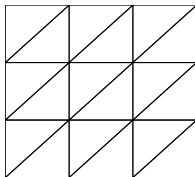
$$\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) = 0,$$
$$\varphi^3 + (1 - \varepsilon)\varphi + 2\Delta \varphi + \Delta^2 \varphi - \mu = 0 \quad \text{on } \Omega \times (0, T).$$

- **Finite Element Method for a Mixed form of three second-order equations:** Backofen, Rätz, and Voigt 2007
- **Finite Difference Method:** Wise, Wang, and Lowengrub 2009; Dong, Feng, Wang, Wise and Zhang, 2018
- **$C^1$  Finite Element Method:** Gomez, Nogueira, 2012
- **Local Discontinuous Galerkin Method:** Guo and Xu, 2016
- **Fourier-spectral Method:** Li and Shen, 2020 (Scalar Auxiliary Variable approach) Yang and Han, 2017 (Invariant Energy Quadratization)

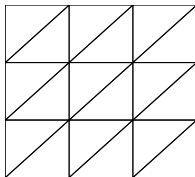
# Numerical Schemes: Our Approach

- Space discretization: Relax the  $C^1$ -continuity, use  $C^0$ -Interior Penalty Method
- Time discretization: Use Eyre's convex splitting scheme known to be uniquely solvable and unconditionally stable

# Spatial Discretization

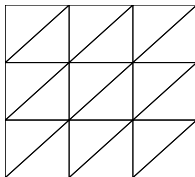


# Spatial Discretization





# Spatial Discretization

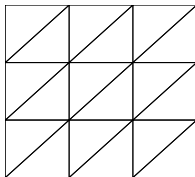


- $\Omega = \bigcup_{K \in \mathcal{T}_h} K,$

**Assume this partition is geometrically-conforming and shape-regular.**

$\mathcal{T}_h$ : collection of all elements  $K$

# Spatial Discretization



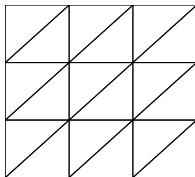
- $\Omega = \bigcup_{K \in \mathcal{T}_h} K,$

**Assume this partition is geometrically-conforming and shape-regular.**

$\mathcal{T}_h$ : collection of all elements  $K$

- $h_K = \text{diameter of triangle } K, h = \max_{K \in \mathcal{T}_h} h_K$

# Spatial Discretization



- $\Omega = \bigcup_{K \in \mathcal{T}_h} K,$

**Assume this partition is geometrically-conforming and shape-regular.**

$\mathcal{T}_h$ : collection of all elements  $K$

- $h_K = \text{diameter of triangle } K, h = \max_{K \in \mathcal{T}_h} h_K$

- $\mathcal{E}_h$ : collection of all edges  $e$  wrt  $\mathcal{T}_h$

# Classical $C^1$ Finite Element Method

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + \mathbf{a}(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

# Classical $C^1$ Finite Element Method

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

Find **finite dimensional** subspaces:

$$V_h \subset H^1(\Omega), \quad Z_h \subset Z,$$

# Classical $C^1$ Finite Element Method

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

Find **finite dimensional** subspaces:

$$V_h \subset H^1(\Omega), \quad Z_h \subset Z,$$

$(\varphi_h, \mu_h) : [0, T] \rightarrow Z_h \times V_h$ :

$$\langle \partial_t \varphi_h, \nu \rangle + (\mathcal{M} \nabla \mu_h, \nabla \nu) = 0 \quad \forall \nu \in V_h$$

$$\left( (\varphi_h)^3 + (1 - \epsilon)\varphi_h, \psi \right) - 2(\nabla \varphi_h, \nabla \psi) + a(\varphi_h, \psi) - (\mu_h, \psi) = 0 \quad \forall \psi \in Z_h$$

holds for almost all  $t \in (0, T)$ .

# Our Approach: Relax $C^1$ continuity

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + \mathbf{a}(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

# Our Approach: Relax $C^1$ continuity

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

Find **finite dimensional** subspaces:

$$V_h \subset H^1(\Omega), \quad Z_h \not\subset Z,$$



# Our Approach: Relax $C^1$ continuity

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

Find **finite dimensional** subspaces:

$$V_h \subset H^1(\Omega), \quad Z_h \not\subset Z,$$

$$Z_h \subset H^1(\Omega)$$

# Our Approach: Relax $C^1$ continuity

Find  $(\varphi, \mu) : [0, T] \rightarrow Z \times H^1(\Omega)$  s.t. for almost all  $t \in (0, T)$

$$\langle \partial_t \varphi, \nu \rangle + (\mathcal{M} \nabla \mu, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega)$$

$$(\varphi^3 + (1 - \epsilon)\varphi, \psi) - 2(\nabla \varphi, \nabla \psi) + a(\varphi, \psi) - (\mu, \psi) = 0 \quad \forall \psi \in Z$$

Find **finite dimensional** subspaces:

$$V_h \subset H^1(\Omega), \quad Z_h \not\subset Z,$$

$$Z_h \subset H^1(\Omega)$$

$(\varphi_h, \mu_h) : [0, T] \rightarrow Z_h \times V_h$ :

$$\langle \partial_t \varphi_h, \nu \rangle + (\mathcal{M} \nabla \mu_h, \nabla \nu) = 0,$$

$$\left( (\varphi_h)^3 + (1 - \epsilon)\varphi_h, \psi \right) - 2(\nabla \varphi_h, \nabla \psi) + a_h^{IP}(\varphi_h, \psi) - (\mu_h, \psi) = 0$$

$\forall \nu \in V_h, \psi \in Z_h$  holds for almost all  $t \in (0, T)$ .

# Spatial Discretization using $C^0$ -IP Method

$$V_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_1(K) \forall K \in \mathcal{T}_h\}$$

# Spatial Discretization using $C^0$ -IP Method

$$V_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_1(K) \forall K \in \mathcal{T}_h\}$$

$$Z_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_2(K) \forall K \in \mathcal{T}_h\},$$

# Spatial Discretization using $C^0$ -IP Method

$$V_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_1(K) \forall K \in \mathcal{T}_h\}$$

$$Z_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_2(K) \forall K \in \mathcal{T}_h\},$$

$R_h : H^1(\Omega) \rightarrow V_h$  is a Ritz projection operator such that

$$(\nabla(R_h\mu - \mu), \nabla\xi) = 0 \quad \forall \xi \in V_h, \quad (R_h\mu - \mu, 1) = 0.$$

# Spatial Discretization using $C^0$ -IP Method

$$V_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_1(K) \forall K \in \mathcal{T}_h\}$$

$$Z_h := \{v \in C(\bar{\Omega}) \mid v|_K \in P_2(K) \forall K \in \mathcal{T}_h\},$$

$R_h : H^1(\Omega) \rightarrow V_h$  is a Ritz projection operator such that

$$(\nabla(R_h\mu - \mu), \nabla\xi) = 0 \quad \forall \xi \in V_h, \quad (R_h\mu - \mu, 1) = 0.$$

$P_h : Z \rightarrow Z_h$  is a Ritz projection operator such that

$$a_h^{IP}(P_h\varphi - \varphi, \xi) + (1 - \epsilon)(P_h\varphi - \varphi, \xi) = 0 \quad \forall \xi \in Z_h, \quad (P_h\varphi - \varphi, 1) = 0.$$

# Spatial Discretization using $C^0$ -IP Method

$a_h^{IP} : Z_h \times Z_h \rightarrow \mathbb{R}$  according to

$$a_h^{IP}(\xi_h, \psi_h) := \sum_{K \in \mathcal{T}_h} \int_K \nabla^2 \xi_h : \nabla^2 \psi_h \, dx + J(\xi_h, \psi_h), \quad \xi_h, \psi_h \in Z_h$$

where

$$\begin{aligned} J(\xi_h, \psi_h) := & \sum_{e \in \mathcal{E}_h} \int_e \left( [\mathbf{n}_e \cdot \nabla \xi_h]_e \{ \mathbf{n}_e \cdot \nabla^2 \psi_h \mathbf{n}_e \}_e + \{ \mathbf{n}_e \cdot \nabla^2 \xi_h \mathbf{n}_e \}_e [\mathbf{n}_e \cdot \nabla \psi_h]_e \right) ds \\ & + \sum_{e \in \mathcal{E}_h} \int_e \frac{\alpha}{h_e} [\mathbf{n}_e \cdot \nabla \xi_h]_e [\mathbf{n}_e \cdot \nabla \psi_h]_e \, ds, \quad \xi_h, \psi_h \in Z_h \end{aligned}$$

and  $\alpha > 0$  is a penalty parameter.

# Boundedness of $a_h^{IP}(\cdot, \cdot)$

## Lemma (Boundedness of $a_h^{IP}(\cdot, \cdot)$ )

There exists positive constants  $C_{cont}$  and  $C_{coer}$  such that for choices of the penalty parameter  $\alpha$  large enough we have

$$\begin{aligned} a_h^{IP}(w_h, v_h) &\leq C_{cont} \|w_h\|_{2,h} \|v_h\|_{2,h} \quad \forall w_h, v_h \in Z_h, \\ C_{coer} \|w_h\|_{2,h}^2 &\leq a_h^{IP}(w_h, w_h) \quad \forall w_h \in Z_h, \end{aligned}$$

where the constants  $C_{cont}$  and  $C_{coer}$  depend only on the shape regularity of  $\mathcal{T}_h$ .

where the  $C^0$ -IP Norm is:

$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [\mathbf{n}_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

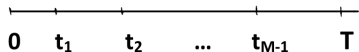


# Time Discretization

We introduce a partition of  $(0,T)$

# Time Discretization

We introduce a partition of  $(0, T)$

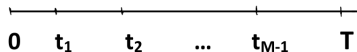


into  $M$  sub-intervals  $(t_{m-1}, t_m)$ :

$$t_m = t_{m-1} + \tau, \text{ for } 1 \leq m \leq M$$

# Time Discretization

We introduce a partition of  $(0, T)$



into  $M$  sub-intervals  $(t_{m-1}, t_m)$ :

$$t_m = t_{m-1} + \tau, \text{ for } 1 \leq m \leq M$$

Notation:  $\varphi^m$  approximate  $\varphi$  at time  $t_m$ .

Numerical time derivative w.r.t.  $\tau$ :

$$\delta_\tau \varphi^m := \frac{\varphi^{m+1} - \varphi^m}{\tau}$$

# Convex Time Splitting Scheme

- Basic Idea:

$$\begin{aligned}\mu &= \delta_\varphi E(\varphi) = \delta_\varphi \left( \underbrace{E^+(\varphi)}_{\text{convex}} + \underbrace{E^-(\varphi)}_{\text{concave}} \right) \\ \implies \mu^m &= \underbrace{(\varphi^m)^3 + (1 - \epsilon)\varphi^m + \Delta^2\varphi^m}_{\delta_\varphi E^+(\varphi^m)} + \underbrace{2\Delta\varphi^{m-1}}_{\delta_\varphi E^-(\varphi^{m-1})}\end{aligned}$$

$$\text{where } E(\varphi) = \int_{\Omega} \left( \frac{\varphi^4}{4} + \frac{1-\epsilon}{2}\varphi^2 + \frac{1}{2}(\Delta\varphi)^2 - |\nabla\varphi|^2 \right) dx.$$

# Convex Time Splitting Scheme

- Basic Idea:

$$\begin{aligned}\mu &= \delta_\varphi E(\varphi) = \delta_\varphi \left( \underbrace{E^+(\varphi)}_{\text{convex}} + \underbrace{E^-(\varphi)}_{\text{concave}} \right) \\ \implies \mu^m &= \underbrace{(\varphi^m)^3 + (1 - \epsilon)\varphi^m + \Delta^2\varphi^m}_{\delta_\varphi E^+(\varphi^m)} + \underbrace{2\Delta\varphi^{m-1}}_{\delta_\varphi E^-(\varphi^{m-1})}\end{aligned}$$

where  $E(\varphi) = \int_{\Omega} \left( \frac{\varphi^4}{4} + \frac{1-\epsilon}{2}\varphi^2 + \frac{1}{2}(\Delta\varphi)^2 - |\nabla\varphi|^2 \right) dx.$

Given  $\varphi^0$ , find  $(\varphi^m, \mu^m)$  for  $1 \leq m \leq M$  by

$$\begin{aligned}\delta_\tau \varphi^m - \nabla \cdot (\mathcal{M} \nabla \mu^m) &= 0, \\ (\varphi^m)^3 + (1 - \epsilon)\varphi^m + \Delta^2\varphi^m + 2\Delta\varphi^{m-1} - \mu^m &= 0,\end{aligned}$$

with boundary conditions

$$\partial_n \varphi^m = \partial_n \Delta \varphi^m = \partial_n \mu^m = 0.$$

# Fully Discrete $C^0$ -IP Method

Given  $\varphi_h^{m-1} \in Z_h$ , find  $\varphi_h^m, \mu_h^m \in Z_h \times V_h$  such that for all  $\nu_h \in V_h, \psi_h \in Z_h$  it holds

$$(\delta_\tau \varphi_h^m, \nu_h) + (\mathcal{M} \nabla \mu_h^m, \nabla \nu_h) = 0$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi_h \right) + a_h^{IP}(\varphi_h^m, \psi_h) - 2(\nabla \varphi_h^{m-1}, \nabla \psi_h) - (\mu_h^m, \psi_h) = 0,$$

where  $\varphi_h^0 := P_h \varphi_0$  and  $\mu_h^0 \in V_h$  is defined as  $\mu_h^0 := R_h \mu_0$ .

## Remark

*The scheme satisfies the discrete conservation property*

$$(\varphi_h^m, 1) = (\varphi_h^0, 1) = (\varphi_0, 1) \text{ for any } 1 \leq m \leq M.$$

# Properties of Scheme

- 1 Unique Solvability:

# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.



# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.
- 2 Unconditional Stability:

# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.
- 2 Unconditional Stability:

$$F(\varphi_h^m) := \frac{1}{4} \|\varphi_h^m\|_{L^4}^4 + \frac{1-\epsilon}{2} \|\varphi_h^m\|_{L^2}^2 - \|\nabla \varphi_h^m\|_{L^2}^2 + \frac{1}{2} a_h^{IP}(\varphi_h^m, \varphi_h^m)$$

# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.
- 2 Unconditional Stability:

$$F(\varphi_h^m) := \frac{1}{4} \|\varphi_h^m\|_{L^4}^4 + \frac{1-\epsilon}{2} \|\varphi_h^m\|_{L^2}^2 - \|\nabla \varphi_h^m\|_{L^2}^2 + \frac{1}{2} a_h^{IP}(\varphi_h^m, \varphi_h^m)$$

Direct consequence of the convex decomposition of the energy.

# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.

- 2 Unconditional Stability:

$$F(\varphi_h^m) := \frac{1}{4} \|\varphi_h^m\|_{L^4}^4 + \frac{1-\epsilon}{2} \|\varphi_h^m\|_{L^2}^2 - \|\nabla \varphi_h^m\|_{L^2}^2 + \frac{1}{2} a_h^{IP}(\varphi_h^m, \varphi_h^m)$$

Direct consequence of the convex decomposition of the energy.

- 3 Optimal error estimates:

# Properties of Scheme

- 1 Unique Solvability:  
Relies on Convexity arguments.

- 2 Unconditional Stability:

$$F(\varphi_h^m) := \frac{1}{4} \|\varphi_h^m\|_{L^4}^4 + \frac{1-\epsilon}{2} \|\varphi_h^m\|_{L^2}^2 - \|\nabla \varphi_h^m\|_{L^2}^2 + \frac{1}{2} a_h^{IP}(\varphi_h^m, \varphi_h^m)$$

Direct consequence of the convex decomposition of the energy.

- 3 Optimal error estimates: **Main Result!**

# Unconditional Unique Solvability

## Definition

Define the functional  $G_h : \mathring{Z}_h \rightarrow \mathbb{R}$

$$G_h(\varphi_h) := \frac{\tau}{2} \left\| \frac{\varphi_h - \varphi_h^{m-1}}{\tau} \right\|_{-1,h}^2 + \frac{1}{2} a_h^{IP}(\varphi_h, \varphi_h) + \frac{1}{4} \|\varphi_h + \bar{\varphi}_0\|_{L^4(\Omega)}^4 \\ + \frac{1-\epsilon}{2} \|\varphi_h + \bar{\varphi}_0\|_{L^2(\Omega)}^2 - 2(\nabla \varphi_h^{m-1}, \nabla \varphi_h),$$

where

$$\|v_h\|_{-1,h} = (\nabla T_h v_h, \nabla T_h v_h)^{1/2} = (v_h, T_h v_h)^{1/2} = (T_h v_h, v_h)^{1/2},$$

with  $T_h : \mathring{Z}_h \rightarrow \mathring{Z}_h$  defined as: given  $\zeta_h \in \mathring{Z}_h$ , find  $T_h \zeta_h \in \mathring{Z}_h$ :

$$(\nabla T_h \zeta_h, \nabla \chi_h) = (\zeta_h, \chi_h) \quad \forall \chi_h \in \mathring{Z}_h.$$

# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

## Proof.



# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

## Proof.





# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

## Proof.

- 1 Shown through the **zero mean formulation** corresponding to our scheme



# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

## Proof.

- 1 Shown through the **zero mean formulation** corresponding to our scheme
- 2 Prove the one-to-one correspondence between solution of our scheme with the solution to the **zero mean formulation**



# Unconditional Unique Solvability

## Theorem

*The fully discrete  $C^0$ -IP scheme is uniquely solvable for any mesh parameters:  $\tau$  and  $h$  and for any  $\varepsilon < 1$ .*

## Proof.

- 1 Shown through the **zero mean formulation** corresponding to our scheme
- 2 Prove the one-to-one correspondence between solution of our scheme with the solution to the **zero mean formulation**
- 3 Existence of the unique solution to the **zero mean formulation** proved through existence of a minimizer for a functional  $G_h$ .



# Unconditional Stability

## Lemma (Discrete Energy Law)

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be a solution of the  $C^0$ -IP method. Then the following energy law holds for any  $h, \tau > 0$ :

$$\begin{aligned} & F(\varphi_h^\ell) + \tau \sum_{m=1}^{\ell} \left\| \mathcal{M}^{1/2} \nabla \mu_h^m \right\|_{L^2(\Omega)}^2 \\ & + \tau^2 \sum_{m=1}^{\ell} \left\{ \frac{(1-\epsilon)}{2} \|\delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 + \|\nabla \delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 \right. \\ & \quad \left. + \frac{1}{4} \|\delta_\tau (\varphi_h^m)^2\|_{L^2(\Omega)}^2 + \frac{1}{2} \|\varphi_h^m \delta_\tau \varphi_h^m\|_{L^2(\Omega)}^2 + \frac{1}{2} a_h^{IP}(\delta_\tau \varphi_h^m, \delta_\tau \varphi_h^m) \right\} \\ & = F(\varphi_h^0), \quad 1 \leq \ell \leq M. \end{aligned}$$

# Uniform *a priori* estimates

## Lemma

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be the unique solution of  $C^0$ -IP scheme. Suppose that  $F(\varphi_h^0) \leq C$  independent of  $h$  and  $\epsilon < \frac{C_{coer}-4}{C_{coer}} < 1$ . For any  $h, \tau > 0$ :

$$\max_{0 \leq m \leq M} \left[ \|\varphi_h^m\|_{L^4(\Omega)}^2 + \|\varphi_h^m\|_{L^2(\Omega)}^2 + \|\varphi_h^m\|_{2,h}^2 \right] \leq C$$

$$\max_{0 \leq m \leq M} \|\varphi_h^m\|_{H^1}^2 \leq C$$

$$\tau \sum_{m=1}^{\ell} \left\| \mathcal{M}^{1/2} \nabla \mu_h^m \right\|_{L^2(\Omega)}^2 \leq C$$

$$\tau^2 \sum_{m=1}^{\ell} \left\{ \|\nabla \delta_{\tau} \varphi_h^m\|_{L^2(\Omega)}^2 + \|(\varphi_h^m)^2 \delta_{\tau} (\varphi_h^m)^2\|_{L^2(\Omega)}^2 + \|\delta_{\tau} \varphi_h^m\|_{2,h}^2 \right\} \leq C$$

for some constant  $C$  that is independent of  $h, \tau$ , and  $T$ .

# Error Estimates

Assume additional regularites:

$$\begin{aligned}\varphi &\in L^\infty(0, T; H^3(\Omega)) \cap L^2(0, T; H^3(\Omega)), \\ \partial_t \varphi &\in L^2(0, T; H^3(\Omega)) \cap L^2(0, T; H_N^{-1}(\Omega)), \\ \partial_{tt} \varphi &\in L^2(0, T; L^2(\Omega)), \\ \mu &\in L^2(0, T; H^2(\Omega)), \\ \partial_t \mu &\in L^2(0, T; L^2(\Omega)).\end{aligned}$$

$C^0$ -IP Norm:

$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [\mathbf{n}_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

Notation:

$$e^{\varphi,m} := \varphi^m - \varphi_h^m, \quad e^{\mu,m} := \mu^m - \mu_h^m.$$

Assumption:  $\mathcal{M} \equiv 1$ .

# Error Analysis: Error Equation

## Weak Form:

$$\langle \partial_t \varphi^m, \nu \rangle + (\nabla \mu^m, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega) \quad (1)$$

$$\left( (\varphi^m)^3 + (1 - \epsilon) \varphi^m, \psi \right) - 2(\nabla \varphi^m, \nabla \psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z. \quad (2)$$

## Fully Discrete $C^0$ -IP Form:

$$(\delta_\tau \varphi_h^m, \nu) + (\nabla \mu_h^m, \nabla \nu) = 0 \quad \forall \nu \in V_h \quad (3)$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi \right) - 2(\nabla \varphi_h^{m-1}, \nabla \psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h \quad (4)$$

# Error Analysis: Error Equation

## Weak Form:

$$\langle \partial_t \varphi^m, \nu \rangle + (\nabla \mu^m, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega) \quad (1)$$

$$\left( (\varphi^m)^3 + (1 - \epsilon) \varphi^m, \psi \right) - 2 (\nabla \varphi^m, \nabla \psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z. \quad (2)$$

## Fully Discrete $C^0$ -IP Form:

$$(\delta_\tau \varphi_h^m, \nu) + (\nabla \mu_h^m, \nabla \nu) = 0 \quad \forall \nu \in V_h \quad (3)$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi \right) - 2 (\nabla \varphi_h^{m-1}, \nabla \psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h \quad (4)$$

- (1) and (3)  $\implies (\delta_\tau e^{\varphi, m}, \nu_h) + (\nabla e^{\mu, m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h),$   
 $\nu_h \in V_h \subset H^1(\Omega).$



# Error Analysis: Error Equation

## Weak Form:

$$\langle \partial_t \varphi^m, \nu \rangle + (\nabla \mu^m, \nabla \nu) = 0 \quad \forall \nu \in H^1(\Omega) \quad (1)$$

$$\left( (\varphi^m)^3 + (1 - \epsilon) \varphi^m, \psi \right) - 2 (\nabla \varphi^m, \nabla \psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z. \quad (2)$$

## Fully Discrete $C^0$ -IP Form:

$$(\delta_\tau \varphi_h^m, \nu) + (\nabla \mu_h^m, \nabla \nu) = 0 \quad \forall \nu \in V_h \quad (3)$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi \right) - 2 (\nabla \varphi_h^{m-1}, \nabla \psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h \quad (4)$$

- (1) and (3)  $\implies (\delta_\tau e^{\varphi, m}, \nu_h) + (\nabla e^{\mu, m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h), \nu_h \in V_h \subset H^1(\Omega).$
- Error equation based on (2) and (4) is not well-defined since  $Z_h \not\subset Z!$

# Error Analysis: Error Equation

**Weak Form:**

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla\varphi^m, \nabla\psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z.$$

**Fully Discrete  $C^0$ -IP Form:**

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla\varphi_h^{m-1}, \nabla\psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h$$

# Error Analysis: Error Equation

**Weak Form:**

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla\varphi^m, \nabla\psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z.$$

**Fully Discrete  $C^0$ -IP Form:**

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla\varphi_h^{m-1}, \nabla\psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h$$

$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3 + (1 - \epsilon)e^{\varphi, m}, \psi \right) - 2(\nabla e^{\varphi, m}, \nabla\psi) + a(\varphi^m, \psi) - a_h^{IP}(\varphi_h^m, \psi) \\ & - (e^{\mu, m}, \psi) = 0 \end{aligned}$$

# Error Analysis: Error Equation

## Weak Form:

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla\varphi^m, \nabla\psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z.$$

## Fully Discrete $C^0$ -IP Form:

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla\varphi_h^{m-1}, \nabla\psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h$$

$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3 + (1 - \epsilon)e^{\varphi, m}, \psi \right) - 2(\nabla e^{\varphi, m}, \nabla\psi) + a(\varphi^m, \psi) - a_h^{IP}(\varphi_h^m, \psi) \\ & - (e^{\mu, m}, \psi) = 0 \end{aligned}$$

Problem:  $\psi \in Z$  is not in  $H^{2+1/2}$  locally!  $\psi \in Z_h$  is not in  $H^2$  globally!

# Error Analysis: Error Equation

## Weak Form:

$$\left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi \right) - 2(\nabla\varphi^m, \nabla\psi) + a(\varphi^m, \psi) - (\mu^m, \psi) = 0 \quad \forall \psi \in Z.$$

## Fully Discrete $C^0$ -IP Form:

$$\left( (\varphi_h^m)^3 + (1 - \epsilon)\varphi_h^m, \psi \right) - 2(\nabla\varphi_h^{m-1}, \nabla\psi) + a_h^{IP}(\varphi_h^m, \psi) - (\mu_h^m, \psi) = 0, \quad \forall \psi \in Z_h$$

$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3 + (1 - \epsilon)e^{\varphi, m}, \psi \right) - 2(\nabla e^{\varphi, m}, \nabla\psi) + a(\varphi^m, \psi) - a_h^{IP}(\varphi_h^m, \psi) \\ & - (e^{\mu, m}, \psi) = 0 \end{aligned}$$

Problem:  $\psi \in Z$  is not in  $H^{2+1/2}$  locally!  $\psi \in Z_h$  is not in  $H^2$  globally!  
Remedy: Lift  $\psi \in Z_h$  into a finite dimensional subspace of  $Z$ .

# Error Analysis

- Remedy: Introduce  $W_h \subset Z$  to be the Hsieh-Clough-Tocher micro finite element space associated with  $\mathcal{T}_h$ .

# Error Analysis

- Remedy: Introduce  $W_h \subset Z$  to be the Hsieh-Clough-Tocher micro finite element space associated with  $\mathcal{T}_h$ .
- Define the enriching operator  $E_h : Z_h \rightarrow W_h \cap Z$  [Brenner, Gudi, Sung '12]

# Error Analysis

- Remedy: Introduce  $W_h \subset Z$  to be the Hsieh-Clough-Tocher micro finite element space associated with  $\mathcal{T}_h$ .
- Define the enriching operator  $E_h : Z_h \rightarrow W_h \cap Z$  [Brenner, Gudi, Sung '12]
- **Weak Form with correction term:** find  $(\varphi^m, \mu^m) \in Z \times H^1(\Omega)$ :

$$(\partial_t \varphi^m, \nu_h) + (\nabla \mu^m, \nabla \nu_h) = 0 \quad \forall \nu_h \in V_h,$$

$$\begin{aligned} a_h^{IP}(\varphi^m, \psi_h) + \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h \right) - 2(\nabla \varphi^m, \nabla \psi_h) - (\mu^m, \psi_h) \\ = \mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) \quad \forall \psi_h \in Z_h \end{aligned}$$

where

$$\begin{aligned} \mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) := a_h^{IP}(\varphi^m, \psi_h - E_h \psi_h) - (\mu^m, \psi_h - E_h \psi_h) \\ + \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h - E_h \psi_h \right) - 2(\nabla \varphi^m, \nabla \psi_h - \nabla E_h \psi_h). \end{aligned}$$



# Error Analysis

- Remedy: Introduce  $W_h \subset Z$  to be the Hsieh-Clough-Tocher micro finite element space associated with  $\mathcal{T}_h$ .
- Define the enriching operator  $E_h : Z_h \rightarrow W_h \cap Z$  [Brenner, Gudi, Sung '12]
- **Weak Form with correction term:** find  $(\varphi^m, \mu^m) \in Z \times H^1(\Omega)$ :

$$\begin{aligned}(\partial_t \varphi^m, \nu_h) + (\nabla \mu^m, \nabla \nu_h) &= 0 \quad \forall \nu_h \in V_h, \\ a_h^{IP}(\varphi^m, \psi_h) + \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h \right) - 2(\nabla \varphi^m, \nabla \psi_h) - (\mu^m, \psi_h) \\ &= \mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) \quad \forall \psi_h \in Z_h\end{aligned}$$

where

$$\begin{aligned}\mathcal{L}(\varphi^m, \mu^m; \psi_h - E_h \psi_h) &:= a_h^{IP}(\varphi^m, \psi_h - E_h \psi_h) - (\mu^m, \psi_h - E_h \psi_h) \\ &+ \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, \psi_h - E_h \psi_h \right) - 2(\nabla \varphi^m, \nabla \psi_h - \nabla E_h \psi_h).\end{aligned}$$

- Solutions to weak form are consistent since  $a_h^{IP}(\varphi, E_h \psi) = a(\varphi, E_h \psi)$  for all  $\psi \in Z_h$ .

# Error Equation

Subtracting fully discrete form from the weak form with the correction term gives:

$$(\delta_\tau e^{\varphi,m}, \nu_h) + (\nabla e^{\mu,m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h), \quad (5)$$

$$\begin{aligned} a_h^{IP}(e^{\varphi,m}, \psi_h) + ((1 - \epsilon)e^{\varphi,m}, \psi_h) - 2(\nabla e^{\varphi,m-1}, \nabla \psi_h) - (e^{\mu,m}, \psi_h) = \\ - \left( (\varphi^m)^3 - (\varphi_h^m)^3, \psi_h \right) - 2(\nabla \varphi^{m-1} - \nabla \varphi^m, \nabla \psi_h) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h). \end{aligned} \quad (6)$$

**Notation:**

$$\begin{aligned} e^{\varphi,m} &= e_P^{\varphi,m} + e_h^{\varphi,m}, & e_P^{\varphi,m} &:= \varphi^m - P_h \varphi^m, & e_h^{\varphi,m} &:= P_h \varphi^m - \varphi_h^m, \\ e^{\mu,m} &= e_R^{\mu,m} + e_h^{\mu,m}, & e_R^{\mu,m} &:= \mu^m - R_h \mu^m, & e_h^{\mu,m} &:= R_h \mu^m - \mu_h^m. \end{aligned}$$

# Error Equation

Subtracting fully discrete form from the weak form with the correction term gives:

$$(\delta_\tau e^{\varphi,m}, \nu_h) + (\nabla e^{\mu,m}, \nabla \nu_h) = (\delta_\tau \varphi^m - \partial_t \varphi^m, \nu_h), \quad (5)$$

$$\begin{aligned} a_h^{IP}(e^{\varphi,m}, \psi_h) + ((1 - \epsilon)e^{\varphi,m}, \psi_h) - 2(\nabla e^{\varphi,m-1}, \nabla \psi_h) - (e^{\mu,m}, \psi_h) = \\ - \left( (\varphi^m)^3 - (\varphi_h^m)^3, \psi_h \right) - 2(\nabla \varphi^{m-1} - \nabla \varphi^m, \nabla \psi_h) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h). \end{aligned} \quad (6)$$

**Notation:**

$$\begin{aligned} e^{\varphi,m} &= e_P^{\varphi,m} + e_h^{\varphi,m}, & e_P^{\varphi,m} &:= \varphi^m - P_h \varphi^m, & e_h^{\varphi,m} &:= P_h \varphi^m - \varphi_h^m, \\ e^{\mu,m} &= e_R^{\mu,m} + e_h^{\mu,m}, & e_R^{\mu,m} &:= \mu^m - R_h \mu^m, & e_h^{\mu,m} &:= R_h \mu^m - \mu_h^m. \end{aligned}$$

Set  $\nu_h = e_h^{\mu,m}$  in (5) and  $\psi_h = \delta_\tau e_h^{\varphi,m}$  in (6).

# Error Equation

$$\begin{aligned} & \|\nabla e_h^{\mu,m}\|_{L^2}^2 + a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) + ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) - 2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) \\ &= (\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) - (\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) + (e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla \varphi^m - \nabla \varphi^{m-1}, \nabla \delta_\tau e_h^{\varphi,m}) - ((\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla e_P^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h) \end{aligned}$$

## Lemma

Let  $(\varphi^m, \mu^m)$  be a weak solution with the additional regularities. Then for any  $h, \tau > 0$  and any  $0 \leq m \leq M$ , we have

$$\|\delta_\tau e_h^{\varphi, m}\|_{-1, h}^2 \leq 4 \|\nabla e_h^{\mu, m}\|_{L^2}^2 + \frac{Ch^2}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s)\|_{H^2}^2 ds + C\tau \int_{t_{m-1}}^{t_m} \|\partial_{ss} \varphi(s)\|_{H^1}^2 ds$$

where the constant  $C$  may depend upon a Poincaré constant but does not depend on  $h$  or  $\tau$ .

where  $\|v_h\|_{-1, h} = (\nabla T_h v_h, \nabla T_h v_h)^{1/2} = (v_h, T_h v_h)^{1/2} = (T_h v_h, v_h)^{1/2}$  and  $T_h$  is the discrete inverse Laplacian.

# Error Analysis: LHS

$$\begin{aligned} & \|\nabla e_h^{\mu,m}\|_{L^2}^2 + a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) + ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) - 2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) \\ &= (\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) - (\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) + (e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla \varphi^m - \nabla \varphi^{m-1}, \nabla \delta_\tau e_h^{\varphi,m}) - ((\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m}) \\ &+ 2(\nabla e_P^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) + \mathcal{L}(\varphi^m, \mu_h^m; \psi_h - E_h \psi_h) \end{aligned}$$

Polarization Property:

$$\begin{aligned} a_h^{IP}(e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) &= \frac{1}{2} \delta_\tau a_h^{IP}(e_h^{\varphi,m}, e_h^{\varphi,m}) + \frac{\tau}{2} a_h^{IP}(\delta_\tau e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) \\ ((1-\epsilon)e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) &= \frac{(1-\epsilon)}{2} \delta_\tau \|e_h^{\varphi,m}\|_{L^2}^2 + \frac{(1-\epsilon)}{2} \|\delta_\tau e_h^{\varphi,m}\|_{L^2}^2 \\ -2(\nabla e_h^{\varphi,m-1}, \nabla \delta_\tau e_h^{\varphi,m}) &= \tau \|\nabla \delta_\tau e_h^{\varphi,m}\|_{L^2}^2 - \delta_\tau (\nabla e_h^{\varphi,m}, \nabla e_h^{\varphi,m}) \end{aligned}$$

# Error Analysis: First 3 RHS terms

$$(\delta_\tau \varphi^m - \partial_t \varphi^m, e_h^{\mu,m}) \leq C\tau \int_{t_{m-1}}^{t_m} \|\partial_{ss}\varphi(s)\|_{L^2}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2,$$

$$\begin{aligned}(\delta_\tau e_P^{\varphi,m}, e_h^{\mu,m}) &\leq C \|\delta_\tau e_P^{\varphi,m}\|_{L^2}^2 + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 \\ &\leq \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|P_h \partial_s \varphi(s) - \partial_s \varphi(s)\|_{L^2}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 \\ &\leq \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds + \frac{1}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2,\end{aligned}$$

$$(e_R^{\mu,m}, \delta_\tau e_h^{\varphi,m}) \leq C \|\nabla e_R^{\mu,m}\|_{L^2}^2 + \frac{1}{36} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2,$$

## Error Analysis: next two RHS terms

$$\begin{aligned} 2(\nabla\varphi^m - \nabla\varphi^{m-1}, \nabla\delta_\tau e_h^{\varphi,m}) &= -2(\tau\Delta\delta_\tau\varphi^m, \delta_\tau e_h^{\varphi,m}) \\ &\leq 2\|\tau\nabla\Delta\delta_\tau\varphi^m\|_{L^2} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &\leq C_\tau \int_{t_{m-1}}^{t_m} \|\partial_s\varphi(s)\|_{H^3}^2 ds + \frac{1}{36} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2. \end{aligned}$$

$$\begin{aligned} \left( (\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m} \right) &\leq \left\| \nabla \left( (\varphi^m)^3 - (\varphi_h^m)^3 \right) \right\|_{L^2} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &= \left\| 3(\varphi^m)^2 \nabla\varphi^m - 3(\varphi_h^m)^2 \nabla\varphi_h^m \right\|_{L^2} \\ &\quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ &= 3 \left\| (\varphi^m + \varphi_h^m) \nabla\varphi^m e^{\varphi,m} + (\varphi_h^m)^2 \nabla e^{\varphi,m} \right\|_{L^2} \\ &\quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \end{aligned}$$



## Error Analysis: next two RHS terms

$$\begin{aligned} & \left( (\varphi^m)^3 - (\varphi_h^m)^3, \delta_\tau e_h^{\varphi,m} \right) \\ & \leq 3 \left( \|\varphi^m + \varphi_h^m\|_{L^6} \|\nabla \varphi^m\|_{L^6} \|e^{\varphi,m}\|_{L^6} + \|\varphi_h^m\|_{L^6}^2 \|\nabla e^{\varphi,m}\|_{L^6} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \left( \|\nabla e_P^{\varphi,m}\|_{L^2} + \|\nabla e_h^{\varphi,m}\|_{L^2} + \|e_P^{\varphi,m}\|_{2,h} + \|e_h^{\varphi,m}\|_{2,h} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \left( \|\nabla e_P^{\varphi,m}\|_{L^2} + \|\nabla e_h^{\varphi,m}\|_{L^2} + \|e_P^{\varphi,m}\|_{2,h} + \|e_h^{\varphi,m}\|_{2,h} \right) \\ & \quad \times \|\delta_\tau e_h^{\varphi,m}\|_{-1,h} \\ & \leq C \|e_P^{\varphi,m}\|_{2,h}^2 + C \|e_h^{\varphi,m}\|_{2,h}^2 + \frac{1}{36} \|\delta_\tau e_h^{\varphi,m}\|_{-1,h}^2. \end{aligned}$$

## Error Analysis: second last term

Using discrete product rule:

$$\begin{aligned}\left(a^{m-1}, \frac{b^m - b^{m-1}}{\tau}\right) &= \frac{1}{\tau} [(a^m, b^m) - (a^{m-1}, b^{m-1})] - \left(\frac{a^m - a^{m-1}}{\tau}, b^m\right) \\ &= \delta_\tau (a^m, b^m) - (\delta_\tau a^m, b^m),\end{aligned}$$

we have the following bound

$$\begin{aligned}2 \left(\nabla e_P^{\varphi, m-1}, \nabla \delta_\tau e_h^{\varphi, m}\right) &= 2\delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) - 2(\nabla \delta_\tau e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) \\ &\leq 2\delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) + C \|\delta_\tau e_P^{\varphi, m}\|_{L^2}^2 + C \|e_h^{\varphi, m}\|_{2,h}^2 \\ &\leq 2\delta_\tau (\nabla e_P^{\varphi, m}, \nabla e_h^{\varphi, m}) + \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds \\ &\quad + C \|e_h^{\varphi, m}\|_{2,h}^2.\end{aligned}$$

## Lemma

Suppose  $(\varphi^m, \mu^m)$  is a weak solution to the PFC equation, with the additional regularities. Then for any  $h, \tau > 0$  and any  $0 \leq m \leq M$  and any  $\beta > 0$ ,

$$\begin{aligned} a_h^{IP}(\varphi^m, e_h^{\varphi, m} - E_h e_h^{\varphi, m}) + \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m, e_h^{\varphi, m} - E_h e_h^{\varphi, m} \right) \\ - 2(\nabla \varphi^m, \nabla (e_h^{\varphi, m} - E_h e_h^{\varphi, m})) - (\mu^m, e_h^{\varphi, m} - E_h e_h^{\varphi, m}) \leq C [\text{Osc}_j(\mu^m)]^2 + \\ C \|e_P^{\varphi, m}\|_{2, h}^2 + \frac{C_{\text{coer}}}{4\beta} \|e_h^{\varphi, m}\|_{2, h}^2 \end{aligned}$$

and

$$\begin{aligned} a_h^{IP}(\delta_\tau \varphi^m, e_h^{\varphi, m-1} - E_h e_h^{\varphi, m-1}) + \left( \delta_\tau \left( (\varphi^m)^3 + (1 - \epsilon)\varphi^m \right), e_h^{\varphi, m-1} - E_h e_h^{\varphi, m-1} \right) \\ - 2\left( \delta_\tau \nabla \varphi^m, \nabla (e_h^{\varphi, m-1} - E_h e_h^{\varphi, m-1}) \right) - \left( \delta_\tau \mu^m, e_h^{\varphi, m-1} - E_h e_h^{\varphi, m-1} \right) \\ \leq C [\text{Osc}_j(\mu_t(t^*))]^2 + C \|e_P^{\varphi, m}\|_{2, h}^2 \end{aligned}$$

$$\begin{aligned}
& \|\nabla e_h^{\mu,m}\|_{L^2}^2 + \frac{1}{2} \delta_\tau a_h^{IP} (e_h^{\varphi,m}, e_h^{\varphi,m}) + \frac{\tau}{2} a_h^{IP} (\delta_\tau e_h^{\varphi,m}, \delta_\tau e_h^{\varphi,m}) \\
& + \frac{(1-\epsilon)}{2} \delta_\tau \|e_h^{\varphi,m}\|_{L^2}^2 + \frac{(1-\epsilon)\tau}{2} \|\delta_\tau e_h^{\varphi,m}\|_{L^2}^2 + \tau \|\nabla \delta_\tau e_h^{\varphi,m}\|_{L^2}^2 \\
& \leq \delta_\tau (\nabla e_h^{\varphi,m}, \nabla e_h^{\varphi,m}) + 2\delta_\tau (\nabla e_P^{\varphi,m}, \nabla e_h^{\varphi,m}) + \frac{6}{12} \|\nabla e_h^{\mu,m}\|_{L^2}^2 + C \|e_h^{\varphi,m}\|_{2,h}^2 \\
& + C \|e_h^{\varphi,m-1}\|_{2,h}^2 + C \|\nabla e_R^{\mu,m}\|_{L^2}^2 + C \|e_P^{\varphi,m}\|_{2,h}^2 + C [\text{Osc}_j(\mu_t(t^*))]^2 \\
& + C\tau \int_{t_{m-1}}^{t_m} \left[ \|\partial_s \varphi(s)\|_{H^3}^2 + \|\partial_{ss} \varphi(s)\|_{L^2}^2 \right] ds + \frac{C}{\tau} \int_{t_{m-1}}^{t_m} \|\partial_s \varphi(s) - P_h \partial_s \varphi(s)\|_{2,h}^2 ds \\
& + \delta_\tau a_h^{IP} (\varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}) + \delta_\tau \left( (\varphi^m)^3 + (1-\epsilon)\varphi^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m} \right) \\
& - 2\delta_\tau (\nabla \varphi^m, \nabla (e_h^{\varphi,m} - E_h e_h^{\varphi,m})) - \delta_\tau (\mu^m, e_h^{\varphi,m} - E_h e_h^{\varphi,m}).
\end{aligned}$$

Applying  $2\tau \sum_{m=1}^{\ell}$ , using the fact that  $e_h^{\varphi,0} = 0$  we obtain

$$\begin{aligned}
& a_h^{IP} \left( e_h^{\varphi,\ell}, e_h^{\varphi,\ell} \right) + (1 - \epsilon) \left\| e_h^{\varphi,\ell} \right\|_{L^2}^2 + \tau \sum_{m=1}^{\ell} \left\| \nabla e_h^{\mu,m} \right\|_{L^2}^2 \\
& + \tau^2 \sum_{m=1}^{\ell} \left[ a_h^{IP} \left( \delta_{\tau} e_h^{\varphi,m}, \delta_{\tau} e_h^{\varphi,m} \right) + (1 - \epsilon) \left\| \delta_{\tau} e_h^{\varphi,m} \right\|_{L^2}^2 + 2 \left\| \nabla \delta_{\tau} e_h^{\varphi,m} \right\|_{L^2}^2 \right] \\
& \leq \frac{C_{coer}}{2\beta} \left\| e_h^{\varphi,\ell} \right\|_{2,h}^2 + \frac{8\beta}{C_{coer}} \left\| e_h^{\varphi,\ell} \right\|_{L^2}^2 + C \left\| e_P^{\varphi,\ell} \right\|_{2,h}^2 + C\tau \sum_{m=1}^{\ell} \left\| e_h^{\varphi,m} \right\|_{2,h}^2 \\
& + C\tau \sum_{m=1}^{\ell} \left[ \left\| \nabla e_R^{\mu,m} \right\|_{L^2}^2 + \left\| e_P^{\varphi,m} \right\|_{2,h}^2 + [\text{Osc}_j(\mu_t(t^*))]^2 \right] \\
& + C\tau^2 \int_{t_0}^{t_{\ell}} \left[ \left\| \partial_s \varphi(s) \right\|_{H^3}^2 + \left\| \partial_{ss} \varphi(s) \right\|_{L^2}^2 \right] ds + C \int_{t_0}^{t_{\ell}} \left\| \partial_s \varphi(s) - P_h \partial_s \varphi(s) \right\|_{2,h}^2 ds \\
& + 2 \left[ C [\text{Osc}_j(\mu^{\ell})]^2 + C \left\| e_P^{\varphi,\ell} \right\|_{2,h}^2 + \frac{C_{coer}}{4\beta} \left\| e_h^{\varphi,\ell} \right\|_{2,h}^2 \right],
\end{aligned}$$

# Error Estimates: Main Result

## Theorem

Suppose  $(\varphi^m, \mu^m)$  is a weak solution to the weak form of the PFC equation, with the additional regularities. Then for any  $\tau, h > 0, \epsilon < \frac{C_{\text{coer}} - 16}{C_{\text{coer}}} < 1$  and any  $0 \leq \tau \leq M$ ,

$$\begin{aligned} & \left\| e_h^{\varphi, \ell} \right\|_{2, h}^2 + C \left\| e_h^{\varphi, \ell} \right\|_{L^2}^2 + C\tau \sum_{m=1}^{\ell} \left\| \nabla e_h^{\mu, m} \right\|_{L^2}^2 + \\ & C\tau^2 \sum_{m=1}^{\ell} \left[ \left\| \delta_{\tau} e_h^{\varphi, \ell} \right\|_{2, h}^2 + (1 - \epsilon) \left\| \delta_{\tau} e_h^{\varphi, m} \right\|_{L^2}^2 + \left\| \nabla \delta_{\tau} e_h^{\varphi, m} \right\|_{L^2}^2 \right] \leq C^*(h^2 + \tau^2) \end{aligned}$$

where  $C^*$  may depend on the oscillations of  $\mu$  and  $\partial_t \mu$  and the final stopping time  $T$  but does not depend upon the spacial step size  $h$  or the time step size  $\tau$ .

# Our Numerical Scheme

# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$



# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$
- Given  $\varphi_h^{m-1} \in Z_h$ , solve for  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$ : Given  $\varphi_h^{m-1} \in Z_h$ , find  $\varphi_h^m, \mu_h^m \in Z_h \times V_h$  such that for all  $\nu_h \in V_h, \psi_h \in Z_h$  it holds

$$(\delta_\tau \varphi_h^m, \nu_h) + (\mathcal{M} \nabla \mu_h^m, \nabla \nu_h) = 0$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi_h \right) + a_h^{IP}(\varphi_h^m, \psi_h) - 2(\nabla \varphi_h^{m-1}, \nabla \psi_h) - (\mu_h^m, \psi_h) = 0,$$

where  $\varphi_h^0 := P_h \varphi_0$  and  $\mu_h^0 \in V_h$  is defined as  $\mu_h^0 := R_h \mu_0$ .

model parameters:  $\epsilon = 0.025$  and  $\mathcal{M} = 1$ , penalty parameter:  $\alpha = 20$ .

# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$
- Given  $\varphi_h^{m-1} \in Z_h$ , solve for  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$ : Given  $\varphi_h^{m-1} \in Z_h$ , find  $\varphi_h^m, \mu_h^m \in Z_h \times V_h$  such that for all  $\nu_h \in V_h, \psi_h \in Z_h$  it holds

$$(\delta_\tau \varphi_h^m, \nu_h) + (\mathcal{M} \nabla \mu_h^m, \nabla \nu_h) = 0$$

$$\left( (\varphi_h^m)^3 + (1 - \epsilon) \varphi_h^m, \psi_h \right) + a_h^{IP}(\varphi_h^m, \psi_h) - 2(\nabla \varphi_h^{m-1}, \nabla \psi_h) - (\mu_h^m, \psi_h) = 0,$$

where  $\varphi_h^0 := P_h \varphi_0$  and  $\mu_h^0 \in V_h$  is defined as  $\mu_h^0 := R_h \mu_0$ .

model parameters:  $\epsilon = 0.025$  and  $\mathcal{M} = 1$ , penalty parameter:  $\alpha = 20$ .

- We use Newton iteration method to solve the discrete nonlinear system with tolerance of the Newton iteration is set to  $10^{-6}$ .  
The initial guess at each time step is taken as the numerical solution at the previous time level.  
One to three Newton's iterative steps are involved at each time step.

# Numerical Experiment I: Accuracy Test

Example (Hu, Wise, Wang, Lowengrub, 2009 )

$$\begin{aligned}\varphi_0(x, y) = & 0.07 - 0.02 \cos\left(\frac{2\pi(x-12)}{32}\right) \sin\left(\frac{2\pi(y-1)}{32}\right) \\ & + 0.02 \cos^2\left(\frac{\pi(x+10)}{32}\right) \cos^2\left(\frac{\pi(y+3)}{32}\right) \\ & - 0.01 \sin^2\left(\frac{4\pi x}{32}\right) \sin^2\left(\frac{4\pi(y-6)}{32}\right)\end{aligned}$$

$\Omega = (0, 32) \times (0, 32)$ ,  $T = 10$ .

$\mathcal{M} \equiv 1$ ,  $\varepsilon = 0.025$ , and the penalty parameter  $\alpha = 20$ .

# Numerical Experiment I: Accuracy Test

$$\|\xi_h\|_{2,h}^2 := \sum_{K \in \mathcal{T}_h} |\xi_h|_{H^2(K)}^2 + \sum_{e \in \mathcal{E}_h} \alpha \|h_e^{-\frac{1}{2}} [n_e \cdot \nabla \xi_h]_e\|_{L^2(e)}^2.$$

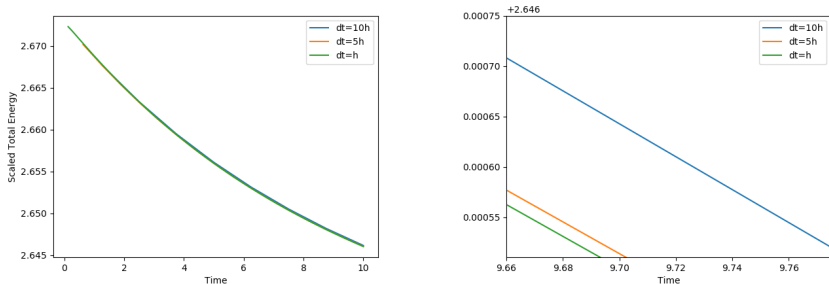
Mesh  $h = 32/512$  with  $\tau$  with  $\tau = 0.05h$  and  $T = 10$  as the 'exact' solution,  $\varphi_{exact}$ .

$$error_\varphi := \varphi_h - \varphi_{exact}$$

where  $\varphi_h$  indicates the solution on the mesh size  $h$ .

$h$	$\ error_\varphi\ _{2,h}$	rate	$\ error_\mu\ _{H^1}$	rate
32/8	0.08412	N/A	0.00522	N/A
32/16	0.05896	0.71329	0.00242	1.07627
32/32	0.03466	0.85058	0.00157	0.76970
32/64	0.01568	1.10514	0.00103	0.76082
32/128	0.00601	1.30482	0.00041	1.25840
32/256	0.00255	1.17707	0.00016	1.27362

# Numerical Experiment I: Unconditional Stability



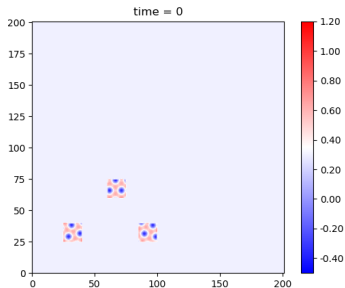
**Figure:** Unconditional stability demonstrated through the time evolution of the scaled total energy  $F/32^2$  for time step sizes  $dt = 10h$ ,  $5h$ ,  $h$  with the spacial step size  $h = 32/256$ .

# Numerical Experiment II: Crystal growth

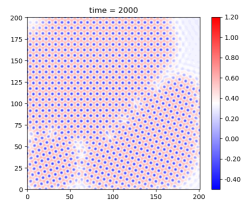
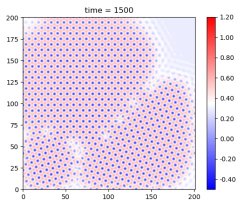
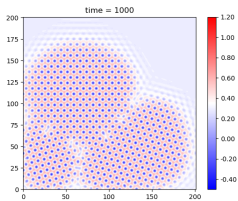
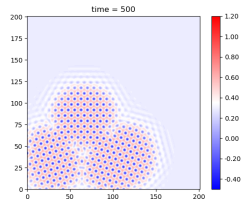
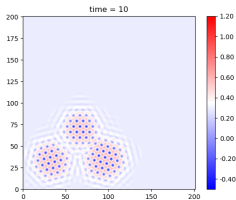
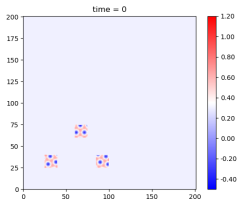
## Example (Gomez, Nogueira, 2012)

$$\varphi_0(x, y) = \bar{\varphi} + C \left[ \cos\left(\frac{q}{\sqrt{3}}y\right) \cos(qx) - 0.5 \cos\left(\frac{2q}{\sqrt{3}}y\right) \right]$$

where  $\bar{\varphi} = 0.285$ ,  $C = 0.466$ ,  $q = 0.66$ ,  $\Omega = (0, 201) \times (0, 201)$ ,  $h = 201/402$ .

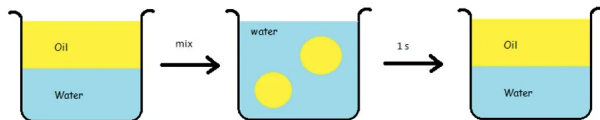


# Motion of liquid-crystal interfaces and grain boundaries



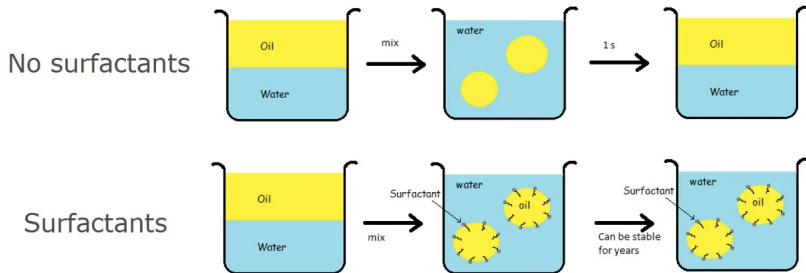
# Phase Field Model for Microemulsions: Motivation

No surfactants





# Phase Field Model for Microemulsions: Motivation



[Picture Courtesy: Biolin Scientific]

# Motivation: Microemulsion systems as skin drug delivery systems



[Source: Walgreens]

# Motivation: Microemulsion systems as skin drug delivery systems



- **Issue:** Topical cream formulations contain petrochemical ingredients (such as petrolatum, silicones).

[Source: Walgreens]

# Motivation: Microemulsion systems as skin drug delivery systems



- **Issue:** Topical cream formulations contain petrochemical ingredients (such as petrolatum, silicones).
- **Goal:** Develop skin drug delivery systems that contain natural and renewable sourced alternatives to these ingredients without compromising on the functionality.

[Source: Walgreens]

# Motivation: Microemulsion systems as skin drug delivery systems



[Source: Walgreens]

- **Issue:** Topical cream formulations contain petrochemical ingredients (such as petrolatum, silicones).
- **Goal:** Develop skin drug delivery systems that contain natural and renewable sourced alternatives to these ingredients without compromising on the functionality.
- **Approach:** Provide computational tools to predict and assess the properties of novel and more sustainable alternatives to the toxic ingredients.

[Source: The Nabi Laboratory of Bioengineered Therapeutics, UTEP]

# Mathematical Model for Microemulsions

- $\Omega \subset \mathbb{R}^2$  be a bounded polygonal domain occupied by the ternary mixture.

# Mathematical Model for Microemulsions

- $\Omega \subset \mathbb{R}^2$  be a bounded polygonal domain occupied by the ternary mixture.
- $\varphi$  : the scalar order parameter indicating the local difference between oil and water concentrations.

# Mathematical Model for Microemulsions

- $\Omega \subset \mathbb{R}^2$  be a bounded polygonal domain occupied by the ternary mixture.
- $\varphi$  : the scalar order parameter indicating the local difference between oil and water concentrations.
- $\varphi = -1$  (water phase),  $\varphi = 1$  (oil phase) and  $\varphi = 0$  (microemulsions phase)



# Mathematical Model for Microemulsions

- $\Omega \subset \mathbb{R}^2$  be a bounded polygonal domain occupied by the ternary mixture.
- $\varphi$  : the scalar order parameter indicating the local difference between oil and water concentrations.
- $\varphi = -1$  (water phase),  $\varphi = 1$  (oil phase) and  $\varphi = 0$  (microemulsions phase)
- $E(\varphi)$  : (Ginzburg-Landau free energy)

$$\underbrace{\int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla \varphi|^2 + \frac{\lambda}{2} (\Delta \varphi)^2 \right\} dx}_{\text{tendency to mix}} + \underbrace{\frac{\beta}{2} \int_{\Omega} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 dx}_{\text{tendency to separate}}$$

[Gompper et. al. 90]

# Mathematical Model for Microemulsions

Conservation Law:

$$\partial_t \varphi + \nabla \cdot \mathbf{j} = 0$$

- $\varphi$  : the scalar order parameter
- $\mathbf{j} = -\mathcal{M}\nabla\mu$ : mass flux
- $\mathcal{M}$ : mobility coefficient,  $\mu = \delta_\varphi E$ : chemical potential

$$E(\varphi) = \int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla\varphi|^2 + \frac{\lambda}{2} (\Delta\varphi)^2 + \frac{\beta}{2} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 \right\} dx$$

# Mathematical Model for Microemulsions

Conservation Law:

$$\partial_t \varphi + \nabla \cdot \mathbf{j} = 0$$

- $\varphi$  : the scalar order parameter
- $\mathbf{j} = -\mathcal{M}\nabla\mu$ : mass flux
- $\mathcal{M}$ : mobility coefficient,  $\mu = \delta_\varphi E$ : chemical potential

$$E(\varphi) = \int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla\varphi|^2 + \frac{\lambda}{2} (\Delta\varphi)^2 + \frac{\beta}{2} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 \right\} dx$$

$$\mu = 3\beta(\varphi^5 - \varphi^3) + \varphi|\nabla\varphi|^2 - \nabla \cdot (\varphi^2\nabla\varphi) + a_0\Delta\varphi + \lambda\Delta^2\varphi.$$

# Mathematical Model for Microemulsions

Conservation Law:

$$\partial_t \varphi + \nabla \cdot \mathbf{j} = 0$$

- $\varphi$  : the scalar order parameter
- $\mathbf{j} = -\mathcal{M}\nabla\mu$ : mass flux
- $\mathcal{M}$ : mobility coefficient,  $\mu = \delta_\varphi E$ : chemical potential

$$E(\varphi) = \int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla\varphi|^2 + \frac{\lambda}{2} (\Delta\varphi)^2 + \frac{\beta}{2} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 \right\} dx$$

$$\mu = 3\beta(\varphi^5 - \varphi^3) + \varphi|\nabla\varphi|^2 - \nabla \cdot (\varphi^2\nabla\varphi) + a_0\Delta\varphi + \lambda\Delta^2\varphi.$$

# Mathematical Model for Microemulsions

Conservation Law:

$$\partial_t \varphi + \nabla \cdot \mathbf{j} = 0$$

- $\varphi$  : the scalar order parameter
- $\mathbf{j} = -\mathcal{M}\nabla\mu$ : mass flux
- $\mathcal{M}$ : mobility coefficient,  $\mu = \delta_\varphi E$ : chemical potential

$$E(\varphi) = \int_{\Omega} \left\{ \frac{(\varphi^2 - a_0)}{2} |\nabla\varphi|^2 + \frac{\lambda}{2} (\Delta\varphi)^2 + \frac{\beta}{2} (\varphi + 1)^2 (\varphi^2 + 0.5) (\varphi - 1)^2 \right\} dx$$

$$\mu = 3\beta(\varphi^5 - \varphi^3) + \varphi|\nabla\varphi|^2 - \nabla \cdot (\varphi^2\nabla\varphi) + a_0\Delta\varphi + \lambda\Delta^2\varphi.$$

$$\partial_t \varphi - \mathcal{M}\Delta \left( 3\beta(\varphi^5 - \varphi^3) + \varphi|\nabla\varphi|^2 - \nabla \cdot ((\varphi^2 - 4)\nabla\varphi) + \lambda\Delta^2\varphi \right) = 0.$$

# Mathematical Model for Microemulsions

Compliment

$$\begin{aligned}\frac{\partial \varphi}{\partial t} - \nabla \cdot (\mathcal{M} \nabla \mu) &= 0, & \text{in } \Omega^T := \Omega \times (0, T), \\ 3\beta(\varphi^5 - \varphi^3) + \varphi |\nabla \varphi|^2 - \nabla \cdot ((\varphi^2 - a_0) \nabla \varphi) + \lambda \Delta^2 \varphi - \mu &= 0, & \text{in } \Omega^T\end{aligned}$$

with natural boundary conditions

$$\partial_n \varphi = \lambda \partial_n \Delta \varphi = \partial_n \mu = 0, \quad \text{on } \partial \Omega^T$$

and the initial value:

$$\varphi(0) = \varphi_0.$$

Notation:

- $H^s(\Omega)$  denote the Sobolev spaces of order  $s \geq 1$ ,
- $Z := \{z \in H^2(\Omega) \mid \partial_n z = 0 \text{ on } \partial \Omega\}$ .

[Pawlow et al., 2011]

# Numerical Schemes: Existing Literature

- Hoppe/Linsemann 2019: Fully implicit backward Euler and  $C^0$ -IP Method quasi-optimal error estimates without any discrete energy law

# Numerical Schemes: Existing Literature

- Hoppe/Linseemann 2019: Fully implicit backward Euler and  $C^0$ -IP Method quasi-optimal error estimates without any discrete energy law
- Diegel/Sharma 2022: Closely related literature is the  $C^0$ -IP framework developed for the Phase Field Crystal Equation based on Eyre's convex splitting scheme.



# Fully Discrete Scheme

Given  $\varphi_h^{m-1} \in Z_h$ , find  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  which satisfies

$$(\delta_\tau \varphi_h^m, \nu_h) + (M \nabla \mu_h^m, \nabla \nu_h) = 0, \quad \forall \nu_h \in V_h$$

$$3\beta ((\varphi_h^m)^5 - (\varphi_h^{m-1})^3, \psi_h) + ((\varphi_h^m)^2 \nabla \varphi_h^m, \nabla \psi_h) + (\varphi_h^m |\nabla \varphi_h^{m-1}|^2, \psi_h) \\ - a_0 (\nabla \varphi_h^{m-1}, \nabla \psi_h) + \lambda a_h^{IP} (\varphi_h^m, \psi_h) - (\mu_h^m, \psi_h) = 0 \quad \forall \psi_h \in Z_h$$

with initial data taken to be  $\varphi_h^0 := P_h \varphi_0$ .

## Theorem

Let  $\lambda \geq \frac{3\beta|\overline{\varphi_0}|^4 C_{P,1}}{2C_{coer}}$ , where  $C_{P,1}$  depends upon a Poincarè constant but does not depend upon  $h$  or  $\tau$ . Then, there exists a solution  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  to the scheme.

# Unconditional Energy Stability

$$F(\varphi) := \frac{\beta}{2} \|\varphi\|_{L^6}^6 - \frac{3\beta}{4} \|\varphi\|_{L^4}^4 + \frac{\beta|\Omega|}{4} + \frac{1}{2} \|\varphi \nabla \varphi\|_{L^2}^2 - \frac{a_0}{2} \|\nabla \varphi\|_{L^2}^2 + \frac{\lambda}{2} a_h^{IP}(\varphi, \varphi).$$

## Theorem (Discrete Energy Law)

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be a solution. Then the following energy law holds for any  $h, \tau > 0$ :

$$F(\varphi_h^\ell) + \tau \sum_{m=1}^{\ell} \left\| \sqrt{\mathcal{M}} \nabla \mu_h^m \right\|_{L^2}^2 \leq F(\varphi_h^0),$$

for all  $1 \leq \ell \leq M$ .

# Uniform *A Priori* Estimates

## Theorem

Let  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$  be the  $C^0$  IP approximation. Suppose that  $F(\varphi_h^0) \leq C$  independent of  $h$  and that  $\lambda > \max \left\{ \frac{3\beta|\overline{\varphi_0}|^4 C_{P,1}}{2C_{coer}}, \frac{a_0 C_{P,2}}{C_{coer}} \right\} > 0$  where  $C_{P,1}, C_{P,2}$  are Poincarè constants and do not depend on  $h$  or  $\tau$ . Then the following estimates hold for any  $\tau, h > 0$ :

$$\begin{aligned} \max_{0 \leq m \leq M} \|\varphi_h^m\|_{2,h}^2 &\leq C \\ \max_{0 \leq m \leq M} \left[ \|\varphi_h^m\|_{L^2}^2 + \|\nabla \varphi_h^m\|_{L^2}^2 + \|\varphi_h^m \nabla \varphi_h^m\|_{L^2}^2 + \|\varphi_h^m\|_{L^\infty}^2 \right] &\leq C^* \\ \tau \sum_{m=1}^{\ell} \left\| \sqrt{\mathcal{M}} \nabla \mu_h^m \right\|_{L^2}^2 &\leq C \end{aligned}$$

for some constants  $C^*, C$  that is independent of  $h, \tau$ , and  $T$ .

## Theorem

Let  $\varphi_h^{m-1} \in Z_h$  be given and

$$\lambda > \max \left\{ \frac{3\beta|\overline{\varphi_0}|^4 C_{P,1}}{2C_{coer}}, \frac{a_0 C_{P,2}}{C_{coer}}, \frac{C^* C_{P,3}}{2C_{coer}} \right\} > 0,$$

where  $C^*$  is the constant from uniform a priori bounds and  $C_{P,1}, C_{P,2}, C_{P,3}$  are all Poincaré constants and do not depend on  $h$  or  $\tau$ .

The solution to the fully discrete scheme is unique for all  $h, \tau > 0$ .

# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$

# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$

# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$
- Given  $\varphi_h^{m-1} \in Z_h$ , solve for  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$ :

$$(\delta_\tau \varphi_h^m, \nu_h) + (M \nabla \mu_h^m, \nabla \nu_h) = 0, \quad \forall \nu_h \in V_h$$

$$3\beta ((\varphi_h^m)^5 - (\varphi_h^{m-1})^3, \psi_h) + ((\varphi_h^m)^2 \nabla \varphi_h^m, \nabla \psi_h) + (\varphi_h^m |\nabla \varphi_h^{m-1}|^2, \psi_h) \\ - a_0 (\nabla \varphi_h^{m-1}, \nabla \psi_h) + \lambda a_h^{IP} (\varphi_h^m, \psi_h) - (\mu_h^m, \psi_h) = 0 \quad \forall \psi_h \in Z_h,$$

model parameters:  $a_0 = 4$  and  $\beta = 5$ , penalty parameter:  $\alpha = 8$ .



# Our Numerical Scheme

- Project  $\varphi_h^0 := P_h \varphi_0$
- Given  $\varphi_h^{m-1} \in Z_h$ , solve for  $(\varphi_h^m, \mu_h^m) \in Z_h \times V_h$ :

$$(\delta_\tau \varphi_h^m, \nu_h) + (M \nabla \mu_h^m, \nabla \nu_h) = 0, \quad \forall \nu_h \in V_h$$

$$3\beta ((\varphi_h^m)^5 - (\varphi_h^{m-1})^3, \psi_h) + ((\varphi_h^m)^2 \nabla \varphi_h^m, \nabla \psi_h) + (\varphi_h^m |\nabla \varphi_h^{m-1}|^2, \psi_h) - a_0 (\nabla \varphi_h^{m-1}, \nabla \psi_h) + \lambda a_h^{IP} (\varphi_h^m, \psi_h) - (\mu_h^m, \psi_h) = 0 \quad \forall \psi_h \in Z_h,$$

model parameters:  $a_0 = 4$  and  $\beta = 5$ , penalty parameter:  $\alpha = 8$ .

- We use Newton iteration method to solve the discrete nonlinear system with tolerance of the Newton iteration is set to  $10^{-6}$ .

The initial guess at each time step is taken as the numerical solution at the previous time level.

One to three Newton's iterative steps are involved at each time step.

# Numerical Experiment I: Accuracy Test

## Example

$$\varphi_0(x, y) = 0.3 \cos(3x) + 0.5 \cos(y)$$

$$\Omega = [0, 2\pi]^2, T = 0.4.$$

$$\mathcal{M} = 10^{-3}, \lambda = 1$$

# Numerical Experiment I: First Order Convergence

- $\varphi_{256}$ : “exact” solution

$N$	$\ \varphi_{256} - \varphi_N\ _{2,h}$	rate	$\ \varphi_{256} - \varphi_N\ _{L^2}$	rate
8	6.1911	–	0.2625	–
16	1.9293	1.6045	0.0624	2.1039
32	0.5601	1.7221	0.0151	2.0685
64	0.1599	1.7516	0.0035	2.1336
128	0.0461	1.7354	0.0008	2.1887

**Table:** Errors and convergence rates of the  $C^0$ -IP method with  $\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ ,  $h = 2\sqrt{2}\pi/N$ ,  $\tau = 0.05/N$ .

# Numerical Experiment I: First Order Convergence

- $\varphi_{256}$ : “exact” solution
- $N$  indicates the number of sub-intervals per side of  $\Omega$

$N$	$\ \varphi_{256} - \varphi_N\ _{2,h}$	rate	$\ \varphi_{256} - \varphi_N\ _{L^2}$	rate
8	6.1911	–	0.2625	–
16	1.9293	1.6045	0.0624	2.1039
32	0.5601	1.7221	0.0151	2.0685
64	0.1599	1.7516	0.0035	2.1336
128	0.0461	1.7354	0.0008	2.1887

**Table:** Errors and convergence rates of the  $C^0$ -IP method with  $\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ ,  $h = 2\sqrt{2}\pi/N$ ,  $\tau = 0.05/N$ .

# Numerical Experiment I: First Order Convergence

- $\varphi_{256}$ : “exact” solution
- $N$  indicates the number of sub-intervals per side of  $\Omega$
- $\tau = 0.05/N$

$N$	$\ \varphi_{256} - \varphi_N\ _{2,h}$	rate	$\ \varphi_{256} - \varphi_N\ _{L^2}$	rate
8	6.1911	–	0.2625	–
16	1.9293	1.6045	0.0624	2.1039
32	0.5601	1.7221	0.0151	2.0685
64	0.1599	1.7516	0.0035	2.1336
128	0.0461	1.7354	0.0008	2.1887

**Table:** Errors and convergence rates of the  $C^0$ -IP method with  $\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ ,  $h = 2\sqrt{2}\pi/N$ ,  $\tau = 0.05/N$ .

# Numerical Experiment II: Energy Dissipation

## Example

$$\varphi_0(x, y) = 0.3 \cos(3x) + 0.5 \cos(y)$$

$$\Omega := [0, 10]^2, T = 5.$$

$\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ , penalty parameter  $\alpha = 8$

$$F(\varphi) = \frac{\beta}{2} \|\varphi\|_{L^6}^6 - \frac{3\beta}{4} \|\varphi\|_{L^4}^4 + \frac{\beta|\Omega|}{4} + \frac{1}{2} \|\varphi \nabla \varphi\|_{L^2}^2 - \frac{a_0}{2} \|\nabla \varphi\|_{L^2}^2 + \frac{\lambda}{2} a_h^{IP}(\varphi, \varphi).$$

# Numerical Experiment II: Energy Dissipation

## Example

$$\varphi_0(x, y) = 0.3 \cos(3x) + 0.5 \cos(y)$$

$$\Omega := [0, 10]^2, T = 5.$$

$\mathcal{M} = 10^{-3}$ ,  $\lambda = 1$ , penalty parameter  $\alpha = 8$

$$F(\varphi) = \frac{\beta}{2} \|\varphi\|_{L^6}^6 - \frac{3\beta}{4} \|\varphi\|_{L^4}^4 + \frac{\beta|\Omega|}{4} + \frac{1}{2} \|\varphi \nabla \varphi\|_{L^2}^2 - \frac{a_0}{2} \|\nabla \varphi\|_{L^2}^2 + \frac{\lambda}{2} a_h^{IP}(\varphi, \varphi).$$

Track the scaled energy  $F(\varphi) - \frac{\beta|\Omega|}{4}$  for time step sizes  $\tau = 0.5, 0.25, 0.0125$ , and 0.0625.

# Numerical Experiment II: Energy Dissipation

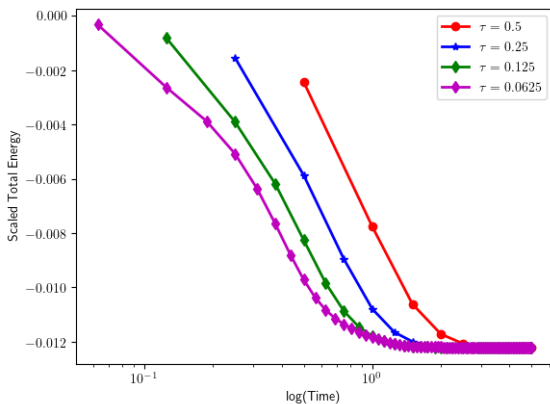


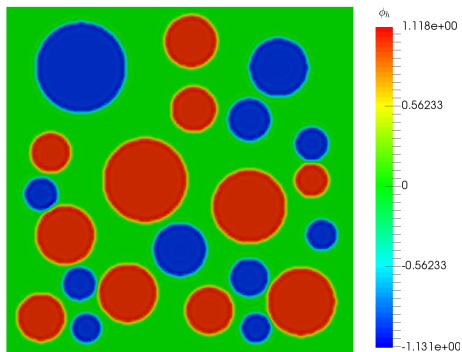
Figure: The time evolution of the scaled total energy  $F(\varphi) - \frac{\beta|\Omega|}{4}$ ,  $h = 10\sqrt{2}/128$



# Numerical Experiment III: Microemulsions Simulation

## Example

$$\Omega = (-5, 5)^2, T = 0.1, \tau = 1.1 \times 10^{-4}, h = 10\sqrt{2}/128, \mathcal{M} = 10, \lambda = 10^{-2}.$$



# Evolution of the profile

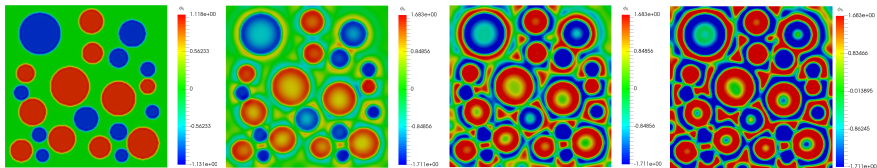


Figure: Profiles at  $m = 0, 1, 2$  and  $3$ .

# Evolution of the profile

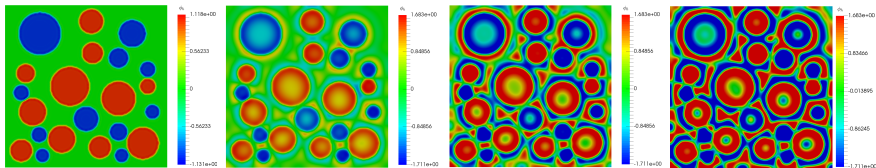


Figure: Profiles at  $m = 0, 1, 2$  and  $3$ .

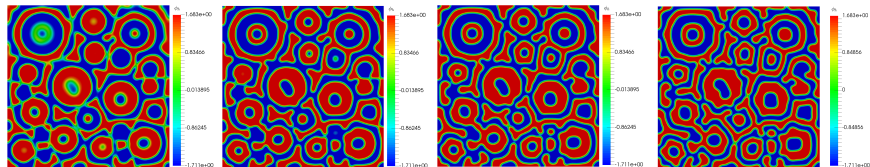


Figure: Profiles at  $m = 4, 11, 20$  and  $25$ .

# Evolution of the profile

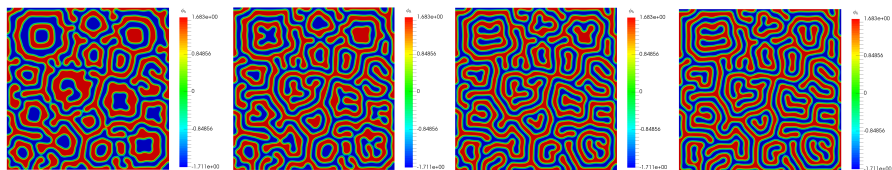


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

# Evolution of the profile

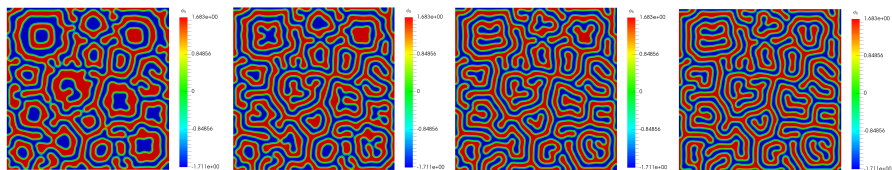


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.

# Evolution of the profile

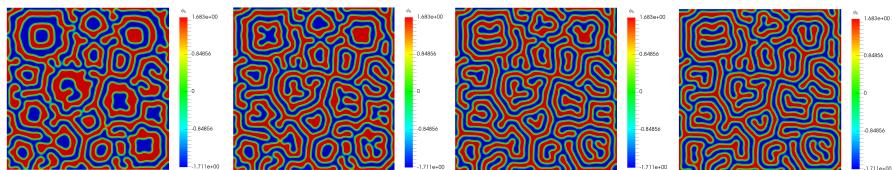


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.
- Optimal time steps for each stage can differ by several orders of magnitude.

# Evolution of the profile

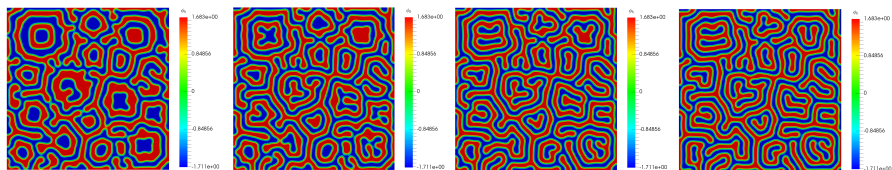


Figure: Profiles at  $m = 37, 74, 158$  and  $511$ .

- Different temporal scales capture different stages of phase-field evolution.
- Optimal time steps for each stage can differ by several orders of magnitude.
- Time-step adaptivity is crucial for accuracy and efficiency of the numerical scheme.

# Conclusions and Ongoing Work



# Conclusions and Ongoing Work

- Numerical schemes using the  $C^0$ -IP framework were presented for the sixth-order phase field models.

**Novel contribution:** provide an error analysis of the  $C^0$ -IP framework.

# Conclusions and Ongoing Work

- Numerical schemes using the  $C^0$ -IP framework were presented for the sixth-order phase field models.  
**Novel contribution:** provide an error analysis of the  $C^0$ -IP framework.
- **Open challenge:** Little is known a priori about the dynamics of the system thus making the task of choosing optimal time step and mesh size parameters difficult.

# Conclusions and Ongoing Work

- Numerical schemes using the  $C^0$ -IP framework were presented for the sixth-order phase field models.  
**Novel contribution:** provide an error analysis of the  $C^0$ -IP framework.
- **Open challenge:** Little is known a priori about the dynamics of the system thus making the task of choosing optimal time step and mesh size parameters difficult.
- **Focus:** Derive a framework which automatically adapts the choice of the method parameters in response to the change in the dynamics of the problem.

# Conclusions and Ongoing Work

- Numerical schemes using the  $C^0$ -IP framework were presented for the sixth-order phase field models.  
**Novel contribution:** provide an error analysis of the  $C^0$ -IP framework.
- **Open challenge:** Little is known a priori about the dynamics of the system thus making the task of choosing optimal time step and mesh size parameters difficult.
- **Focus:** Derive a framework which automatically adapts the choice of the method parameters in response to the change in the dynamics of the problem.

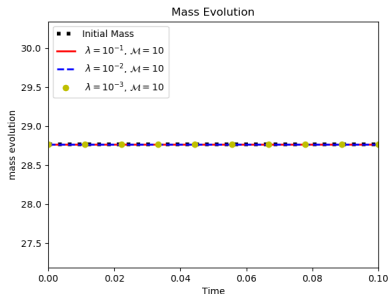
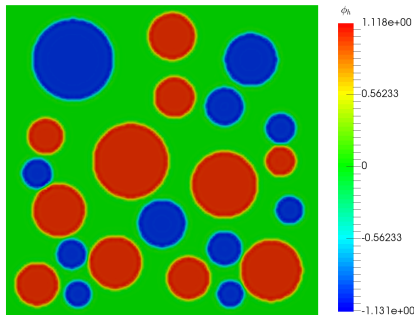
Thank you for your attention!

# Numerical Experiment: Discrete Mass Conservation

Initial Conditions:

## Example

$\Omega = (-5, 5)^2$ ,  $T = 0.1$ ,  $\tau = 1.1 \times 10^{-4}$ ,  $h = 10\sqrt{2}/128$ ,  $\mathcal{M} = 10$ ,  
 $\lambda = 10^{-1}, 10^{-2}, 10^{-3}$ .



# Effect of the decreasing $\lambda$

Theory suggests that increasing  $\lambda$  guarantees the existence and stability of the solution.

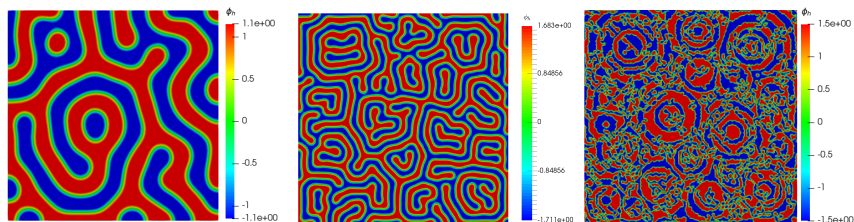


Figure:  $\lambda = 10^{-1}, 10^{-2}, 10^{-3}$

$$\Omega = (-5, 5)^2, T = 0.1, \tau = 1.1 \times 10^{-4}, h = 10\sqrt{2}/128, \mathcal{M} = 10,$$